# SAMPLING PER MODE SIMULATION FOR SWITCHING DIFFUSIONS 

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- Switching jump diffusion
- Splitting technique
- Some issues
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- Multilevel Feynman-Kac distributions
- Dynamical evolution
(3) SAMPLING PER MODE ALGORITHM
- Particle Methods
- Sampling per Mode algorithm
(4) Asymptotic Behaviour
- Asymptotic Behaviour
- Law of Large Numbers
- Central Limit Theorem
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(3) CONCLUSION
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- We use a splitting technique adapted to the context of switching diffusions: the sampling per mode algorithm introduced by Krystul in [Krystul, 2006]

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## SWITCHING JUMP DIFFUSION

- Strong Markov process $Z=\left\{\left(X_{t}, \theta_{t}\right) ; t \geq 0\right\}$ with value in $\mathbb{R}^{d} \times \mathbb{M}$ with a finite set $\mathbb{M}=\{1, \cdots, M\}$,
- Strong Markov process $Z=\left\{\left(X_{t}, \theta_{t}\right) ; t \geq 0\right\}$ with value in $\mathbb{R}^{d} \times \mathbb{M}$ with a finite set $\mathbb{M}=\{1, \cdots, M\}$,
- the continuous component is described as a d-dimensional SDE

$$
d X_{t}=b\left(X_{t}, \theta_{t}\right) d t+\sigma\left(X_{t}, \theta_{t}\right) d B_{t}
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- and the discrete mode as a pure jump process

$$
\mathbb{P}\left(\theta_{t+\Delta t}=j \mid \theta_{t}=i, X_{t}=x\right)=\lambda_{i j}(x) \Delta t+o(\Delta t), i \neq j,
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with jump intensities depending on the continuous component.

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■ Let $A \subset \mathbb{R}^{d}$ be a closed critical region in which $X_{t}$ could enter but with a very small probability.

- If $T_{A}$ denotes the hitting time of $A$, we would like to estimate $\mathbb{P}\left(T_{A} \leq T\right)$ with $T$ a deterministic or a stopping time.


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■ With $B=A \times \mathbb{M}$ and $B_{k}=D_{k} \times \mathbb{M}$, we define for $k=1, \cdots, n$

$$
T_{k}=\inf \left\{t \geq 0: Z_{t} \in B_{k}\right\}=\inf \left\{t \geq 0: X_{t} \in D_{k}\right\}
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which satisfy $0=T_{0} \leq T_{1} \leq \cdots \leq T_{n}=T_{B}$.

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- Then

$$
\mathbb{P}\left(T_{A} \leq T\right)=\mathbb{P}\left(T_{B} \leq T\right)=\prod_{k=1}^{n} \mathbb{P}\left(T_{k} \leq T \mid T_{k-1} \leq T\right)
$$

where conditional probabilities are not very small.

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## SOME ISSUES

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- to sample more particles from mode with higher probability,
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- Increasing the number of particles should improve the estimate but only at the cost of increased simulation time,
- Idea: keep constant the number of particles in each visited mode at each resampling step,
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## MULTILEVEL FEYNMAN-KAC DISTRIBUTIONS

- To capture the behaviour of $Z$ between each thresholds, we consider the random excursions $\mathcal{Z}_{k}$ of $Z$ between $T_{k-1}$ and $T_{k} \wedge T$

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- and we introduce the selection functions,

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g_{k}\left(\mathcal{Z}_{k}\right)=1_{\left\{Z_{T_{k} \wedge T} \in B_{k}\right\}}, \quad g_{k}^{j}\left(\mathcal{Z}_{k}\right)=1_{\left\{Z_{T_{k} \wedge T} \in D_{k} \times\{j\}\right\}}, \quad j \in \mathbb{M},
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1_{\left\{T_{k} \leq T\right\}}=g_{k}\left(\mathcal{Z}_{k}\right), \text { and } 1_{\left\{T_{k} \leq T, \theta_{T_{k}}=j\right\}}=g_{k}^{j}\left(\mathcal{Z}_{k}\right)
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- We can interpret the rare event probability in terms of the Feynman-Kac measures defined by
$\gamma_{k}(f)=\mathbb{E}\left[f\left(\mathcal{Z}_{k}\right) g_{k-1}\left(\mathcal{Z}_{k-1}\right)\right]=\mathbb{E}\left[f\left(\left(X_{t}, \theta_{t}\right), T_{k-1} \leq t \leq T_{k} \wedge T\right) 1_{\left\{T_{k-1} \leq T\right\}}\right]$
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- and the corresponding normalized measures defined by

$$
\begin{aligned}
& \eta_{k}(f)=\frac{\gamma_{k}(f)}{\gamma_{k}(1)}=\mathbb{E}\left[f\left(\left(X_{t}, \theta_{t}\right), T_{k-1} \leq t \leq T_{k} \wedge T\right) \mid T_{k-1} \leq T\right] \\
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- In particular, for $f \equiv 1$

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- and for $f=g_{k}$ or $f=g_{k}^{j}$, it holds

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- We have the key formulas

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- Then, we recover

$$
\mathbb{P}\left(T_{n} \leq T\right)=\prod_{p=0}^{n} \eta_{p}\left(g_{p}\right)
$$

## MULTILEVEL FEYNMAN-KAC DISTRIBUTIONS

■ In order to keep trace of the discrete mode, we construct for any $j \in \mathbb{M}$

$$
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& \gamma_{k}^{j}(f)=\mathbb{E}\left[f\left(\mathcal{Z}_{k}\right) g_{k-1}^{j}\left(\mathcal{Z}_{k-1}\right)\right]=\mathbb{E}\left[f\left(Z_{t}, T_{k-1} \leq t \leq T_{k} \wedge T\right) 1_{\left\{T_{k-1} \leq T, \theta_{T_{k-1}}=j\right\}}\right] \\
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$$

- We have the decompositions

$$
\widehat{\eta}_{k}=\sum_{j \in \mathbb{M}} \omega_{k}^{j} \widehat{\eta}_{k}^{\prime}, \quad \eta_{k+1}=\sum_{j \in \mathbb{M}} \omega_{k}^{j} \eta_{k+1}^{j},
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where

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■ Using the Markov property of $\mathcal{Z}$ (with Markov kernel $\mathcal{M}_{k}$ ), we obtain

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- and the nonlinear measure-valued transformations

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\widehat{\eta}_{k}(f)=\frac{\eta_{k}\left(f g_{k}\right)}{\eta_{k}\left(g_{k}\right)}:=\psi_{k}\left(\eta_{k}\right)(f), \quad \widehat{\eta}_{k}^{j}(f)=\frac{\eta_{k}\left(f g_{k}^{j}\right)}{\eta_{k}\left(g_{k}^{j}\right)}:=\psi_{k}^{j}\left(\eta_{k}\right)(f)
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- Using the Markov property of $\mathcal{Z}$ (with Markov kernel $\mathcal{M}_{k}$ ), we obtain

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- so, the following two separate selection/mutation transitions

$$
\eta_{k} \xrightarrow{\text { selection }} \widehat{\eta}_{k}:=\Psi_{k}\left(\eta_{k}\right) \xrightarrow{\text { mutation }} \eta_{k+1}=\widehat{\eta}_{k} \mathcal{M}_{k+1} .
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by

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\eta_{k}^{N}:=\frac{1}{N} \sum_{i=1}^{N} \delta_{\xi_{k}^{i}} \xrightarrow{\text { selection }} \widehat{\eta}_{k}^{N}:=\frac{1}{N} \sum_{i=1}^{N} \delta_{\widehat{\xi}_{k}^{i}} \xrightarrow{\text { mutation }} \eta_{k+1}=\frac{1}{N} \sum_{i=1}^{N} \delta_{\xi_{k+1}^{i}}
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■ Obviously, the total number of particles can change at each time some mode is not visited, or empty mode is visited afresh.
- Let $\widehat{N}_{k}$ and $N_{k}$ denote the total numbers of particles $\widehat{\xi}_{k}$ and $\xi_{k}$, and $\omega_{k}^{j, N}$ the weights associated with the modes, we have the evolution scheme

$$
\left(N_{k},\left(\omega_{k-1}^{j, N}\right)_{j \in J_{k-1}}, \xi_{k}\right) \rightarrow\left(\widehat{N}_{k},\left(\omega_{k}^{j}\right)_{j \in J_{k}}, \widehat{\xi}_{k}\right) \rightarrow\left(N_{k+1},\left(\omega_{k}^{j, N}\right)_{j \in J_{k}}, \xi_{k+1}\right)
$$

where $J_{k}$ denotes the set of non empty modes at step $k$

## INITIALIZATION

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- Here $J_{0}^{j}$ is the set of the indices of the particles in the mode $j$.
- If $\widehat{N}_{k}=0$ the particle system dies, otherwise independently of each other, each particle $\widehat{\xi}_{k}^{\kappa}$ evolves randomly according to the Markov transition $\mathcal{M}_{k+1}$

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- So $\eta_{k+1}^{N}=\sum_{j \in J_{k}} \omega_{k}^{j, N} \eta_{k+1}^{j, N}$, with $\eta_{k+1}^{j, N}=\frac{1}{N^{j}} \sum_{\kappa \in J_{k}^{j}} \delta_{\xi_{k+1}^{k}}$, where $J_{k}^{j}$ is the set of the labels of the particles in mode $j \in J_{k}$.


## SELECTION/RESAMPLING $\left(N_{k+1}, \omega_{k}^{N}, \xi_{k+1}\right) \rightarrow\left(\widehat{N}_{k+1}, \omega_{k+1}^{N}, \widehat{\xi}_{k+1}\right)$

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- The total number of particles is $\widehat{N}_{k+1}=\sum_{j \in J_{k+1}} N^{j}$.




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## MAIN THEOREMS: LAW OF LARGE NUMBERS

## THEOREM (LAW OF LARGE NUMBERS)

For any $n \geq 0$ and any bounded function $f$, we have

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## Theorem (CEntral Limit Theorem)

Let $N \rightarrow \infty$ in such a way that $\rho_{j}=N^{j} / N$ are "preserved" for all $j \in \mathbb{M}$. Then, the random variable

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\sqrt{N}\left(1_{\left\{N_{n+1}>0\right\}} \gamma_{n+1}^{N}(1)-\mathbb{P}\left(T_{n}<T\right)\right)
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$$
\Omega_{q}=\sum_{j \in \mathbb{M}}\left(\omega_{q-1}^{j}\right)^{2} \rho_{j}^{-1}=1+\chi^{2}\left(\omega_{q-1}, \rho\right),
$$

and

$$
\Delta_{q}^{n}(t, z)=\mathbb{P}\left(T_{n} \leq T \mid T_{q}=t, Z_{T_{q}}=z\right) .
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## CONCLUSION

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- A better comprehension of the expression of $W_{n+1}$ could help the choice of the $N^{j}$ regarding the cost of the algorithm.
- This algorithm is implemented in a software developed by National Aerospace Laboratory (NLR) and used to evaluate the safety characteristics of an arbitrary (new) operational Air Traffic Management concept [Blom, 2009].


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