# On the Stability and the Approximation of Branching Distribution Flows, with Applications to Nonlinear Multiple Target Filtering* 

F. Caron, P. Del Moral, M. Pace ${ }^{\dagger}$ B.-N. Vo ${ }^{\ddagger}$

October 6, 2010


#### Abstract

We analyse the exponential stability properties of a class of measure-valued equations arising in nonlinear multi-target filtering problems. We also prove the uniform convergence properties w.r.t. the time parameter of a rather general class of stochastic filtering algorithms, including sequential Monte Carlo type models and mean field particle interpretation models. We illustrate these results in the context of the Bernoulli and the Probability Hypothesis Density filter, yielding what seems to be the first results of this kind in this subject.

Mathematics Subject Classification: Primary: 93E11, 65C35, 37L15; Secondary: $60 \mathrm{~J} 85,65 \mathrm{C} 05,65 \mathrm{C} 35$.

Keywords: Nonlinear multi-target filtering, Bernoulli filter, Probability hypothesis density filter, interacting particle systems, particle filters, concentration inequalities, semigroup stability, functional contraction inequalities.


## 1 Introduction

Let $\left(E_{n}\right)_{n \geq 0}$ be a sequence of measurable spaces equipped with the $\sigma$-fields $\left(\mathcal{E}_{n}\right)_{n \geq 0}$, and for each with $n \geq 0$, denote $\mathcal{M}\left(E_{n}\right), \mathcal{M}_{+}\left(E_{n}\right)$ and $\mathcal{P}\left(E_{n}\right)$ the set of all finite signed measures, the subset of positive measures and the subset of probability measures, respectively, over the space $E_{n}$. The aim of this work is to present a stochastic interacting particle interpretation for numerical solutions of the general measure-valued dynamical systems $\gamma_{n} \in \mathcal{M}_{+}\left(E_{n}\right)$ defined by the following non-linear equation

$$
\begin{equation*}
\gamma_{n}\left(d x_{n}\right)=\left(\gamma_{n-1} Q_{n, \gamma_{n-1}}\right)\left(d x_{n}\right):=\int_{E_{n-1}} \gamma_{n-1}\left(d x_{n-1}\right) Q_{n, \gamma_{n-1}}\left(x_{n-1}, d x_{n}\right) \tag{1.1}
\end{equation*}
$$

[^0]with initial measure $\gamma_{0} \in \mathcal{M}_{+}\left(E_{0}\right)$, and positive and bounded integral operators $Q_{n, \gamma}$ from $E_{n-1}$ into $E_{n}$, indexed by the time parameter $n \geq 1$ and the set of measures $\gamma \in \mathcal{M}_{+}\left(E_{n}\right)$.

This class of measure-valued equations arises in a natural way in the analysis of the first moments evolution of nonlinear branching processes, as well as in signal processing and more particularly in multiple targets tracking models. A pair of filtering models is discussed in some details in section 1.1.1 and in section 1.1.2. In the context of multiple targets tracking problems these measure-valued equations represents the first-order statistical moments of the conditional distributions of the target occupation measures given observation random measures obscured by clutter, detection uncertainty and data association uncertainty.

As in most of the filtering problems encountered in practice, the initial distribution of the targets is usually unknown. It is therefore essential to check wether or not the filtering equation "forgets" any erroneous initial distribution. For a thorough discussion on the stability properties of traditional nonlinear filtering problems with a detailed overview of theoretical developments on this subject, we refer to the book [6] and to the more recent article by M. L. Kleptsyna and A. Y. Veretennikov [15]. Besides the fact that significant progress has been made in the recent years in the rigorous derivation of multiple target tracking nonlinear equations (see for instance [4, 17, [27, 22]), up to our knowledge the stability and the robustness properties of these measure-valued models have never been addressed so far in the literature on the subject. One aim of this paper is to study one such important property: the exponential stability properties of multiple target filtering models. We present an original and general perturbation type technique combining the continuity property and the stability analysis of nonlinear semigroups of the form (1.1). A more thorough presentation of these results is provided in section 1.2 dedicated to the statement of the main results of the present article. The detailed presentation of this perturbation technique can be found in section 3 .

On the other hand, while the integral equation (1.1) appears to be simple at first glance, numerical solutions are computationally intensive, often requiring integrations in high dimensional spaces. One natural way to solve the non-linear integral equation (1.1) is to use find a judicious probabilistic interpretation of the normalized distributions flow given below

$$
\eta_{n}\left(d x_{n}\right):=\gamma_{n}\left(d x_{n}\right) / \gamma_{n}(1)
$$

To describe with some conciseness these stochastic models, it is important to observe that the pair process $\left(\gamma_{n}(1), \eta_{n}\right) \in\left(\mathbb{R}_{+} \times \mathcal{P}\left(E_{n}\right)\right)$ satisfies an evolution equation of the following form

$$
\begin{equation*}
\left(\gamma_{n}(1), \eta_{n}\right)=\Gamma_{n}\left(\gamma_{n-1}(1), \eta_{n-1}\right) \tag{1.2}
\end{equation*}
$$

Let the mappings $\Gamma_{n}^{1}: \mathbb{R}_{+} \times \mathcal{P}\left(E_{n}\right) \rightarrow \mathbb{R}_{+}$and $\Gamma_{n}^{2}: \mathbb{R}_{+} \times \mathcal{P}\left(E_{n}\right) \rightarrow \mathcal{P}\left(E_{n}\right)$, denote the first and the second components of $\Gamma_{n}$ respectively. By construction, we notice that the total mass process can be computed using the recursive formula

$$
\begin{equation*}
\gamma_{n+1}(1)=\gamma_{n}\left(G_{n, \gamma_{n}}\right)=\eta_{n}\left(G_{n, \gamma_{n}}\right) \gamma_{n}(1) \quad \text { with } \quad G_{n, \gamma_{n}}:=Q_{n+1, \gamma_{n}}(1) \tag{1.3}
\end{equation*}
$$

Suppose that we are given an approximation $\left(\gamma_{n}^{N}(1), \eta_{n}^{N}\right)$ of the pair $\left(\gamma_{n}(1), \eta_{n}\right)$ at some time horizon $n$, where $N$ stands for some precision parameter; that is $\left(\gamma_{n}^{N}(1), \eta_{n}^{N}\right)$ converges (in some sense) to $\left(\gamma_{n}(1), \eta_{n}\right)$, as $N \rightarrow \infty$. Then, the $N$-approximation of the measure $\gamma_{n}$ is given by $\gamma_{n}^{N}=\gamma_{n}^{N}(1) \times \eta_{n}^{N}$. The central idea behind any approximation model is to ensure
that the total mass process at time $(n+1)$ defined by

$$
\begin{equation*}
\gamma_{n+1}^{N}(1)=\eta_{n}^{N}\left(G_{n, \gamma_{n}^{N}}\right) \gamma_{n}^{N}(1) \tag{1.4}
\end{equation*}
$$

can be "easily" computed in terms of the $N$-approximation measures $\gamma_{n}^{N}$. Assuming that the initial mass $\gamma_{0}(1)=\gamma_{0}^{N}(1)$ is known, the next step is to find some strategy to approximate the quantities $\Gamma_{n+1}^{2}\left(\gamma_{n}^{N}(1), \eta_{n}^{N}\right)$ by some $N$-approximation measures $\eta_{n+1}^{N}$, and to set $\gamma_{n+1}^{N}=$ $\gamma_{n+1}^{N}(1) \times \eta_{n+1}^{N}$. The local fluctuations of $\eta_{n}^{N}$ around the measures $\Gamma_{n}^{2}\left(\gamma_{n-1}^{N}(1), \eta_{n-1}^{N}\right)$ is defined in terms of a collection of random fields $W_{n}^{N}$ :

$$
\begin{equation*}
W_{n}^{N}:=\sqrt{N}\left[\eta_{n}^{N}-\Gamma_{n}^{2}\left(\gamma_{n-1}^{N}(1), \eta_{n-1}^{N}\right)\right] \Longleftrightarrow \eta_{n}^{N}=\Gamma_{n}^{2}\left(\gamma_{n-1}^{N}(1), \eta_{n-1}^{N}\right)+\frac{1}{\sqrt{N}} W_{n}^{N} \tag{1.5}
\end{equation*}
$$

which satisfies for any $r \geq 1$ and any test function $f$ with uniform norm $\|f\| \leq 1$,

$$
\begin{equation*}
\mathbb{E}\left(\left|W_{n}^{N}(f)\right|^{r} \mid \mathcal{F}_{n-1}^{N}\right)^{\frac{1}{r}} \leq a_{r} \tag{1.6}
\end{equation*}
$$

where $\mathcal{F}_{n-1}^{N}=\sigma\left(\eta_{p}^{N}, 0 \leq p<n\right)$ is the $\sigma$-field generated by the random measures $\eta_{p}^{N}, 0 \leq$ $p<n$, while $b$ and $a_{r}$ are universal constants whose values do not depend on the precision parameter $N$. The stochastic analysis of the resulting particle approximation model relies on the analysis of the propagation of the local sampling errors defined in (1.5). The main objective is to control, at any time horizon $n$, the fluctuations of the random measures $\left(\gamma_{n}^{N}, \eta_{n}^{N}\right)$ around their limiting values $\left(\gamma_{n}, \eta_{n}\right)$ defined by the following random fields:

$$
\begin{equation*}
V_{n}^{\gamma, N}:=\sqrt{N}\left[\gamma_{n}^{N}-\gamma_{n}\right] \quad \text { with } \quad V_{n}^{\eta, N}:=\sqrt{N}\left[\eta_{n}^{N}-\eta_{n}\right] \tag{1.7}
\end{equation*}
$$

The construction of the $N$-approximation measures $\eta_{n}^{N}$ is far from being unique. In the present article, we devise three different classes of stochastic particle approximation models. These stochastic algorithms are discussed in section 4. The first one is a mean field particle interpretation of the flow of probability measures $\eta_{n}$, and it is presented in section 4.1. The second model is an interacting particle association model while the third one is a combination of these two approximation algorithms. These pair of approximation models are respectively discussed in section 4.2 and in section 4.3. In the context of multi-target tracking models, the first two approximation models are closely related to the the sequential Monte Carlo technique presented in the series of articles [23, 24, 28, 29, 30, 31], and respectively, the Gaussian mixture Probability Hypothesis Density filter discussed in the article by B.-N. Vo, and W.-K. Ma [25, 26], and the the Rao-Blackwellized Particle multi-target filters presented by S. Sarkka, A. Vehtari, and J. Lampinen in [20, 21]. These modern stochastic algorithms are rather simple to implement and computationally tractable, and they exhibit excellent performance.

Nevertheless, despite advances in recent years [3, 14, 24], these Monte Carlo particle type multi-target filters remain poorly understood theoretically. One aim of this article is to present a novel class of stochastic algorithms with a refined analysis including uniform convergence results w.r.t. the time parameter. We also illustrate these results in the context of multi-target tracking models, yielding what seems to be the first uniform results of this type in this subject.

The rest of the article is organized as follows: In section 1.1 we illustrate the abstract measure-valued equations (1.1) with two recent multi-target filters models, namely
the Bernoulli filter and the Probability Hypothesis Density filter (abbreviate PHD filter). Section 1.2 is devoted to the statement of our main results. In section 2, we describe the semigroups and the continuity properties of the nonlinear equation 1.1. We show that this semigroup analysis can be applied to analyse the convergence of the Bernoulli and the PHD approximation filters. Section 3 is devoted to the stability properties of nonlinear measurevalued processes of the form (1.2). We present a perturbation technique and a series of functional contraction inequalities. In the next three sections, we illustrate these results in the context of Feynman-Kac models, as well as Bernoulli and PHD models. Section 4 is concerned with the detailed presentation and the convergence analysis of three different classes of particle type approximation models, including mean field type particle approximations and particle association stochastic algorithms. Finally, the appendix of the article contains most of technical proofs in the text.

### 1.1 Measure-valued systems in Multi-target tracking

The measure-valued process given by (1.1) is a generalisation of Feynman-Kac measures. Its continuous time version naturally arise in the modeling and analysis of the first moments of spatial branching process [6, 11].

Our major motivation for studying this class of measure-valued system stems from advanced signal processing, more specifically, multiple target tracking. Driven primarily in the early 1970's by aerospace applications such as radar, sonar, guidance, navigation, and air traffic control, today multi-target filtering has found applications in many diverse disciplines, see for example the texts [1], [2] [18] and references therein. These nonlinear filtering problems deal with jointly estimating the number and states of several interacting targets given a sequence of partial observations corrupted by noise, false measurements as well as miss-detection. This rapidly developing subject is, arguably, one of the most interesting contact points between the theory of spatial branching processes, mean field particle systems and advanced signal processing.

The first connections between stochastic branching processes and multi-target tracking seem to go back to the article by S. Mori, et. al. [19] published in 1986. However it was Mahler's systematic treatment of multi-sensor multi-target filtering using random finite sets theory [13, 12, 16, 17] that lead to the development novel multi-target filters and sparked world wide interests. To motivate the article, we briefly outline two recent multi-target filters that do not fit the standard Feynman-Kacs framework, but fall under the umbrella of the measure-valued equation (1.1). The first is the Bernoulli filter for joint detection and tracking of a single target while the second is the Probability Hypothesis Density filter.

### 1.1.1 Bernoulli filtering

A basic problem in target tracking is that the target of interest may not always be present and exact knowledge of target existence/presence cannot be determined from observations due to clutter and detection uncertainty [18]. The Bernoulli filter is a generalisation of the standard Bayes filter, which accommodates presence and absence of the target [27]. In a Bernoulli model, the birth of the target at time $n+1$ is modelled by a measure $\mu_{n+1}$ on $E_{n+1}$. The target enters the scene with a probability $\mu_{n+1}(1)<1$ and its state is distributed according to the normalised measure $\mu_{n+1} / \mu_{n+1}(1)$. At time $n$, a target $X_{n}$ has
a probability $s_{n}\left(X_{n}\right)$ of surviving to the next time and evolve to a new state according to a given elementary Markov transition $M_{n+1}$ from $E_{n}$ into $E_{n+1}$. At time $n+1$, the target (if it exists) generates with probability $d_{n+1}\left(X_{n+1}\right)$ an observation $Y_{n+1}$ on some auxiliary state space, say $E_{n+1}^{Y}$ with likelihood function $l_{n+1}\left(X_{n+1}, y\right)$. This so-called Bernoulli observation point process is superimposed with an additional and independent Poisson point process with intensity function $h_{n}>0$ to form the occupation (or counting) measure observation process $\mathcal{Y}_{n+1}=\sum_{1 \leq i \leq N_{n+1}^{Y}} \delta_{Y_{n+1}^{i}}$.

In its original form, the Bernoulli filter jointly propagates the probability existence of the target and the distribution of the target state [27]. Combining the probability of existence and the state distribution into a single measure, it can be shown that the Bernoulli filter satisfies the integral equation (1.1), with the probability of existence of the target given by the mass $\gamma_{n}(1)$ and the distribution of the target state given by the normalised measure $\eta_{n}=\gamma_{n} / \gamma_{n}(1)$. The integral operator for the Bernoulli filter takes the following form

$$
\begin{equation*}
Q_{n+1, \gamma_{n}}\left(x_{n}, d x_{n+1}\right):=\frac{s_{n}\left(x_{n}\right) g_{n}\left(x_{n}\right) M_{n+1}\left(x_{n}, d x_{n+1}\right)+\left(\gamma_{n}(1)^{-1}-1\right) \mu_{n+1}\left(d x_{n+1}\right)}{\left(1-\gamma_{n}(1)\right)+\gamma_{n}\left(g_{n}\right)} \tag{1.8}
\end{equation*}
$$

where $g_{n}$ is a likelihood function given by

$$
\begin{equation*}
g_{n}\left(x_{n}\right):=\left(1-d_{n}\left(x_{n}\right)\right)+d_{n}\left(x_{n}\right) \mathcal{Y}_{n}\left(l_{n}\left(x_{n}, \cdot\right) / h_{n}\right) \tag{1.9}
\end{equation*}
$$

### 1.1.2 PHD filtering

A more challenging problem arises when the number of targets varies randomly in time, obscured by clutter, detection uncertainty and data association uncertainty. Suppose that at a given time $n$ there are $N_{n}^{X}$ targets $\left(X_{n}^{i}\right)_{1 \leq i \leq N_{n}^{X}}$ each taking values in some measurable state space $E_{n}$. A target $X_{n}^{i}$, at time $n$, survives to the next time step with probability $s_{n}\left(X_{n}^{i}\right)$ and evolves to a new state according to a given elementary Markov transition $M_{n+1}^{\prime}$ from $E_{n}$ into $E_{n+1}$. In addition $X_{n}^{i}$ can spawn new targets at the next time, usually modelled by a spatial Poisson process with intensity measure $B_{n+1}\left(X_{n}^{i}, \cdot\right)$ on $E_{n+1}$. At the same time, an independent collection of new targets is added to the current configuration. This additional and spontaneous branching process is often modeled by a spatial Poisson process with a prescribed intensity measure $\mu_{n+1}$ on $E_{n+1}$. Each target $X_{n+1}^{i}$ generates with probability $d_{n+1}\left(X_{n+1}^{i}\right)$ an observation $Y_{n+1}^{i}$ on some auxiliary state space, say $E_{n+1}^{Y}$, with probability density function $g_{n+1}\left(X_{n+1}^{i}, y\right)$. In addition to this partial observation point process we also observe an additional and independent Poisson point process with intensity function $h_{n}$. Multi-target tracking concerns the estimation of the random measures $\mathcal{X}_{n+1}=\sum_{1 \leq i \leq N_{n}^{X}} \delta_{X_{n}^{i}}$, given the observation occupation measures $\mathcal{Y}_{p}=\sum_{1 \leq i \leq N_{p}^{Y}} \delta_{Y_{p}^{i}}$.

The multi-target tracking problem is computationally intractable in general and the Probability Hypothesis Density PHD (filter), is an approximation that propagates the firstorder statistical moment, or intensity, of the multi-target state forward in time [17]. The PHD filter satisfies the integral equation (1.1), with the integral operator given below

$$
\begin{equation*}
Q_{n+1, \gamma_{n}}\left(x_{n}, d x_{n+1}\right)=g_{n, \gamma_{n}}\left(x_{n}\right) M_{n+1}\left(x_{n}, d x_{n+1}\right)+\gamma_{n}(1)^{-1} \mu_{n+1}\left(d x_{n+1}\right) \tag{1.10}
\end{equation*}
$$

where $M_{n+1}$ is a Markov kernel defined by

$$
\begin{equation*}
M_{n+1}\left(x_{n}, d x_{n+1}\right):=\frac{s_{n}\left(x_{n}\right) M_{n+1}^{\prime}\left(x_{n}, d x_{n+1}\right)+B_{n+1}\left(x_{n}, d x_{n+1}\right)}{s_{n}\left(x_{n}\right)+b_{n}\left(x_{n}\right)} \tag{1.11}
\end{equation*}
$$

with the branching rate $b_{n}\left(x_{n}\right)=B_{n+1}(1)\left(x_{n}\right)$. The likelihood function $g_{n, \gamma_{n}}$ is given by

$$
\begin{equation*}
g_{n, \gamma_{n}}:=r_{n} \times \widehat{g}_{n, \gamma_{n}} \quad \text { with } \quad r_{n}:=\left(s_{n}+b_{n}\right) \tag{1.12}
\end{equation*}
$$

and

$$
\begin{equation*}
\widehat{g}_{n, \gamma_{n}}\left(x_{n}\right):=\left(1-d_{n}\left(x_{n}\right)\right)+d_{n}\left(x_{n}\right) \int \mathcal{Y}_{n}(d y) \frac{g_{n}\left(x_{n}, y_{n}\right)}{h_{n}\left(y_{n}\right)+\gamma_{n}\left(d_{n} g_{n}\left(., y_{n}\right)\right)} \tag{1.13}
\end{equation*}
$$

Since its inception by Mahler [17] in 2003, the PHD filter has attracted substantial interest to date. The development of numerical solutions for the PHD filter [24], [26] have opened the door to numerous novel extensions and applications. More details on the derivation of the PHD filter using random finite sets, Poisson techniques or random measures theoretic approaches can be found in the series of articles [4, 17, 22].

### 1.2 Statement of the main results

To describe with some conciseness the main result of this article, we need to introduce some notation. We let $\operatorname{Osc}_{1}\left(E_{n}\right)$, be the set of $\mathcal{E}_{n}$-measurable functions $f$ on $E_{n}$ with oscillations $\operatorname{osc}(f)=\sup _{x, x^{\prime}}\left|f(x)-f\left(x^{\prime}\right)\right| \leq 1$. We denote by $\mu(f)=\int \mu(d x) f(x)$ the Lebesgue integral of $f$ w.r.t. some measure $\mu \in \mathcal{M}\left(E_{n}\right)$, and we let $\|\mu-\nu\|_{\text {tv }}$ be the total variation distance between two probability measures $\nu$ and $\mu$ on $E_{n}$.

We assume that the following pair of regularity conditions are satisfied.
$\left(H_{1}\right)$ : There exists a series of compact sets $I_{n} \subset(0, \infty)$ such that the initial mass value $\gamma_{0}(1) \in I_{0}$, and for any $m \in I_{n} \eta \in \mathcal{P}\left(E_{n}\right)$, we have

$$
\theta_{-, n}(m) \leq \eta\left(G_{n, m \eta}\right) \leq \theta_{+, n}(m) \text { for some pair of positive functions } \theta_{+-, n} .
$$

The main implication of condition $\left(H_{1}\right)$ comes from the fact that the total mass processes $\gamma_{n}(1)$ and their $N$-approximation models $\gamma_{n}^{N}(1)$ are finite and they evolves at every time $n$ in a series of compact sets

$$
I_{n} \subset\left[m_{n}^{-}, m_{n}^{+}\right] \subset(0, \infty)
$$

with the sequence of parameters $m_{n}^{+/-}$defined by the recursive equations $m_{n+1}^{-}=m_{n}^{-} \theta_{-, n}\left(m_{n}^{-}\right)$ and $m_{n+1}^{+}=m_{n}^{+} \theta_{+, n}\left(m_{n}^{+}\right)$, with the initial conditions $m_{0}^{-}=m_{0}^{+}=\gamma_{0}(1)$.
$\left(H_{2}\right):$ For any $n \geq 1, f \in \operatorname{Osc}_{1}\left(E_{n}\right)$, and any $(m, \eta),\left(m^{\prime}, \eta^{\prime}\right) \in\left(I_{n} \times \mathcal{P}\left(E_{n}\right)\right)$, the one step mappings $\Gamma_{n}=\left(\Gamma_{n}^{1}, \Gamma_{n}^{2}\right)$ defined in (1.2) satisfy the following Lipschitz type inequalities:

$$
\begin{align*}
\left|\Gamma_{n}^{1}(m, \eta)-\Gamma_{n}^{1}\left(m^{\prime}, \eta^{\prime}\right)\right| & \leq c(n)\left|m-m^{\prime}\right|+\int\left|\left[\eta-\eta^{\prime}\right](\varphi)\right| \Sigma_{n,\left(m^{\prime}, \eta^{\prime}\right)}^{1}(d \varphi)  \tag{1.14}\\
\left|\left[\Gamma_{n}^{2}(m, \eta)-\Gamma_{n}^{2}\left(m^{\prime}, \eta^{\prime}\right)\right](f)\right| & \leq c(n)\left|m-m^{\prime}\right|+\int\left|\left[\eta-\eta^{\prime}\right](\varphi)\right| \Sigma_{n,\left(m^{\prime}, \eta^{\prime}\right)}^{2}(f, d \varphi) \tag{1.15}
\end{align*}
$$

for some finite constants $c(n)<\infty$, and some collection of bounded measures $\Sigma_{n,\left(m^{\prime}, \eta^{\prime}\right)}^{1}$ and $\Sigma_{n,\left(m^{\prime}, \eta^{\prime}\right)}^{2}(f,$.$) on \mathcal{B}\left(E_{n}\right)$ such that

$$
\int \operatorname{osc}(\varphi) \Sigma_{n,(m, \eta)}^{1}(d \varphi) \leq \delta\left(\Sigma_{n}^{1}\right) \quad \text { and } \quad \int \operatorname{osc}(\varphi) \Sigma_{n,(m, \eta)}^{2}(f, d \varphi) \leq \delta\left(\Sigma_{n}^{2}\right)
$$

for some finite constant $\delta\left(\Sigma_{n}^{i}\right)<\infty, i=1,2$, whose values do dot depend on the parameters $(m, \eta) \in\left(I_{n} \times \mathcal{P}\left(E_{n}\right)\right.$ and $f \in \operatorname{Osc}_{1}\left(E_{n}\right)$.

Condition $\left(H_{2}\right)$ is a rather basic and weak continuity type property. It states that the one step transformations of the flow of measures (1.2) are weakly Lipschitz, in the sense that the mass variations and the integral differences w.r.t. some test function $f$ can be controlled by the different initial masses and measures w.r.t. a collection of integrals of a possibly infinite number of test functions. It is satisfied for a large class of one step transformations $\Gamma_{n}$. In section 2.3, we will verify that it is satisfied for the general class of Bernoulli and the PHD filters discussed in section 1.1.1 and section 1.1.2,

We are now in position to state the main results of this article. The first one is concerned with the exponential stability properties of the semigroup $\Gamma_{p, n}=\left(\Gamma_{p, n}^{1}, \Gamma_{p, n}^{2}\right)$, with $0 \leq p \leq n$ associated with the one step transformations of the flow (1.2). A more precise description and the complete proof of the next theorem is provided in section 3 .

Theorem 1.1 We let $\Phi_{p, n, \nu}^{1}$ and $\Phi_{p, n, m}^{2}$ be the semigroups associated with the one step transformations of the flow of total masses $\Phi_{n, \nu_{n-1}}^{1}:=\Gamma_{n}^{1}\left(., \nu_{n-1}\right)$ and measures $\Phi_{n, m_{n-1}}^{2}:=$ $\Gamma_{n}^{2}\left(m_{n-1},.\right)$, with a fixed collection of measures $\nu:=\left(\nu_{n}\right)_{n \geq 0} \in \prod_{n \geq 0} \mathcal{P}\left(E_{n}\right)$ and masses $m:=\left(m_{n}\right)_{n \geq 0} \in \prod_{n \geq 0} I_{n}$. When these semigroups are exponentially stable (in the sense that they forget exponentially fast their initial conditions) and when the pair of mappings $\nu_{n-1} \mapsto \Phi_{n, \nu_{n-1}}^{1}$ and $m_{n-1} \mapsto \Phi_{n, m_{n-1}}^{2}$ are sufficiently regular then we have the following contraction inequalities

$$
\left|\Gamma_{p, n}^{1}\left(u^{\prime}, \eta^{\prime}\right)-\Gamma_{p, n}^{1}(u, \eta)\right| \vee\left\|\Gamma_{p, n}^{2}\left(u^{\prime}, \eta^{\prime}\right)-\Gamma_{p, n}^{2}(u, \eta)\right\|_{\mathrm{tv}} \leq c e^{-\lambda(n-p)}
$$

for any $p \leq n, u, u^{\prime} \in I_{p}, \eta, \eta^{\prime} \in \mathcal{P}\left(E_{p}\right)$, and some finite constants $c<\infty$ and $\lambda>0$ whose values do not depend on the time parameters $p \leq n$.

The second theorem is concerned with estimating the approximation error associated with a $N$-approximation model satisfying condition (1.6). The first part of the theorem is proved in section 1.2 . The proof of the uniform estimates is discussed in section 3.1 (see for instance lemma 3.4).

Theorem 1.2 Under the assumptions $\left(H_{1}\right)$ and $\left(H_{2}\right)$, the semigroup $\Gamma_{p, n}$ satisfies the same Lipschitz type inequalities as those stated in (1.14) and (1.15) for some collection of measures $\Sigma_{p, n}^{1}$ and $\Sigma_{p, n}^{2}(f,$.$) on \mathcal{B}\left(E_{p}\right)$. In addition, for any $N$-approximation model satisfying condition (1.6) we have the estimates:

$$
\begin{equation*}
\mathbb{E}\left(\left|V_{n}^{\gamma, N}(1)\right|^{r}\right)^{\frac{1}{r}} \leq a_{r} \sum_{p=0}^{n} \delta\left(\Sigma_{p, n}^{1}\right) \quad \text { and } \quad \mathbb{E}\left(\left|V_{n}^{\eta, N}(f)\right|^{r}\right)^{\frac{1}{r}} \leq a_{r} \sum_{p=0}^{n} \delta\left(\Sigma_{p, n}^{2}\right) \tag{1.16}
\end{equation*}
$$

for any $r \geq 1$, and $N \geq 1$, with some constants $a_{r}<\infty$ whose values only depend on $r$. Furthermore, under the regularity conditions of theorem 1.1 the couple of estimates stated above are uniform w.r.t. the time horizon; that is, we have that

$$
\forall i=1,2 \quad \sup _{n \geq 0} \sum_{p=0}^{n} \delta\left(\Sigma_{p, n}^{i}\right)<\infty
$$

These rather abstract theorems apply to a general class of discrete generation measurevalued equations of the form (1.1). We illustrate the application of this pair of theorems in the analysis of the stability properties and the approximation convergence of the pair of multiple target filters presented in this introductory section.

These results can basically be stated as follows:

- The Bernoulli filter presented in section 1.1.1 with a sufficiently mixing prediction and almost equal survival and spontaneous births rates $s_{n} \sim \mu_{n}(1)$ is exponentially stable.
- The PHD filter presented in section 1.1 .2 is exponentially stable for small clutter intensities and sufficiently high detection probability and spontaneous birth rates.
- In both situations, the estimation error of any $N$-approximation model satisfying condition (1.6) does not accumulate over time. Furthermore, the uniform rates of convergence provided in theorem 1.2 allows to design stochastic algorithms with prescribed performance index at any time horizon.

We end this section with some direct consequences of theorem 1.2 ,
Firstly, we observe that the mean error estimates stated in the above theorem clearly implies the almost sure convergence results

$$
\lim _{N \rightarrow \infty} \eta_{n}^{N}(f)=\eta_{n}(f) \quad \text { and } \quad \lim _{N \rightarrow \infty} \gamma_{n}^{N}(f)=\gamma_{n}(f)
$$

for any bounded function $f$ on $E_{n}$. Furthermore, with some information on the constants $a_{r}$, these $\mathbb{L}_{r}$-mean error bounds can be turned to exponential concentration inequalities. To be more precise, by lemma 7.3 .3 in [6], the collection of constants $a_{r}$ in theorem 1.2 , can be chosen so that

$$
\begin{equation*}
a_{2 r}^{2 r} \leq b^{2 r}(2 r)!2^{-r} / r!\quad \text { and } \quad a_{2 r+1}^{2 r+1} \leq b^{2 r+1}(2 r+1)!2^{-r} / r! \tag{1.17}
\end{equation*}
$$

for some $b<\infty$, whose values do not depend on $r$. Arguing as in [5], the above $\mathbb{L}_{r}$-mean error bounds we can establish the following non asymptotic Gaussian tail estimates:

$$
\mathbb{P}\left(\left|\left[\eta_{n}^{N}-\eta_{n}\right](f)\right| \geq \frac{b_{n}}{\sqrt{N}}+\epsilon\right) \leq \exp \left(-\frac{N \epsilon^{2}}{2 b_{n}^{2}}\right) \quad \text { with } \quad b_{n} \leq b \sum_{p=0}^{n} \delta\left(\Sigma_{p, n}^{2}\right)
$$

It is worth noting that the above constructions allows us to consider with further work branching particle models in path spaces. These path space models arise in the analysis of the historical process associated with a branching models as well as the analysis of a filtering problem of the whole signal path given a series of observations. For a more thorough discussion on these path space models, we refer the reader to [6, 7]. The last referenced article also provides a detailed study on the long time behavior of the first moments evolution of nonlinear spatial branching processes, including time uniform estimates for particle interpretation models. In this context, the spatial branching rates are homogeneous w.r.t. the time parameter and the resulting model can be represented by Feynman-Kac particle absorption model of the same form as the one discussed in section 3.2 (see for instance [9]). The limiting normalized measures associated with these particle absorption models are often referred as quasi-invariant measures or Yaglom's distributions. The stability properties of
these absorption models, their particle interpretations, and their connections with Feynman-Kac-Schroedinger semigroups is rather well understood, see for instance [5, 7, 9], as well as section 12.4 in [6], and references therein.

We end this section with some standard notation used in the paper:
We denote respectively by $\mathcal{M}(E), \mathcal{P}(E)$, and $\mathcal{B}(E)$, the set of all finite positive measures $\mu$ on some measurable space $(E, \mathcal{E})$, the convex subset of all probability measures, and the Banach space of all bounded and measurable functions $f$ equipped with the uniform norm $\|f\|$. We denote by $f^{-}$and $f^{+}$the infimum and the supremum of a function $f$. For measurable subsets $A \in \mathcal{E}$, in various instances we slightly abuse notation and we denote $\mu(A)$ instead of $\mu\left(1_{A}\right)$; and we set $\delta_{a}$ the Dirac measure at $a \in E$. We recall that a bounded and positive integral operator $Q$ from a measurable space ( $E_{1}, \mathcal{E}_{1}$ ) into an auxiliary measurable space $\left(E_{2}, \mathcal{E}_{2}\right)$ is an operator $f \mapsto Q(f)$ from $\mathcal{B}\left(E_{2}\right)$ into $\mathcal{B}\left(E_{1}\right)$ such that the functions

$$
x \mapsto Q(f)(x):=\int_{E_{2}} Q(x, d y) f(y)
$$

are $\mathcal{E}_{1}$-measurable and bounded for some measures $Q(x,.) \in \mathcal{M}\left(E_{2}\right)$. These operators also generate a dual operator $\mu \mapsto \mu Q$ from $\mathcal{M}\left(E_{1}\right)$ into $\mathcal{M}\left(E_{2}\right)$ defined by $(\mu Q)(f):=\mu(Q(f))$. A Markov kernel is a positive and bounded integral operator $M$ with $M(1)=1$. We denote by $Q_{p, n}=Q_{p+1} Q_{p+2} \ldots Q_{n}$, with $p \leq n$ the semigroup associated with a given sequence of bounded and positive integral operator $Q_{n}$ from some measurable spaces ( $E_{n-1}, \mathcal{E}_{n-1}$ ) into $\left(E_{n}, \mathcal{E}_{n}\right)$. For $p=n$, we use the convention $Q_{n, n}=I d$, the identity operator.

We associate with a bounded positive potential function $G: x \in E \mapsto G(x) \in[0, \infty)$, the Bayes-Boltzmann-Gibbs transformations

$$
\Psi_{G}: \eta \in \mathcal{M}(E) \mapsto \Psi_{G}(\eta) \in \mathcal{P}(E) \quad \text { with } \quad \Psi_{G}(\eta)(d x):=\frac{1}{\eta(G)} G(x) \eta(d x)
$$

provided $\eta(G)>0$. We recall that $\Psi_{G}(\eta)$ can be expressed in terms of a Markov transport equation

$$
\begin{equation*}
\eta S_{\eta}=\Psi_{G}(\eta) \tag{1.1}
\end{equation*}
$$

for some selection type transition $S_{\eta}(x, d y)$. For instance, we can take

$$
\begin{equation*}
S_{\eta}(x, d y):=\frac{\epsilon}{\eta(G)} \delta_{x}(d y)+\left(1-\frac{\epsilon}{\eta(G)}\right) \Psi_{(G-\epsilon)}(\eta)(d y) \tag{1.19}
\end{equation*}
$$

for any $\epsilon \geq 0$ s.t. $G(x) \geq \epsilon$. Notice that for $\epsilon=0$, we have $S_{\eta}(x, d y)=\Psi_{G}(\eta)(d y)$. We can also choose

$$
\begin{equation*}
S_{\eta}(x, d y):=\epsilon G(x) \delta_{x}(d y)+(1-\epsilon G(x)) \Psi_{G}(\eta)(d y) \tag{1.20}
\end{equation*}
$$

for any $\epsilon \geq 0$ that may depend on the current measure $\eta$, and s.t. $\epsilon G(x) \leq 1$. For instance, we can choose $1 / \epsilon$ to be the $\eta$-essential maximum of the potential function $G$. Finally, in the context of Bernoulli and PHD filtering we set $\bar{\mu}_{n+1}=\mu_{n+1} / \mu_{n+1}(1)$, for any $n \geq 0$, the normalized spontaneous birth measures.

## 2 Semigroup description

### 2.1 The Bernoulli filter semigroup

By construction, we notice that the mass process and the normalized measures are given by the rather simple recursive formulae

$$
\begin{equation*}
\gamma_{n+1}(1)=\frac{\gamma_{n}(1) \eta_{n}\left(g_{n}\right)}{\left(1-\gamma_{n}(1)\right)+\gamma_{n}(1) \eta_{n}\left(g_{n}\right)} \Psi_{g_{n}}\left(\eta_{n}\right)\left(s_{n}\right)+\frac{\left(1-\gamma_{n}(1)\right)}{\left(1-\gamma_{n}(1)\right)+\gamma_{n}(1) \eta_{n}\left(g_{n}\right)} \mu_{n+1}(1) \tag{2.1}
\end{equation*}
$$

and $\eta_{n+1}:=\alpha_{n}\left(\gamma_{n}\right) \Psi_{g_{n} s_{n}}\left(\eta_{n}\right) M_{n+1}+\left(1-\alpha_{n}\left(\gamma_{n}\right)\right) \bar{\mu}_{n+1}$, with $\alpha_{n}: \gamma \in \mathcal{M}\left(E_{n}\right) \mapsto \alpha_{n}(\gamma)$ defined by

$$
\alpha_{n}(\gamma)=\frac{\gamma\left(g_{n} s_{n}\right)}{\gamma\left(s_{n} g_{n}\right)+(1-\gamma(1)) \mu_{n+1}(1)}
$$

By construction, if we set $\gamma=m \times \eta$ then

$$
\Gamma_{n+1}^{1}(m, \eta)=\frac{\gamma\left(g_{n}\right)}{(1-m)+\gamma\left(g_{n}\right)} \Psi_{g_{n}}(\eta)\left(s_{n}\right)+\frac{(1-m)}{(1-m)+\gamma\left(g_{n}\right)} \mu_{n+1}(1)
$$

and $\quad \Gamma_{n+1}^{2}(m, \eta)=\Psi_{g_{n} s_{n}}(\eta) M_{n+1, \gamma}$, with the Markov transitions $M_{n+1, \gamma}$ defined below

$$
\begin{equation*}
M_{n+1, \gamma}(x, .):=\alpha_{n}(\gamma) M_{n+1}(x, .)+\left(1-\alpha_{n}(\gamma)\right) \bar{\mu}_{n+1} \tag{2.2}
\end{equation*}
$$

Next we provide an alternative interpretation of the mapping $\Gamma_{n+1}^{2}$. Firstly, observe that

$$
\begin{equation*}
\Psi_{g_{n} s_{n}}(\eta) M_{n+1, \gamma}(f)=\frac{\eta\left(Q_{n+1, m}(f)\right)}{\eta\left(Q_{n+1, m}(1)\right)} \tag{2.3}
\end{equation*}
$$

with the integral operator

$$
Q_{n+1, m}(f)(x):=m g_{n}(x) s_{n}(x) M_{n+1}(f)(x)+(1-m) \mu_{n+1}(f)
$$

This implies that $\Gamma_{n+1}^{2}(m, \eta)=\Psi_{\widehat{G}_{n, m}}(\eta) \widehat{M}_{n+1, m}$, with the potential function

$$
\begin{equation*}
\widehat{G}_{n, m}=m g_{n} s_{n}+(1-m) \mu_{n+1}(1) \tag{2.4}
\end{equation*}
$$

and the Markov transitions

$$
\begin{equation*}
\widehat{M}_{n+1, m}(f):=\frac{m g_{n} s_{n}}{m g_{n} s_{n}+(1-m) \mu_{n+1}(1)} M_{n+1}(f)+\frac{(1-m) \mu_{n+1}(1)}{m g_{n} s_{n}+(1-m) \mu_{n+1}(1)} \bar{\mu}_{n+1}(f) \tag{2.5}
\end{equation*}
$$

The condition $\left(H_{1}\right)$ is clearly not met for the Bernoulli filter 1.8) when $s_{n}=0$ and $\mu_{n+1}(1)=$ 0 , since in this situation $\gamma_{n}=0$ for any $n \geq 1$. Nevertheless, this condition is met with $I_{n} \subset(0,1]$ and $m \theta_{+, n}(m)=1$, as long as $s_{n}$ and $\mu_{n+1}(1)$ are uniformly bounded from below. It is also met for $s_{n}=0$, as long as $0<\mu_{n+1}(1)<1$ and the likelihood function given in (1.9) is uniformly bounded. The condition is also met for $\mu_{n+1}(1)=0$, as long as $\gamma_{0}(1)>0$, and the likelihood function given in (1.9) and the function $s_{n}$ are uniformly lower bounded.

We prove these assertions using the fact that

$$
\begin{equation*}
\gamma_{n+1}(1)=\widehat{\gamma}_{n}(1) \Psi_{g_{n}}\left(\eta_{n}\right)\left(s_{n}\right)+\left(1-\widehat{\gamma}_{n}(1)\right) \mu_{n+1}(1) \tag{2.6}
\end{equation*}
$$

with the updated mass parameters $\widehat{\gamma}_{n}(1):=\frac{\gamma_{n}(1) \eta_{n}\left(g_{n}\right)}{\left(1-\gamma_{n}(1)\right)+\gamma_{n}(1) \eta_{n}\left(g_{n}\right)} \in[0,1]$. If we set $s_{n}^{-}:=$ $\inf _{E_{n}} s_{n}$ and $s_{n}^{+}=\sup _{E_{n}} s_{n}$ then we have $\gamma_{n}(1) \in\left[m_{n}^{-}, m_{n}^{+}\right]$for any $n \geq 1$, with parameters $m_{n}^{-}=\mu_{n}(1) \wedge s_{n-1}^{-}$and $m_{n}^{+}=\mu_{n}(1) \vee s_{n-1}^{+}(\leq 1)$. If $s_{n}$ and $\mu_{n+1}(1)$ are uniformly bounded from below then we have $m_{n}^{-}>0$. In addition, for the constant mapping $s_{n}=\mu_{n+1}(1)$, the total mass process is constant $\gamma_{n+1}(1)=m_{n+1}^{+}=m_{n+1}^{-}=\mu_{n+1}(1)$. Furthermore, in this situation the flow of normalized measures is given by the updating-prediction transformation defined by $\eta_{n+1}=\Psi_{g_{n}^{(s)}}\left(\eta_{n}\right) M_{n+1}^{(s)}$, with the likelihood function $g_{n}^{(s)}$ and the Markov transitions $M_{n+1}^{(s)}$ defined by

$$
\begin{equation*}
g_{n}^{(s)}:=s_{n} g_{n}+\left(1-s_{n}\right) \quad \text { and } \quad M_{n+1}^{(s)}(f):=\frac{s_{n} g_{n} M_{n+1}(f)+\left(1-s_{n}\right) \bar{\mu}_{n+1}(f)}{s_{n} g_{n}+\left(1-s_{n}\right)} \tag{2.7}
\end{equation*}
$$

When $\mu_{n+1}(1)=0$, the flow of normalized measures is again given by a simple updatingprediction equation

$$
\begin{equation*}
\eta_{n+1}=\Psi_{g_{n} s_{n}}\left(\eta_{n}\right) M_{n+1} \quad \text { and } \quad \gamma_{n+1}(1)=\Psi_{g_{n}}\left(\eta_{n}\right)\left(s_{n}\right) \times \theta_{\eta_{n}\left(g_{n}\right)}\left(\gamma_{n}(1)\right) \tag{2.8}
\end{equation*}
$$

with the increasing mappings $\theta_{a}$ defined below

$$
\begin{equation*}
x \in[0,1] \mapsto \theta_{a}(x):=a x /[a x+(1-x)] \tag{2.9}
\end{equation*}
$$

In addition, if $s_{n}^{-}>0$ then $m_{n+1}^{-} \geq s_{n}^{-} \times \frac{g_{n}^{-} m_{n}^{-}}{g_{n}^{-} m_{n}^{-}+\left(1-m_{n}^{-}\right)}>0$, as long as $g_{n}^{-}:=\inf _{E_{n}} g_{n}>$ 0 , and $\gamma_{0}(1)>0$. We prove this inequality using the fact that the mapping $(a, x) \in$ $\left[0, \infty\left[\times[0,1] \mapsto \theta_{a}(x)\right.\right.$ is increasing in both coordinates. In the case where $s_{n}=1$, using the fact that and $\theta_{a} \circ \theta_{b}=\theta_{a b}$, we prove that

$$
\gamma_{n+1}(1)=\theta_{\eta_{n}\left(g_{n}\right)}\left(\gamma_{n}(1)\right)=\theta_{\prod_{p=0}^{n} \eta_{p}\left(g_{p}\right)}\left(\gamma_{0}(1)\right)
$$

Conversely, when $\gamma_{0}(1)<1$ and $0<\mu_{n+1}(1)<1$ and $s_{n}=0$, for any $n \geq 0$, then we have a constant flow of normalized measures $\eta_{n}=\bar{\mu}_{n}$, with $n \geq 1$, and the total mass process is such that

$$
\left.\gamma_{n}(1) \in\right] 0,1\left[\Longrightarrow \gamma_{n+1}(1)=\mu_{n+1}(1) \times\left[1-\theta_{\bar{\mu}_{n}\left(g_{n}\right)}\left(\gamma_{n}(1)\right)\right] \in\right] 0,1[
$$

with the convention $\bar{\mu}_{0}=\eta_{0}$, for $n=0$. In addition, if $\mu_{n+1}(1)=1$ then we have

$$
\gamma_{2(n+1)}(1)=\theta_{\prod_{p=0}^{n}\left(b_{2 p} / b_{2 p+1}\right)}\left(\gamma_{0}(1)\right) \quad \text { and } \quad \gamma_{2 n+1}(1)=\theta_{b_{2 n}^{-1} \prod_{p=0}^{n-1}\left(b_{2 p+1} / b_{2 p}\right)}\left(\gamma_{0}(1)\right)
$$

for any $n \geq 0$, with the parameters $b_{n}:=\bar{\mu}_{n}\left(g_{n}\right)$. We prove these formuae using the the fact that $1-\theta_{a}(x)=\theta_{1 / a}(1-x)$, and $\theta_{a} \circ \theta_{b}=\theta_{a b}$. This again implies that $m_{n}^{-}>0$ as long as $\gamma_{0}(1)>0$ and the likelihood function are uniformly lower bounded.

### 2.2 The PHD filter semigroup

By construction, if we set $\gamma=m \times \eta$ then we find that

$$
\Gamma_{n+1}^{1}(m, \eta)=\gamma\left(g_{n, \gamma}\right)+\mu_{n+1}(1) \quad \text { and } \quad \Gamma_{n+1}^{2}(m, \eta)=\Psi_{g_{n, \gamma}}(\eta) M_{n+1, \gamma}
$$

In the above display, $M_{n+1, \gamma}$ is the collection of Markov transitions defined below

$$
M_{n+1, \gamma}(x, .):=\alpha_{n}(\gamma) M_{n+1}(x, .)+\left(1-\alpha_{n}(\gamma)\right) \bar{\mu}_{n+1} \quad \text { with } \quad \alpha_{n}(\gamma)=\frac{\gamma\left(g_{n, \gamma}\right)}{\gamma\left(g_{n, \gamma}\right)+\mu_{n+1}(1)}
$$

The interpretation of the updating transformation $\Psi_{g_{n, \gamma}}(\eta)$ in terms of a Markov transport equation is non unique. For instance, using (1.12) this Bolzmann-Gibbs transformation can be decomposed into two parts. The first one relates to the undetectable targets and the second is associated with non clutter observations. An alternative description is provided below. We consider a virtual auxiliary observation point $c$ (corresponding to undetectable targets) and set $\mathcal{Y}_{n}^{c}=\mathcal{Y}_{n}+\delta_{c}$. We also denote by $g_{n, \gamma}^{c}(., y)$ the function defined below

$$
g_{n}^{\gamma}(., y)= \begin{cases}r_{n}\left(1-d_{n}\right) & \text { if } y=c \\ r_{n} \frac{d_{n} g_{n}\left(., y_{n}\right)}{h_{n}(y)+\gamma\left(d_{n} g_{n}(., y)\right)} & \text { if } y \neq c\end{cases}
$$

In this notation, the updating transformation $\Psi_{g_{n, \gamma}}(\eta)$ can be rewritten in the following form

$$
\Psi_{g_{n, \gamma}}(\eta)=\Psi_{\bar{g}_{n, \gamma}}(\eta) \quad \text { with } \quad \bar{g}_{n, \gamma}=\int \mathcal{Y}_{n}^{c}(d y) g_{n}^{\gamma}(., y)
$$

The averaged potential function $\bar{g}_{n, \gamma}$ allows us to measure the likelihood of signal states w.r.t. the current observation measure $\mathcal{Y}_{n}^{c}$. Using (1.18), the Bolzmann-Gibbs transformation $\Psi_{\bar{g}_{n, \gamma}}(\eta)$ can be interpreted as non linear Markov transport equation of the following form

$$
\begin{equation*}
\Psi_{\bar{g}_{n, \gamma}}(\eta)=\eta S_{n, \gamma} \quad \text { and } \quad \Gamma^{2}(m, \eta)=\eta K_{n+1, \gamma} \quad \text { with } \quad K_{n+1, \gamma}=S_{n, \gamma} M_{n+1, \gamma} \tag{2.10}
\end{equation*}
$$

for some Markov transitions $S_{n, \gamma}$ from $E_{n}$ into itself.
We also notice that condition $\left(H_{1}\right)$ holds as long as the functions $s_{n}, b_{n}$, and $g_{n}\left(., y_{n}\right)$ are uniformly bounded and $\mu_{n}(1)>0$. It is also met when $\mu_{n}(1)=0$, as long as $r_{n}=\left(s_{n}+b_{n}\right)$ is uniformly lower bounded and $\mathcal{Y}_{n} \neq 0$ or $d_{n}<1$.

### 2.3 Lipschitz regularity properties

Firstly, we mention that condition $\left(H_{2}\right)$ can be replaced by the following regularity condition:
$\left(H_{2}^{\prime}\right):$ For any $n \geq 1, f \in \operatorname{Osc}_{1}\left(E_{n}\right)$, and any $(m, \eta),\left(m^{\prime}, \eta^{\prime}\right) \in\left(I_{n} \times \mathcal{P}\left(E_{n}\right)\right)$, the integral operators $Q_{n, m \eta}$ satisfy the following Lipschitz type inequalities:

$$
\begin{equation*}
\left\|Q_{n, m \eta}(f)-Q_{n, m^{\prime} \eta^{\prime}}(f)\right\| \leq c(n)\left|m-m^{\prime}\right|+\int\left|\left[\eta-\eta^{\prime}\right](\varphi)\right| \Sigma_{n,\left(m^{\prime}, \eta^{\prime}\right)}(f, d \varphi) \tag{2.11}
\end{equation*}
$$

for some collection of bounded measures $\Sigma_{n,\left(m^{\prime}, \eta^{\prime}\right)}(f,$.$) on \mathcal{B}\left(E_{n}\right)$ such that

$$
\int \operatorname{osc}(\varphi) \Sigma_{n,(m, \eta)}(f, d \varphi) \leq \delta\left(\Sigma_{n}\right)
$$

for some finite constant $\delta\left(\Sigma_{n}\right)<\infty$, whose values do dot depend on the parameters $(m, \eta) \in$ $\left(I_{n} \times \mathcal{P}\left(E_{n}\right)\right)$ and $f \in \operatorname{Osc}_{1}\left(E_{n}\right)$.

We prove $\left(H_{2}^{\prime}\right) \Rightarrow 1.14$ using the decompositions

$$
m \eta Q_{n, m \eta}-m^{\prime} \eta^{\prime} Q_{n, m^{\prime} \eta^{\prime}}=m \eta\left[Q_{n, m \eta}-Q_{n, m^{\prime} \eta^{\prime}}\right]+\left[m \eta-m^{\prime} \eta^{\prime}\right] Q_{n, m^{\prime} \eta^{\prime}}
$$

and of course $\left[m \eta-m^{\prime} \eta^{\prime}\right]=\left[m-m^{\prime}\right] \eta+m^{\prime}\left[\eta-\eta^{\prime}\right]$. To prove $\left.\left(H_{2}^{\prime}\right) \Rightarrow 1.15\right)$, we let $\gamma=m \eta$ and $\gamma^{\prime}=m^{\prime} \eta^{\prime}$ and we use the decomposition

$$
\left[\Gamma_{n}^{2}(m, \eta)-\Gamma_{n}^{2}\left(m^{\prime}, \eta^{\prime}\right)\right](f)=\frac{1}{\gamma Q_{n, \gamma}(1)}\left[\gamma Q_{n, \gamma}-\gamma^{\prime} Q_{n, \gamma^{\prime}}\right]\left(f-\Gamma_{n}^{2}\left(m^{\prime}, \eta^{\prime}\right)(f)\right)
$$

The Bernoulli filter (1.8) satisfies $\left(H_{2}^{\prime}\right)$, as long as the likelihood functions $g_{n}$ given in (1.9) are uniformly bounded above. In this situation, (2.11) is met with

$$
\left\|Q_{n, m \eta}(f)-Q_{n, m^{\prime} \eta^{\prime}}(f)\right\| \leq c(n)\left|m-m^{\prime}\right|+c^{\prime}(n)\left|\left[\eta-\eta^{\prime}\right]\left(g_{n}\right)\right|
$$

for some finite constant $c^{\prime}(n)<\infty$.
The PHD equation satisfies $\left(H_{2}^{\prime}\right)$, as long as the functions $h_{n}(y)+g_{n, y}^{\prime}$ with $g_{n, y}^{\prime}:=$ $d_{n} g_{n}(., y)$ are uniformly bounded above and below. To prove this claim, we simply use the fact that

$$
\left\|\widehat{g}_{n, \gamma}-\widehat{g}_{n, \gamma^{\prime}}\right\| \leq c_{n}\left[\left|m^{\prime}-m\right|+\int \mathcal{Y}_{n}(d y)\left|\left[\eta^{\prime}-\eta\right]\left(g_{n, y}^{\prime}\right)\right|\right]
$$

This estimate is a direct consequence of the following one

$$
\widehat{g}_{n, \gamma}(x)-\widehat{g}_{n, \gamma^{\prime}}(x)=\int \mathcal{Y}_{n}(d y) \frac{g_{n, y}^{\prime}(x)}{h_{n}(y)+\gamma\left(g_{n, y}^{\prime}\right)} \frac{\left[\gamma^{\prime}-\gamma\right]\left(g_{n, y}^{\prime}\right)}{h_{n}(y)+\gamma^{\prime}\left(g_{n, y}^{\prime}\right)}
$$

Next, we provide a pivotal regularity property of the semigroup $\left(\Gamma_{p, n}\right)_{0 \leq p \leq n}$ associated with the one step transformations of the flow (1.2).

Proposition 2.1 We assume that conditions $\left(H_{1}\right)$ and $\left(H_{2}\right)$ are satisfied. Then, for any $0 \leq p \leq n, f \in \operatorname{Osc}_{1}\left(E_{n}\right)$, and any $(m, \eta),\left(m^{\prime}, \eta^{\prime}\right) \in\left(I_{p} \times \mathcal{P}\left(E_{p}\right)\right)$, we have the following Lipschitz type inequalities:

$$
\begin{aligned}
\left|\Gamma_{p, n}^{1}(m, \eta)-\Gamma_{p, n}^{1}\left(m^{\prime}, \eta^{\prime}\right)\right| & \leq c_{p}(n)\left|m-m^{\prime}\right|+\int\left|\left[\eta-\eta^{\prime}\right](\varphi)\right| \Sigma_{p, n,\left(m^{\prime}, \eta^{\prime}\right)}^{1}(d \varphi) \\
\left|\left[\Gamma_{p, n}^{2}(m, \eta)-\Gamma_{p, n}^{2}\left(m^{\prime}, \eta^{\prime}\right)\right](f)\right| & \leq c_{p}(n)\left|m-m^{\prime}\right|+\int\left|\left[\eta-\eta^{\prime}\right](\varphi)\right| \Sigma_{p, n,\left(m^{\prime}, \eta^{\prime}\right)}^{2}(f, d \varphi)
\end{aligned}
$$

for some finite constants $c_{p}(n)<\infty$, and some collection of bounded measures $\Sigma_{p, n,\left(m^{\prime}, \eta^{\prime}\right)}^{1}$ and $\Sigma_{p, n,\left(m^{\prime}, \eta^{\prime}\right)}^{2}(f$, . $)$ on $\mathcal{B}\left(E_{p}\right)$ such that

$$
\begin{equation*}
\int \operatorname{osc}(\varphi) \Sigma_{p, n,(m, \eta)}^{1}(d \varphi) \leq \delta\left(\Sigma_{p, n}^{1}\right) \quad \text { and } \quad \int \operatorname{osc}(\varphi) \Sigma_{p, n,(m, \eta)}^{2}(f, d \varphi) \leq \delta\left(\Sigma_{p, n}^{2}\right) \tag{2.12}
\end{equation*}
$$

for some finite constant $\delta\left(\Sigma_{p, n}^{i}\right)<\infty, i=1,2$, whose values do dot depend on the parameters $(m, \eta) \in\left(I_{p} \times \mathcal{P}\left(E_{p}\right)\right.$ and $f \in \operatorname{Osc}_{1}\left(E_{n}\right)$.

## Proof:

To prove this proposition, we use a backward induction on the parameter $1 \leq p \leq n$. For $p=(n-1)$, we have $\Gamma_{n-1, n}^{i}=\Gamma_{n}^{i}$, with $i=1,2$, so that the desired result is satisfied for $p=(n-1)$. We further assume that the estimates hold at a given rank $p<n$. To prove the estimates at rank $(p-1)$, we recall that

$$
\Gamma_{p-1, n}(m, \eta)=\Gamma_{p, n}\left(\Gamma_{p}(m, \eta)\right) \Rightarrow \forall i=1,2 \quad \Gamma_{p-1, n}^{i}(m, \eta)=\Gamma_{p, n}^{i}\left(\Gamma_{p}(m, \eta)\right)
$$

Under the induction hypothesis

$$
\begin{aligned}
\left|\Gamma_{p-1, n}^{1}(m, \eta)-\Gamma_{p-1, n}^{1}\left(m^{\prime}, \eta^{\prime}\right)\right| \leq & c_{p}(n)\left|\Gamma_{p}^{1}(m, \eta)-\Gamma_{p}^{1}\left(m^{\prime}, \eta^{\prime}\right)\right| \\
& +\int\left|\left[\Gamma_{p}^{2}(m, \eta)-\Gamma_{p}^{2}\left(m^{\prime}, \eta^{\prime}\right)\right](\varphi)\right| \Sigma_{p, n, \Gamma_{p}\left(m^{\prime}, \eta^{\prime}\right)}^{1}(d \varphi)
\end{aligned}
$$

Using the fact that

$$
\begin{aligned}
\left|\Gamma_{p}^{1}(m, \eta)-\Gamma_{p}^{1}\left(m^{\prime}, \eta^{\prime}\right)\right| & \leq c(p)\left|m-m^{\prime}\right|+\int\left|\left[\eta-\eta^{\prime}\right](\varphi)\right| \Sigma_{p,\left(m^{\prime}, \eta^{\prime}\right)}^{1}(d \varphi) \\
\left|\left[\Gamma_{p}^{2}(m, \eta)-\Gamma_{p}^{2}\left(m^{\prime}, \eta^{\prime}\right)\right](\varphi)\right| & \leq c(p)\left|m-m^{\prime}\right|+\int\left|\left[\eta-\eta^{\prime}\right](\psi)\right| \Sigma_{p,\left(m^{\prime}, \eta^{\prime}\right)}^{2}(\varphi, d \psi)
\end{aligned}
$$

the end of the proof is now clear. The analysis of $\Gamma_{p-1, n}^{2}$ follows the same line of arguments and is omitted. This ends the proof of the proposition.

### 2.4 Proof of theorem 1.2

This section is mainly concerned with the proof of the couple of estimates 1.16 stated in theorem 1.2 . We use the decomposition

$$
\begin{align*}
\left(\gamma_{n}^{N}(1), \eta_{n}^{N}\right)-\left(\gamma_{n}(1), \eta_{n}\right)= & {\left[\Gamma_{0, n}\left(\gamma_{0}^{N}(1), \eta_{0}^{N}\right)-\Gamma_{0, n}\left(\gamma_{0}(1), \eta_{0}\right)\right] } \\
& +\sum_{p=1}^{n}\left[\Gamma_{p, n}\left(\gamma_{p}^{N}(1), \eta_{p}^{N}\right)-\Gamma_{p-1, n}\left(\gamma_{p-1}^{N}(1), \eta_{p-1}^{N}\right)\right] \tag{2.13}
\end{align*}
$$

and the fact that

$$
\Gamma_{p-1, p}\left(\gamma_{p-1}^{N}(1), \eta_{p-1}^{N}\right)=\left(\gamma_{p}^{N}(1), \Gamma_{p-1, p}^{2}\left(\gamma_{p-1}^{N}(1), \eta_{p-1}^{N}\right)\right)
$$

to show that

$$
\begin{aligned}
\gamma_{n}^{N}(1)-\gamma_{n}(1)= & {\left[\Gamma_{0, n}^{1}\left(\gamma_{0}^{N}(1), \eta_{0}^{N}\right)-\Gamma_{0, n}^{1}\left(\gamma_{0}(1), \eta_{0}\right)\right] } \\
& +\sum_{p=1}^{n}\left[\Gamma_{p, n}^{1}\left(\gamma_{p}^{N}(1), \eta_{p}^{N}\right)-\Gamma_{p, n}^{1}\left(\gamma_{p}^{N}(1), \Gamma_{p-1, p}^{2}\left(\gamma_{p-1}^{N}(1), \eta_{p-1}^{N}\right)\right)\right]
\end{aligned}
$$

Recalling that $\gamma_{0}^{N}(1)=\gamma_{0}(1)$, using proposition 2.1, we find that

$$
\sqrt{N}\left|\gamma_{n}^{N}(1)-\gamma_{n}(1)\right| \leq \sum_{p=0}^{n} c_{p}(n) \int \mid\left[W_{p}^{N}(\varphi) \mid \Sigma_{p, n}^{(N, 1)}(d \varphi)\right.
$$

with the predictable measure $\Sigma_{p, n}^{(N, 1)}=\Sigma_{p, n,(m, \eta)}^{1}$ associated with the parameters $(m, \eta)=$ $\left(\gamma_{p}^{N}(1), \Gamma_{p-1, p}^{2}\left(\gamma_{p-1}^{N}(1), \eta_{p-1}^{N}\right)\right)$, with $0<p \leq n$; and for $p=0$, we set $\Sigma_{0, n}^{(N, 1)}=\Sigma_{0, n,\left(\gamma_{0}(1), \eta_{0}\right)}$. Combing the generalized Minkowski's inequality with (1.6) we have

$$
\mathbb{E}\left(\left|\int\right| W_{p}^{N}(\varphi)\left|\Sigma_{p, n}^{(N, 1)}(d \varphi)\right|^{r} \mid \mathcal{F}_{p-1}^{(N)}\right)^{\frac{1}{r}} \leq a_{r} \delta\left(\Sigma_{p, n}^{1}\right)
$$

for some constants $a_{r}$ whose values only depend on the time parameter. This implies that

$$
\mathbb{E}\left(\left|\gamma_{n}^{N}(1)-\gamma_{n}(1)\right|^{r}\right)^{\frac{1}{r}} \leq a_{r} \sum_{p=0}^{n} \delta\left(\Sigma_{p, n}^{1}\right)
$$

The normalized occupation measures can be analyzed in the same way. This ends the proof of the theorem 1.2 .

## 3 Functional contraction inequalities

### 3.1 Stability properties

This section is concerned with the long time behavior of nonlinear measure-valued processes of the form (1.2). The complexity of these models depend in part on the interaction function between the flow of masses $\gamma_{n}(1)$ and the flow of probability measures $\eta_{n}=\gamma_{n} / \gamma_{n}(1)$. One natural way to start the analysis of these models is to study the stability properties of the measure-valued semigroup associated with a fixed flow of masses, and vice versa. These two mathematical objects are defined below.

Definition 3.1 We associate with a flow of masses $m=\left(m_{n}\right)_{n \geq 0} \in \prod_{n \geq 0} I_{n}$ and probability measures $\nu:=\left(\nu_{n}\right)_{n \geq 0} \in \prod_{n \geq 0} \mathcal{P}\left(E_{n}\right)$ the pair of semigroups

$$
\begin{equation*}
\Phi_{p, n, \nu}^{1}:=\Phi_{n, \nu_{n-1}}^{1} \circ \ldots \circ \Phi_{1, \nu_{0}}^{1} \quad \text { and } \quad \Phi_{p, n, m}^{2}:=\Phi_{n, m_{n-1}}^{2} \circ \ldots \circ \Phi_{1, m_{0}}^{2} \tag{3.1}
\end{equation*}
$$

with $0 \leq p \leq n$, and the one step transformations

$$
\begin{aligned}
\Phi_{n, \nu_{n-1}}^{1} & : u \in I_{n-1} \mapsto \Phi_{n, \nu_{n-1}}^{1}(u):=\Gamma_{n}^{1}\left(u, \nu_{n-1}\right) \in I_{n} \\
\Phi_{n, m_{n-1}}^{2} & : \eta \in \mathcal{P}\left(E_{n-1}\right) \mapsto \Phi_{n, m_{n-1}}^{2}(\eta):=\Gamma_{n}^{2}\left(m_{n-1}, \eta\right) \in \mathcal{P}\left(E_{n}\right)
\end{aligned}
$$

By construction, using a simple induction on the time parameter $n$, we find that

$$
\begin{gathered}
\left(m_{0}, \nu_{0}\right)=\left(\gamma_{0}(1), \eta_{0}\right) \text { and } \forall n \geq 1 \quad m_{n}=\Phi_{n, \nu_{n-1}}^{1}\left(m_{n-1}\right) \text { and } \nu_{n}=\Phi_{n, m_{n-1}}^{2}\left(\nu_{n-1}\right) \\
\mathfrak{\Downarrow} \\
\forall n \geq 0 \quad\left(m_{n}, \nu_{n}\right)=\left(\gamma_{n}(1), \eta_{n}\right)
\end{gathered}
$$

In the cases that are of particular interest, the semigroups $\Phi_{p, n, \nu}^{1}$ and $\Phi_{p, n, m}^{2}$ will have a Feynman-Kac representation. These models are rather well understood. A brief review on their contraction properties is provided in section 3.2. Further details can be found in the monograph [6. The first basic regularity property of these models which are needed is the following weak Lipschitz type property :
$(\operatorname{Lip}(\Phi))$ For any $p \leq n, u, u^{\prime} \in I_{p}, \eta, \eta^{\prime} \in \mathcal{P}\left(E_{p}\right)$ and $f \in \operatorname{Osc}_{1}\left(E_{n}\right)$ the following Lipschitz inequalities

$$
\begin{align*}
\left|\Phi_{p, n, \nu}^{1}(u)-\Phi_{p, n, \nu}^{1}\left(u^{\prime}\right)\right| & \leq a_{p, n}^{1}\left|u-u^{\prime}\right|  \tag{3.2}\\
\left|\left[\Phi_{p, n, m}^{2}(\eta)-\Phi_{p, n, m}^{2}\left(\eta^{\prime}\right)\right](f)\right| & \leq a_{p, n}^{2} \int\left|\left[\eta-\eta^{\prime}\right](\varphi)\right| \Omega_{p, n, \eta^{\prime}}^{2}(f, d \varphi) \tag{3.3}
\end{align*}
$$

for some finite constants $a_{p, n}^{i}<\infty$, with $i=1,2$, and some collection of Markov transitions $\Omega_{p, n, \eta^{\prime}}^{2}$ from $\operatorname{Osc}_{1}\left(E_{n}\right)$ into $O s c_{1}\left(E_{p}\right)$, with $p \leq n$, whose values only depend on the parameters $p, n$, resp. $p, n$ and $\eta^{\prime}$.

The semigroups $\Phi_{p, n, \nu}^{1}$ and $\Phi_{p, n, m}^{2}$ may or may bot be asymptotically stable depending on whether $a_{p, n}^{i}$ tends to 0 , as $(n-p) \rightarrow \infty$. In section 3.3 we provide a set of easily checked regularity conditions under which the semigroups associated with the Bernoulli models discussed in 2.1 are asymptotically stable.

The second step in the study of the stability properties of the semigroups associated with the flow (1.2) is the following continuity property:
$(\operatorname{Cont}(\Phi))$ For any $n \geq 1, u, u^{\prime} \in I_{n-1}, \eta, \eta^{\prime} \in \mathcal{P}\left(E_{n-1}\right)$ and any $f \in \operatorname{Osc}_{1}\left(E_{n}\right)$

$$
\begin{align*}
\left|\Phi_{n, \eta}^{1}(u)-\Phi_{n, \eta^{\prime}}^{1}(u)\right| & \leq \tau_{n}^{1} \int\left|\left[\eta-\eta^{\prime}\right](\varphi)\right| \Omega_{n, \eta^{\prime}}^{1}(d \varphi)  \tag{3.4}\\
\left|\left[\Phi_{n, u}^{2}(\eta)-\Phi_{n, u^{\prime}}^{2}(\eta)\right](f)\right| & \leq \tau_{n}^{2}\left|u-u^{\prime}\right| \tag{3.5}
\end{align*}
$$

for some finite constants $\tau_{n}^{i}<\infty$, with $i=1,2$, and some collection probability measures $\Omega_{n, \nu^{\prime}}^{1}$ on $\operatorname{Osc}_{1}\left(E_{n-1}\right)$, whose values only depend on the parameters $n$, resp. $n$ and $\nu^{\prime}$.

This elementary continuity condition allows us to enter the contraction properties of the semigroups $\Phi_{p, n, \nu}^{1}$ and $\Phi_{p, n, m}^{2}$ in the stability analysis of the flow of measures 1.2. The resulting functional contraction inequalities will be described in terms of the following collection of parameters.

Definition 3.2 When the couple of conditions $(\operatorname{Lip}(\Phi))$ and $(\operatorname{Cont}(\Phi))$ stated above are satisfied, for any $i=1,2$ and $p \leq n$ we set

$$
\begin{equation*}
\bar{a}_{p, n}^{i}=\tau_{p+1}^{i} a_{p+1, n}^{i} \quad b_{p, n}=\sum_{p<q<n} \bar{a}_{p, q}^{1} \bar{a}_{q, n}^{2} \quad \text { and } \quad b_{p, n}^{\prime}=\sum_{p \leq q<n} a_{p, q}^{1} \bar{a}_{q, n}^{2} \tag{3.6}
\end{equation*}
$$

The main result of this section is the following proposition.
Proposition 3.3 If conditions $(\operatorname{Lip}(\Phi))$ and $(\operatorname{Cont}(\Phi))$ are satisfied, then for any $p \leq n$, $u, u^{\prime} \in I_{p}, \eta, \eta^{\prime} \in \mathcal{P}\left(E_{p}\right)$ and $f \in O s c_{1}\left(E_{n}\right)$ we have the following Lipschitz inequalities

$$
\begin{aligned}
\left|\Gamma_{p, n}^{1}\left(u^{\prime}, \eta^{\prime}\right)-\Gamma_{p, n}^{1}(u, \eta)\right| & \leq c_{p, n}^{1,1}\left|u-u^{\prime}\right|+c_{p, n}^{1,2} \int\left|\left[\eta-\eta^{\prime}\right](\varphi)\right| \Sigma_{p, n, u^{\prime}, \eta^{\prime}}^{1}(d \varphi) \\
\left.\mid \Gamma_{p, n}^{2}\left(u^{\prime}, \eta^{\prime}\right)(f)-\Gamma_{p, n}^{2}(u, \eta)(f)\right] & \leq c_{p, n}^{2,1}\left|u-u^{\prime}\right|+c_{p, n}^{2,2} \int\left|\left[\eta-\eta^{\prime}\right](\varphi)\right| \Sigma_{p, n, u^{\prime} \eta^{\prime}}^{2}(f, d \varphi)
\end{aligned}
$$

for some probability measures $\Sigma_{p, n, u^{\prime}, \eta^{\prime}}^{1}(d \varphi)$ and Markov transitions $\Sigma_{p, n, m^{\prime} \eta^{\prime}}^{2}$, with the collection of parameters

$$
\begin{aligned}
& c_{p, n}^{1,1}=a_{p, n}^{1}+\sum_{p \leq q<n} c_{p, q}^{2,1} \bar{a}_{q, n}^{1} \quad \text { and } \quad c_{p, n}^{1,2}=\sum_{p \leq q<n} c_{p, q}^{2,2} \bar{a}_{q, n}^{1} \\
& c_{p, n}^{2,1}=b_{p, n}^{\prime}+\sum_{l=1}^{n-p} \sum_{p \leq r_{1}<\ldots r_{l}<n} b_{p, r_{1}}^{\prime} \prod_{1 \leq k \leq l} b_{r_{k}, r_{k+1}} \\
& c_{p, n}^{2,2}=a_{p, n}^{2}+\sum_{l=1}^{n-p} \sum_{p \leq r_{1}<\ldots r_{l}<n} a_{p, r_{1}}^{2} \prod_{1 \leq k \leq l} b_{r_{k}, r_{k+1}}, \quad \text { with the convention } r_{l+1}=n .
\end{aligned}
$$

In particular, the collection of parameters $\delta\left(\Sigma_{p, n}^{i}\right)_{i=1,2}, p \leq n$ introduced in 1.16) and 2.12) are such that $\delta\left(\Sigma_{p, n}^{1}\right) \leq c_{p, n}^{1,2}$ and $\delta\left(\Sigma_{p, n}^{2}\right) \leq c_{p, n}^{2,2}$.

The proof of this proposition is rather technical and it is postponed to section 5.3 in the appendix. Now we conclude this section with a direct application of the above estimates. The proof of the theorem 1.1 stated in the introduction and the uniform estimates discussed in theorem 1.2 are a direct consequence of the following lemma.

Lemma 3.4 Suppose that $\tau^{i}=\sup _{n \geq 1} \tau_{n}^{i}<\infty$, and $a_{p, n}^{i} \leq c_{i} e^{-\lambda_{i}(n-p)}$, for any $p \leq n$, and some finite parameters $c_{i}<\infty$ and $\lambda_{i}>0$, with $i=1,2$, satisfying the following condition

$$
\lambda_{1} \neq \lambda_{2} \quad \text { and } \quad c_{1} c_{2} \tau^{1} \tau^{2} \leq\left(1-e^{-\left(\lambda_{1} \wedge \lambda_{2}\right)}\right) \quad\left(e^{-\left(\lambda_{1} \wedge \lambda_{2}\right)}-e^{-\left(\lambda_{1} \vee \lambda_{2}\right)}\right)
$$

Then, for any $i, j \in\{1,2\}$ we have

$$
c_{p, n}^{i, j} \leq c^{i, j} e^{-\lambda(n-p)} \quad \text { with } \quad \lambda=\left(\lambda_{1} \wedge \lambda_{2}\right)-\log \left(1+c \tau^{1} \tau^{2} \frac{e^{\left(\lambda_{1} \wedge \lambda_{2}\right)}}{e^{-\left(\lambda_{1} \wedge \lambda_{2}\right)}-e^{-\left(\lambda_{1} \vee \lambda_{2}\right)}}\right)>0
$$

and the parameters $c^{i, j}$ defined below

$$
\begin{aligned}
& c^{2,2}=c_{2} \\
& c^{1,1}=c_{1}\left(1+c^{2,1} \tau^{1} /\left(e^{-\lambda}-e^{-\lambda_{1}}\right)\right)
\end{aligned} \begin{aligned}
& c^{2,1}=c_{1} c_{2} \tau^{2} /\left(e^{-\left(\lambda_{1} \wedge \lambda_{2}\right)}-e^{-\left(\lambda_{1} \vee \lambda_{2}\right)}\right) \\
& c^{1,2}=c_{1} c_{2} \tau^{1} /\left(e^{-\lambda}-e^{-\lambda_{1}}\right)
\end{aligned}
$$

In particular, for any $N$-approximation models $\left(\gamma_{n}^{N}(1), \eta_{n}^{N}\right)$ of the flow $\left(\gamma_{n}(1), \eta_{n}\right)$ satisfying condition (1.6), the $\mathbb{L}_{r}$-mean error estimates presented in (1.16) are uniform w.r.t. the time parameter

$$
\sup _{n \geq 0} \mathbb{E}\left(\left|V_{n}^{\gamma, N}(1)\right|^{r}\right)^{\frac{1}{r}} \leq a_{r} c^{1,2} /\left(1-e^{-\lambda}\right) \quad \text { and } \quad \sup _{n \geq 0} \mathbb{E}\left(\left|V_{n}^{\eta, N}(f)\right|^{r}\right)^{\frac{1}{r}} \leq a_{r} c^{2,2} /\left(1-e^{-\lambda}\right)
$$

with some constants $a_{r}<\infty$ whose values only depend on $r$.

## Proof:

Under the premise of the lemma

$$
b_{p, n} \leq c \tau \sum_{p<q<n} e^{-\lambda_{1}(q-(p+1))} e^{-\lambda_{2}(n-(q+1))} \quad \text { and } \quad b_{p, n}^{\prime} \leq c \tau^{2} \sum_{p \leq q<n} e^{-\lambda_{1}(q-p)} e^{-\lambda_{2}(n-(q+1))}
$$

with $c=c_{1} c_{2}$ and $\tau=\tau^{1} \tau^{2}$. We further assume that $\lambda_{1}>\lambda_{2}$ and we set $\Delta=\left|\lambda_{1}-\lambda_{2}\right|$.

$$
b_{p, n} \leq c \tau e^{-\lambda_{2}((n-1)-(p+1))} \sum_{p<q<n} e^{-\Delta(q-(p+1))} \leq c \tau e^{-\lambda_{2}((n-1)-(p+1))} /\left(1-e^{-\Delta}\right)
$$

In the same way, if $\lambda_{2}>\lambda_{1}$ we have

$$
b_{p, n} \leq c \tau e^{-\lambda_{1}((n-1)-(p+1))} \sum_{p<q<n} e^{-\Delta(n-(q+1))} \leq c \tau e^{-\lambda_{1}((n-1)-(p+1))} /\left(1-e^{-\Delta}\right)
$$

This implies that $b_{p, n} \leq c \tau e^{-\left(\lambda_{1} \wedge \lambda_{2}\right)((n-1)-(p+1))} /\left(1-e^{-\Delta}\right)$. In much the same way, it can be shown that

$$
\begin{equation*}
b_{p, n}^{\prime}=\leq c \tau^{2} e^{-\left(\lambda_{1} \wedge \lambda_{2}\right)((n-1)-p)} /\left(1-e^{-\Delta}\right) \tag{3.7}
\end{equation*}
$$

We are now in a position to estimate the parameters $c_{p, n}^{i, j}$. Firstly, we observe that

$$
c_{p, n}^{2,2} \leq c_{2} e^{-\lambda_{2}(n-p)}+c_{2} \sum_{l=1}^{n-p}\left(\frac{c \tau^{1} \tau^{2} e^{2\left(\lambda_{1} \wedge \lambda_{2}\right)}}{1-e^{-\Delta}}\right)^{l} \sum_{p \leq r_{1}<\ldots r_{l}<n} e^{-\lambda_{2}\left(r_{1}-p\right)} e^{-\left(\lambda_{1} \wedge \lambda_{2}\right)\left(n-r_{1}\right)}
$$

When $\lambda_{1}>\lambda_{2}$, we find that

$$
c_{p, n}^{2,2} \leq c_{2} e^{-\lambda_{2}(n-p)} \sum_{l=0}^{n-p}\left(\frac{c \tau e^{2 \lambda_{2}}}{1-e^{-\Delta}}\right)^{l}\binom{n-p}{l}
$$

and therefore

$$
c_{p, n}^{2,2} \leq c_{2} e^{-\lambda_{2}(n-p)}\left(1+c \tau \frac{e^{2 \lambda_{2}}}{1-e^{-\Delta}}\right)^{n-p} \Rightarrow c_{p, n}^{2,2}=c_{2} e^{-\lambda(n-p)}
$$

with

$$
\lambda=\lambda_{2}-\log \left(1+c \tau \frac{e^{\lambda_{2}}}{e^{-\lambda_{2}}-e^{-\lambda_{1}}}\right)>0 \quad \text { as long as } \quad c \tau \leq\left(1-e^{-\lambda_{2}}\right) \quad\left(e^{-\lambda_{2}}-e^{-\lambda_{1}}\right) .
$$

When $\lambda_{2}>\lambda_{1}$ we have $\lambda_{2}=\lambda_{1}+\Delta$, we find that

$$
c_{p, n}^{2,2} \leq c_{2} e^{-\lambda_{2}(n-p)}+c_{2} e^{-\lambda_{1}(n-p)} \sum_{l=1}^{n-p}\left(\frac{c \tau e^{2 \lambda_{1}}}{1-e^{-\Delta}}\right)^{l} \sum_{p \leq r_{1}<\ldots r_{l}<n} e^{-\Delta\left(r_{1}-p\right)}
$$

from which it follows that

$$
c_{p, n}^{2,2} \leq c_{2} e^{-\lambda_{1}(n-p)}\left(1+c \tau \frac{e^{2 \lambda_{1}}}{1-e^{-\Delta}}\right)^{n-p}
$$

Using a similar line of argument as above, we have $c_{p, n}^{2,2} \leq c_{2} e^{-\lambda(n-p)}$, with

$$
\lambda=\lambda_{1}-\log \left(1+c \tau \frac{e^{\lambda_{1}}}{e^{-\lambda_{1}}-e^{-\lambda_{2}}}\right)>0 \quad \text { as long as } \quad c \tau \leq\left(1-e^{-\lambda_{1}}\right)\left(e^{-\lambda_{1}}-e^{-\lambda_{2}}\right)
$$

We conclude that $c_{p, n}^{2,2} \leq c_{2} e^{-\lambda(n-p)}$, with

$$
\lambda=\left(\lambda_{1} \wedge \lambda_{2}\right)-\log \left(1+c \tau \frac{e^{\left(\lambda_{1} \wedge \lambda_{2}\right)}}{e^{-\left(\lambda_{1} \wedge \lambda_{2}\right)}-e^{-\left(\lambda_{1} \vee \lambda_{2}\right)}}\right)>0
$$

as long as $c \tau \leq\left(1-e^{-\left(\lambda_{1} \wedge \lambda_{2}\right)}\right) \quad\left(e^{-\left(\lambda_{1} \wedge \lambda_{2}\right)}-e^{-\left(\lambda_{1} \vee \lambda_{2}\right)}\right)$. Using 3.7) we also show that

$$
c_{p, n}^{2,1} \leq c^{2,1} e^{-\lambda(n-p)} \quad \text { with } \quad c^{2,1}=c \tau^{2} \frac{1}{e^{-\left(\lambda_{1} \wedge \lambda_{2}\right)}-e^{-\left(\lambda_{1} \vee \lambda_{2}\right)}}
$$

Using these estimates, we find

$$
c_{p, n}^{1,1} \leq c_{1} e^{-\lambda_{1}(n-p)}+c^{2,1} c_{1} \tau^{1} \sum_{p \leq q<n} e^{-\lambda(q-p)} e^{-\lambda_{1}(n-(q+1))}
$$

Since $\lambda_{1}>\lambda$ we find that

$$
c_{p, n}^{1,1} \leq c_{1} e^{-\lambda_{1}(n-p)}+c^{2,1} c_{1} \tau^{1} e^{-\lambda((n-1)-p)} /\left(1-e^{-\Delta^{\prime}}\right) \quad \text { with } \quad \Delta^{\prime}=\lambda_{1}-\lambda>0
$$

This yields

$$
c_{p, n}^{1,1} \leq c^{1,1} e^{-\lambda(n-p)} \quad \text { with } \quad c^{1,1}:=c_{1}\left(1+c^{2,1} \tau^{1} /\left(e^{-\lambda}-e^{-\lambda_{1}}\right)\right)
$$

Finally, we observe that

$$
c_{p, n}^{1,2}=c \tau^{1} \sum_{p \leq q<n} e^{-\lambda(q-p)} e^{-\lambda_{1}(n-(q+1))} \leq c \tau^{1} e^{-\lambda((n-1)-p)} /\left(1-e^{-\Delta^{\prime}}\right)
$$

which implies that $c_{p, n}^{1,2} \leq c^{1,2} e^{-\lambda(n-p)}$, with $c^{1,2}:=c \tau^{1} /\left(e^{-\lambda}-e^{-\lambda_{1}}\right)$. This ends the proof of the lemma.

### 3.2 Feynman-Kac models

We let $Q_{p, n}$, with $0 \leq p \leq n$, be the Feynman-Kac semi-group associated with a sequence of bounded and positive integral operator $Q_{n}$ from some measurable spaces ( $E_{n-1}, \mathcal{E}_{n-1}$ ) into $\left(E_{n}, \mathcal{E}_{n}\right)$. For any $n \geq 1$, we denote by $G_{n-1}$ and $M_{n}$ the potential function on $E_{n-1}$ and the Markov transition from $E_{n-1}$ into $E_{n}$ defined below

$$
G_{n-1}(x)=Q_{n}(1)(x) \quad \text { and } \quad M_{n}(f)(x)=\frac{Q_{n}(f)(x)}{Q_{n}(1)(x)}
$$

We also denote by $\Phi_{p, n}, 0 \leq p \leq n$, the nonlinear semigroup from $\mathcal{P}\left(E_{p}\right)$ into $\mathcal{P}\left(E_{n}\right)$ defined below

$$
\begin{equation*}
\forall \eta \in \mathcal{P}\left(E_{p}\right), \forall f \in \mathcal{B}\left(E_{n}\right) \quad \Phi_{p, n}(\eta)(f)=\eta Q_{p, n}(f) / \eta Q_{p, n}(1) \tag{3.8}
\end{equation*}
$$

As usual we use the convention $\Phi_{n, n}=I d$, for $p=n$. It is important to observe that this semigroup is alternatively defined by the formulae

$$
\Phi_{p, n}(\eta)(f)=\frac{\eta\left(G_{p, n} P_{p, n}(f)\right)}{\eta\left(G_{p, n}\right)} \quad \text { with } \quad G_{p, n}=Q_{p, n}(1) \quad \text { and } \quad P_{p, n}\left(f_{n}\right)=Q_{p, n}\left(f_{n}\right) / Q_{p, n}(1)
$$

The next two parameters

$$
\begin{equation*}
r_{p, n}=\sup _{x, x^{\prime} \in E_{p}} \frac{G_{p, n}(x)}{G_{p, n}\left(x^{\prime}\right)} \quad \text { and } \quad \beta\left(P_{p, n}\right)=\sup _{x_{p}, y_{p} \in E_{p}}\left\|P_{p, n}\left(x_{p}, .\right)-P_{p, n}\left(y_{p}, .\right)\right\|_{\mathrm{tv}} \tag{3.9}
\end{equation*}
$$

measure respectively the relative oscillations of the potential functions $G_{p, n}$ and the contraction properties of the Markov transition $P_{p, n}$. Various estimates in the forthcoming sections will be expressed in terms of these parameters. For instance and for further use in several places in this article, we have the following Lipschitz regularity property.

Proposition 3.5 ([8]) For any $f_{n} \in O s c_{1}\left(E_{n}\right)$ we have

$$
\begin{equation*}
\left|\left[\Phi_{p, n}\left(\eta_{p}\right)-\Phi_{p, n}\left(\mu_{p}\right)\right]\left(f_{n}\right)\right| \leq 2 r_{p, n} \beta\left(P_{p, n}\right) \quad\left|\left[\eta_{p}-\mu_{p}\right] \bar{P}_{p, n}^{\mu_{p}}\left(f_{n}\right)\right| \tag{3.10}
\end{equation*}
$$

for some function $\bar{P}_{p, n}^{\mu_{p}}\left(f_{n}\right) \in O c_{1}\left(E_{p}\right)$ that doesn't depends on the measure $\eta_{p}$.
Our next objective is to estimate the the contraction coefficients $r_{p, n}$ and $\beta\left(P_{p, n}\right)$ in terms of the mixing type properties of the semigroup

$$
M_{p, n}\left(x_{p}, d x_{n}\right):=M_{p+1} M_{p+2} \ldots M_{n}\left(x_{p}, d x_{n}\right)
$$

associated with the Markov operators $M_{n}$. We introduce the following regularity condition.
$(M G)_{m}$ There exists an integer $m \geq 1$ and a sequence $\left(\epsilon_{p}(M)\right)_{p \geq 0} \in(0,1)^{\mathbb{N}}$ and some finite constant $r_{p}$ such that for any $p \geq 0$ and any $\left(x, x^{\prime}\right) \in E_{p}^{2}$ we have

$$
\begin{equation*}
M_{p, p+m}\left(x_{p}, .\right) \geq \epsilon_{p}(m) \quad M_{p, p+m}\left(x_{p}^{\prime}, .\right) \quad \text { and } \quad G_{p}(x) \leq r_{p} G_{n}\left(x^{\prime}\right) \tag{3.11}
\end{equation*}
$$

It is well known that the above condition is satisfied for any aperiodic and irreducible Markov chains on finite spaces, as well as for regular Markov processes on compact spaces, and for non compact spaces this condition is related to the tails of the transition distributions on the boundaries of the state space (see for instance [6]). Under condition $(M)_{m}$ we have for any $n \geq m \geq 1$, and $p \geq 1$,

$$
\begin{align*}
r_{p, p+n} & \leq \epsilon_{p}(m)^{-1} \prod_{0 \leq k<m} r_{p+k}  \tag{3.12}\\
\beta\left(P_{p, p+n}\right) & \leq \prod_{k=0}^{\lfloor n / m\rfloor-1}\left(1-\epsilon_{p+k m}^{(m)}\right) \quad \text { with } \quad \epsilon_{p}^{(m)}:=\epsilon_{p}^{2}(m) \prod_{0<k<m} r_{p+k}^{-1} \tag{3.13}
\end{align*}
$$

Notice that these estimates are also valid for any $n \geq 0$. Several contraction inequalities can be deduced from these estimates (see for instance chapter 4 of the book [6]). To give a flavor of these results, we further assume that $(M)_{m}$ is satisfied with $m=1$, and we have $\epsilon=\inf _{n} \epsilon_{n}(1)>0$. In this case, we can show that $r_{p, p+n} \leq r_{p} / \epsilon$ and $\beta\left(P_{p, p+n}\right) \leq\left(1-\epsilon^{2}\right)^{n}$. We end this short section with a direct consequence of proposition 3.5.
Corollary 3.6 Consider the Bernoulli semigroup presented in section 2.1. For constant mappings $s_{n}=\mu_{n+1}(1)$, the first component mapping is constant $\Phi_{n+1, \nu_{n}}^{1}(u)=s_{n}$ and the second component mapping $\Phi_{n+1, m_{n}}^{2}(\eta)=\Psi_{g_{n}^{(s)}}(\eta) M_{n+1}^{(s)}$ induces a Feynman-Kac semigroup with the likelihood function $g_{n}^{(s)}$ and the Markov transitions $M_{n+1}^{(s)}$ defined in 2.7). In this situation, the condition (3.2) is clearly met with $a_{p, n}^{1}=0$, for any $p<n$, We further assume that the semigroup of associated with the Markov transitions $M_{n}$ satisfies the mixing property stated in the l.h.s. of (3.11) for some integer $m \geq 1$ and some parameter $\left.\left.\epsilon_{p}(m) \in\right] 0,1\right]$. In this situation, the condition (3.3) is also met with the collection of parameters $a_{p, n}^{2}$ given below

$$
\begin{aligned}
a_{p, n}^{2} & \leq 2 \rho_{p}(m) \prod_{k=0}^{\lfloor(n-p) / m\rfloor-1}\left(1-\epsilon_{p+k m}^{(m, s)}\right) \\
\rho_{p}(m) & :=\epsilon_{p}^{-1}(m) \prod_{p \leq k<p+m} r_{k}^{2}\left(s_{k}\right) r_{k}(1) \text { and } \epsilon_{p}^{(s, m)}=\epsilon_{p}^{2}(m) r_{p}\left(s_{p}\right) / \prod_{p \leq k<p+m} r_{k}\left(s_{k}\right)^{3} r_{k}(1)^{2}
\end{aligned}
$$

with the collection of parameters $r_{n}\left(s_{n}\right)$ defined by $r_{n}\left(s_{n}\right):=\frac{s_{n} g_{n}^{+}+\left(1-s_{n}\right)}{s_{n} g_{n}^{-}+\left(1-s_{n}\right)}\left(\leq r_{n}(1)\right)$.

### 3.3 Bernoulli models

This section is concerned with the contraction properties of the semigroups $\Phi_{p, n, \nu}^{1}$ and $\Phi_{p, n, m}^{2}$ associated with the Bernoulli filter discussed in section 2.1. Before proceeding, we provide a brief discussion on the oscillations of the likelihood functions $g_{n}$ given below

$$
g_{n}\left(x_{n}\right)=\left(1-d_{n}\left(x_{n}\right)\right)+d_{n}\left(x_{n}\right) \mathcal{Y}_{n}\left(l_{n}\left(x_{n}, \cdot\right) / h_{n}\right)
$$

in terms of some $[0,1]$-valued detection probability functions $d_{n}$, some local likelihood functions $l_{n}$, and some positive clutter intensity function $h_{n}$. The oscillations of these likelihood functions strongly depend on the nature of the functions $\left(d_{n}, h_{n}, l_{n}\right)$.

Assuming that $h_{n}^{-}>0$ we have

$$
\begin{equation*}
\left(1-d_{n}^{\circ,-}\right)+d_{n}^{\circ,-} \frac{l_{n}^{-}}{h_{n}^{+}} \mathcal{Y}_{n}(1) \leq g_{n}^{-} \leq g_{n}^{+} \leq\left(1-d_{n}^{\circ,+}\right)+d_{n}^{\circ,+} \frac{l_{n}^{+}}{h_{n}^{-}} \mathcal{Y}_{n}(1) \tag{3.14}
\end{equation*}
$$

with the parameters

$$
\begin{aligned}
& d_{n}^{0,+}=d_{n}^{+} 1_{l_{n}^{+} \mathcal{Y}_{n}(1) \geq h_{n}^{-}}+d_{n}^{-} 1_{l_{n}^{+} \mathcal{Y}_{n}(1)<h_{n}^{-}} \\
& d_{n}^{\circ,-}=d_{n}^{-} 1_{l_{n}^{-}} \mathcal{Y}_{n}(1) \geq h_{n}^{+}+d_{n}^{+} 1_{l_{n}^{-}} \mathcal{Y}_{n}(1)<h_{n}^{+}
\end{aligned}
$$

The semigroup contraction inequalities developed in this section will be expressed in terms of the following parameters

$$
\delta_{n}(s g):=\frac{g_{n}^{+} s_{n}^{+}}{g_{n}^{-} s_{n}^{-}}, \quad \delta_{n}(g):=\frac{g_{n}^{+}}{g_{n}^{-}} \quad \text { and } \quad \delta_{n}^{\prime}(g):=\frac{1}{g_{n}^{-}} \wedge g_{n}^{+}
$$

For time homogeneous models $\left(d_{n}, h_{n}, l_{n}\right)=(d, h, l)$, with constant detection probability $d_{n}(x)=d$ and uniformly bounded number of observations $\sup _{n} \mathcal{Y}_{n}(1) \leq \mathcal{Y}^{+}(1)<\infty$ we have the following estimates

$$
(1-d) \leq g_{n}^{-} \leq g_{n}^{+} \leq(1-d)+d \frac{l^{+}}{h^{-}} \mathcal{Y}^{+}(1)
$$

In this situation, we have

$$
\delta_{n}(g) \leq 1+\frac{d}{1-d} \frac{l^{+}}{h^{-}} \mathcal{Y}^{+}(1)
$$

For small clutter intensity function with $h^{-}>0$ and $l^{-}>0$ we also have the observation free estimates $\frac{g_{n}^{+}}{g_{n}^{-}} \leq \frac{l^{+} h^{+}}{l^{-} h^{-}}$, from which we find that the upper bound

$$
\begin{equation*}
\delta(g):=\sup _{n \geq 0} \delta_{n}(g) \leq \inf \left\{1+\frac{d}{1-d} \frac{l^{+}}{h^{-}} \mathcal{Y}(1), \frac{l^{+} h^{+}}{l^{-} h^{-}}\right\} \tag{3.15}
\end{equation*}
$$

and for $d<1$

$$
\begin{equation*}
\delta^{\prime}(g):=\sup _{n \geq 0} \delta_{n}^{\prime}(g) \leq \sup \left\{(1-d)+d \frac{l^{+}}{h^{-}} \mathcal{Y}(1), \frac{1}{1-d}\right\} \tag{3.16}
\end{equation*}
$$

To be more precise, if we set $\inf _{n} \mathcal{Y}_{n}(1)=\mathcal{Y}^{-}(1)$ then

$$
1 \leq \frac{l^{-}}{h^{+}} \mathcal{Y}(1)^{-} \Rightarrow \delta^{\prime}(g) \leq(1-d)+d \frac{l^{+}}{h^{-}} \mathcal{Y}^{+}(1)
$$

In addition, if we have $d(1-d) \mathcal{Y}(1) \leq h^{-} / l^{+}$and $d<1$ then we find the the observation free estimates

$$
d \mathcal{Y}(1) l^{+} / h^{-} \leq 1 /(1-d) \Rightarrow \delta^{\prime}(g) \leq(1-d)+\frac{1}{1-d}
$$

Conversely, we have the observation free estimates

$$
\frac{l^{+}}{h^{-}} \mathcal{Y}(1)^{+} \leq 1 \Rightarrow \delta^{\prime}(g) \leq \frac{1}{(1-d)+d \frac{l^{-}}{h^{+}} \mathcal{Y}^{-}(1)} \leq \frac{1}{1-d}
$$

We are now in position to state the main result of this section.
Theorem 3.7 If $\left.\mu_{n+1}(1) \in\right] 0,1\left[, 0<s_{n}^{-} \leq s_{n}^{+}<1\right.$, and the semigroup $M_{p, n}$ satisfies the condition stated in the l.h.s. of (3.11) for some integer $m \geq 1$ and some positive constant $\epsilon_{p}(m)$, then the condition $(\operatorname{Lip}(\Phi))$ is met with

$$
a_{p, n}^{1} \leq 2 \epsilon_{p}^{-1} \delta_{p}^{\prime}(g) \prod_{p \leq k<p+n}\left(1-\epsilon_{k}^{2}\right) \quad \text { and } \quad a_{p, n}^{2} \leq 2 \rho_{p}(m) \prod_{k=0}^{\lfloor n / m\rfloor-1}\left(1-\epsilon_{p+k m}^{(m)}\right)
$$

with some parameters

$$
\epsilon_{n} \geq \inf \left\{\frac{s_{n}^{-}}{\mu_{n+1}(1)}, \frac{\mu_{n+1}(1)}{s_{n}^{+}}, \frac{1-s_{n}^{+}}{1-\mu_{n+1}(1)}, \frac{1-\mu_{n+1}(1)}{1-s_{n}^{-}}\right\}
$$

and

$$
\rho_{p}(m) \leq \epsilon_{p}(m)^{-1} \prod_{0 \leq k<m} \delta_{p+k}(s g)^{3} \quad \text { and } \quad \epsilon_{p}^{(m)} \geq \epsilon_{p}(m)^{2} \delta_{p}(s g)^{-4} \prod_{0<k<m} \delta_{p+k}(s g)^{-5}
$$

In addition condition $(\operatorname{Cont}(\Phi))$ is met with

$$
\tau_{n+1}^{1} \leq \delta_{n}(g)\left[\left(s_{n}^{+}-s_{n}^{-}\right)+\left\|s_{n}-\mu_{n+1}(1)\right\|\right] \quad \text { and } \quad \tau_{n+1}^{2} \leq \delta_{n}^{\prime}(g) \sup \left\{\frac{\mu_{n+1}}{s_{n}^{-}}, \frac{s_{n}^{+}}{\mu_{n+1}(1)}\right\}
$$

The proof of the theorem is postponed to section 5.2. To give a flavour of these estimates we examine time homogeneous models $\left(d_{n}, h_{n}, l_{n}, s_{n}, \mu_{n}\right)=(d, h, l, s, \mu)$, with constant detection and survival probabilities $d_{n}(x)=d, s_{n}(x)=s$, and uniformly bounded number of observations $\sup _{n} \mathcal{Y}_{n}(1) \leq \mathcal{Y}(1)<\infty$. In this situation, we have $\left(\epsilon_{p}(m), \epsilon_{p}^{(s)}(m)\right)=$ $\left(\epsilon(m), \epsilon^{(s)}(m)\right)$ and using the estimates 3.15 we prove the following bounds

$$
\tau_{n+1}^{1} \leq \delta(g)|s-\mu(1)| \quad \text { and } \quad \tau_{n+1}^{2} \leq \delta^{\prime}(g) \frac{\mu(1) \vee s}{\mu(1) \wedge s}
$$

as well as

$$
a_{0, n}^{1} \leq 2 \epsilon^{-1} \delta^{\prime}(g)\left(1-\epsilon^{2}\right)^{n} \quad \text { and } \quad a_{0, n}^{2} \leq 2 \epsilon(m)^{-1} \delta(g)^{3 m}\left(1-\epsilon(m)^{2} \delta(g)^{-5 m+1}\right)^{\lfloor n / m\rfloor}
$$

with some parameter $\epsilon$ such that

$$
\inf \left\{\frac{s}{\mu(1)}, \frac{\mu(1)}{s}, \frac{1-s}{1-\mu(1)}, \frac{1-\mu(1)}{1-s}\right\} \leq \epsilon \leq 1
$$

It is also readily verified that the assumptions of lemma 3.4 are satisfied with the parameters

$$
\begin{array}{ll}
\tau^{1} \leq \delta(g)|s-\mu(1)| & \tau^{2} \leq \delta^{\prime}(g)((\mu(1) \vee s) /(\mu(1) \wedge s)) \\
c_{1}=2 \epsilon^{-1} \delta^{\prime}(g) & c_{2}=2 \epsilon(m)^{-1}\left(1-\epsilon(m)^{2} \delta(g)^{-5 m+1}\right)^{-1} \delta(g)^{3 m}
\end{array}
$$

and the Lyapunov constants

$$
\lambda_{1}=-\log \left(1-\epsilon^{2}\right) \quad \text { and } \quad \lambda_{2}=-\frac{1}{m} \log \left(1-\epsilon(m)^{2} \delta(g)^{-5 m+1}\right)
$$

We notice that $\epsilon$ tends to 1 and $\tau^{1}$ tends to 0 , as $|s-\mu(1)|$ tends to 0 . Thus, there exists some $\varsigma \geq 0$ such that

$$
\lambda_{1}>\lambda_{2} \quad \text { and } \quad c_{1} c_{2} \tau^{1} \tau^{2}<\left(1-e^{-\lambda_{2}}\right) \quad\left(e^{-\lambda_{2}}-e^{-\lambda_{1}}\right)
$$

as long as $|s-\mu(1)| \leq \varsigma$. We summarize this discussion with the following corollary.
Corollary 3.8 Consider the time homogeneous model discussed above. Under the assumptions of theorem 3.7, for any $N$-approximation models $\left(\gamma_{n}^{N}(1), \eta_{n}^{N}\right)$ of the Bernoulli model $\left(\gamma_{n}(1), \eta_{n}\right)$ satisfying condition (1.6), the $\mathbb{L}_{r}$-mean error estimates presented in (1.16) are uniform w.r.t. the time parameter

$$
\sup _{n \geq 0} \mathbb{E}\left(\left|V_{n}^{\gamma, N}(1)\right|^{r}\right)^{\frac{1}{r}} \leq a_{r} c^{1,2} /\left(1-e^{-\lambda}\right) \quad \text { and } \quad \sup _{n \geq 0} \mathbb{E}\left(\left|V_{n}^{\eta, N}(f)\right|^{r}\right)^{\frac{1}{r}} \leq a_{r} c^{2,2} /\left(1-e^{-\lambda}\right)
$$

with the parameters $\left(c^{1,2}, c^{2,2}, \lambda\right)$ defined in lemma 3.4, and some finite constants $a_{r}<\infty$ whose values only depend on $r$.

Remark 3.9 When $\mu_{n+1}(1)=0$ we have seen in 2.8) that

$$
\Phi_{n+1, \nu_{n}}^{1}(u)=\Psi_{g_{n}}\left(\nu_{n}\right)\left(s_{n}\right) \times \theta_{\nu_{n}\left(g_{n}\right)}(u) \quad \text { and } \quad \Phi_{n+1, m_{n}}^{2}(\eta)=\Psi_{g_{n} s_{n}}(\eta) M_{n+1}
$$

with the collection of mappings $\theta_{a}$, with $a \in[0, \infty[$, defined in (2.9). Using the fact that

$$
\left|\Phi_{n+1, \nu_{n}}^{1}(u)-\Phi_{n+1, \nu_{n}}^{1}\left(u^{\prime}\right)\right|=\frac{\Psi_{g_{n}}\left(\nu_{n}\right)\left(s_{n}\right) \nu_{n}\left(g_{n}\right)}{\left[\nu\left(g_{n}\right) u+(1-u)\right]\left[\nu\left(g_{n}\right) u^{\prime}+\left(1-u^{\prime}\right)\right]}\left|u-u^{\prime}\right|
$$

one proves that (3.2) is met with $a_{p, n}^{1} \leq \prod_{p \leq k<n}\left(s_{k}^{+} g_{k}^{+}\right) /\left(1 \wedge g_{k}^{-}\right)^{2}$. We also notice that the second component mapping $\Phi_{n+1, m_{n}}^{2}$ doesn't depends on the parameter $m_{n}$, and it induces a Feynman-Kac semigroup of the same form as the one studied in section 3.2. Assuming that the mixing condition stated in the l.h.s. of (3.11) is satisfied some integer $m \geq 1$ and some parameter $\epsilon_{p}(m)>0$, one can prove that (3.3) is met with the collection of parameters $a_{p, n}^{2}$ given below

$$
a_{p, n}^{2} \leq 2 \rho_{p}(m) \prod_{k=0}^{\lfloor(n-p) / m\rfloor-1}\left(1-\epsilon_{p+k m}^{(m)}\right) \quad \text { with } \quad \rho_{p}(m)=\epsilon_{p}^{-1}(m) \prod_{p \leq q<p+m} \delta_{q}(s g)
$$

and the collection of parameters $\epsilon_{p}^{(m)}=\epsilon_{p}^{(m)}=\epsilon_{p}^{2}(m) / \prod_{p<q<p+m} \delta_{q}(\mathrm{sg})$.

### 3.4 PHD Models

This section is concerned with the contraction properties of the semigroups $\Phi_{p, n, \nu}^{1}$ and $\Phi_{p, n, m}^{2}$ associated with the PHD filter discussed in section 1.1.2 and in section 2.2.

The analysis of these nonlinear models is much more involved than the one of the Bernoulli models. We simplify the analysis and we further assume that the clutter intensity function, the detectability rate as well as the survival and the spawning rates introduced in section 2.2 are time homogeneous and constants functions, and we set

$$
\left(b_{n}(x), h_{n}(x), s_{n}(x), r_{n}(x)\right)=(b, h, s, r)
$$

To simplify the presentation, we also assume that the state spaces, the Markov transitions of the targets, the likelihood functions and the spontaneous birth measures are time homogeneous, that is we have that $E_{n}=E, E_{n}^{Y}=E^{Y}, M_{n}=M, g_{n}(x, y)=g(x, y)$ and $\mu_{n+1}=\mu$. Without further mention, we suppose that $r(1-d)<1, \mu(1)>0, r>0$, and for any $y \in E^{Y}$ we have

$$
0 \leq g^{-}(y):=\inf _{x \in E} g(x, y) \leq g^{+}(y):=\sup _{x \in E} g(x, y)<\infty
$$

Given a mapping $\theta$ from $E^{Y}$ into $\mathbb{R}$, we set $\mathcal{Y}^{-}(\theta):=\inf _{n} \mathcal{Y}_{n}(\theta)$ and $\mathcal{Y}^{+}(\theta):=\sup _{n} \mathcal{Y}_{n}(\theta)$.
We recall from (1.10) that the PDH filter is defined by the measure-valued equation $\gamma_{n+1}=\gamma_{n} Q_{n+1, \gamma_{n}}$, with the integral operator

$$
Q_{n+1, \gamma_{n}}\left(x_{n}, d x_{n+1}\right)=g_{n, \gamma_{n}}\left(x_{n}\right) M_{n+1}\left(x_{n}, d x_{n+1}\right)+\gamma_{n}(1)^{-1} \mu_{n+1}\left(d x_{n+1}\right)
$$

and the function $g_{n, \gamma_{n}}$ defined below

$$
g_{n, \gamma_{n}}(x)=r(1-d)+r d \int \mathcal{Y}_{n}(d y) \frac{g(x, y)}{h+d \gamma_{n}(g(., y))}
$$

We also notice that the total mass process and the normalized distribution flow are given by the following equations

$$
\begin{aligned}
\gamma_{n+1}(1) & =\Phi_{n+1, \eta_{n}}^{1}\left(\gamma_{n}(1)\right) \\
& =\gamma_{n}(1) r(1-d)+\int \mathcal{Y}_{n}(d y) w_{\gamma_{n}(1)}\left(\eta_{n}, y\right)+\mu(1) \\
\eta_{n+1}(1) & =\Phi_{n+1, \gamma_{n}(1)}^{2}\left(\eta_{n}\right) \\
& \propto \gamma_{n}(1) r(1-d) \eta_{n} M+\int \mathcal{Y}_{n}(d y) w_{\gamma_{n}(1)}\left(\eta_{n}, y\right) \Psi_{g(\cdot, y)}\left(\eta_{n}\right) M+\mu(1) \bar{\mu}
\end{aligned}
$$

with the probability measure $\bar{\mu}$ and weight functions $w$ defined below

$$
\bar{\mu}(d x)=\mu(d x) / \mu(1) \quad \text { and } \quad w_{u}(\eta, y):=r\left(1-\frac{h}{h+\operatorname{du\eta }(g(\cdot, y))}\right)
$$

For null clutter parameter $h=0$, we already observe that the total mass transformation $\Phi_{n+1, \eta_{n}}^{1}$ doesn't depend on the flow of probability measures $\eta_{n}$ and it is simply given by

$$
\Phi_{n+1, \eta_{n}}^{1}\left(\gamma_{n}(1)\right)=\gamma_{n}(1) r(1-d)+r \mathcal{Y}_{n}(1)+\mu(1)
$$

In this particular situation, we have

$$
\gamma_{n}^{N}(1)=\gamma_{n}(1)=(r(1-d))^{n} \gamma_{0}(1)+\sum_{0 \leq k<n}(r(1-d))^{n-1-k}\left(r \mathcal{Y}_{k}(1)+\mu(1)\right)
$$

Now, we easily show that the pair of conditions (3.2) and (3.4) are satisfied with the parameters $a_{p, n}^{1}=(r(1-d))^{n-p}$ and $\tau_{n}^{1}=0$. In more general situations, the total mass process is not explicitly known. Some useful estimates are provided by the following lemma.

Lemma 3.10 We assume that the number of observations is uniformly bounded; that is, we have that $\mathcal{Y}^{+}(1)<\infty$. In this situation, the total mass process $\gamma_{n}(1)$ and any approximation model $\gamma_{n}^{N}(1)$ given by the recursion (1.4) (with the initial condition $\gamma_{0}^{N}(1)=\gamma_{0}(1)$ ) take values in a sequence of compact sets $I_{n} \subset\left[m^{-}, m^{+}\right]$with

$$
m^{-}:=\frac{\mu(1)}{1-r(1-d)}\left(1+r d \mathcal{Y}^{-}\left(\frac{g^{-}}{h+d \mu(1) g^{-}}\right)\right) \quad \text { and } \quad m^{+}:=\gamma_{0}(1)+\frac{r \mathcal{Y}^{+}(1)+\mu(1)}{1-r(1-d)}
$$

## Proof:

Using the fact that $\gamma_{n}(1) \geq \mu(1)$ we prove that $r\left(1-\frac{h}{h+d \mu(1) g^{-}(y)}\right) \leq w_{\gamma_{n}(1)}\left(\eta_{n}, y\right) \leq r$, from which we conclude that

$$
\gamma_{n}(1) r(1-d)+r \mathcal{Y}_{h, n}(1)+\mu(1) \leq \Phi_{n+1, \eta_{n}}^{1}\left(\gamma_{n}(1)\right) \leq \gamma_{n}(1) r(1-d)+r \mathcal{Y}_{n}(1)+\mu(1)
$$

with the random measures

$$
\mathcal{Y}_{h, n}(d y):=\mathcal{Y}_{n}(d y) \frac{d \mu(1) g^{-}(y)}{h+d \mu(1) g^{-}(y)}
$$

For any sequence of probability measures $\nu:=\left(\nu_{n}\right)_{n \geq 0} \in \mathcal{P}(E)^{\mathbb{N}}$, and any starting mass $u \in[0, \infty[$ one conclude that

$$
(r(1-d))^{n} u+\frac{r \mathcal{Y}_{h}^{-}(1)+\mu(1)}{1-r(1-d)} \leq \Phi_{0, n, \nu}^{1}(u) \leq(r(1-d))^{n} u+\frac{r \mathcal{Y}^{+}(1)+\mu(1)}{1-r(1-d)}
$$

This implies that $\gamma_{n}(1), \gamma_{n}^{N}(1) \in I_{n} \subset\left[m^{-}, m^{+}\right]$with

$$
m^{-}:=\frac{r \mathcal{Y}_{h}^{-}(1)+\mu(1)}{1-r(1-d)}=\frac{\mu(1)}{1-r(1-d)}\left(1+r d \mathcal{Y}^{-}\left(\frac{g^{-}}{h+d \mu(1) g^{-}}\right)\right)
$$

The end of the proof of the lemma is now completed.
We are now in position to state the main result of this section.
Theorem 3.11 We assume that the number of observations is uniformly bounded; that is, we have that $\mathcal{Y}^{+}(1)<\infty$. In this situation, the condition $(\operatorname{Lip}(\Phi))$ is met with the Lipschitz constants $a_{p, n}^{i} \leq \prod_{p \leq k<n} a_{k, k+1}^{i}$, with $i=1,2$, and the sequence of parameters $\left(a_{n, n+1}^{i}\right)_{n \geq 0}$, $i=1,2$, defined below

$$
\begin{aligned}
& a_{n, n+1}^{1} \leq r(1-d)+r d h \mathcal{Y}_{n}\left(\frac{g^{+}}{\left[h+d m^{-} g^{-}\right]^{2}}\right) \\
& a_{n, n+1}^{2} \leq m^{+} \frac{\beta(M)\left[(1-d)+d \mathcal{Y}_{n}\left(\frac{g^{+}}{h+d m^{+} g^{+}} \frac{g^{+}}{g^{-}}\right)\right]+h d \mathcal{Y}_{n}\left(\frac{g^{+}-g^{-}}{\left(h+d m^{-} g^{-}\right)^{2}}\right)}{(1-d) m^{-}+d m^{-} \mathcal{Y}_{n}\left(\frac{g^{-}}{h+d m^{-} g^{-}}\right)+\mu(1) / r}
\end{aligned}
$$

In addition, condition $(\operatorname{Cont}(\Phi))$ is met with the sequence of parameters

$$
\tau_{n+1}^{1} \leq r d h m^{+} \mathcal{Y}_{n}\left(\frac{g^{+}-g^{-}}{\left[h+d m^{-} g^{-}\right]^{2}}\right) \quad \tau_{n+1}^{2} \leq \frac{(1-d)+h d \mathcal{Y}_{n}\left(\frac{g^{+}}{\left(h+d m^{-} g^{-}\right)^{2}}\right)}{(1-d) m^{-}+d m^{-} \mathcal{Y}_{n}\left(\frac{g^{-}}{h+d m^{-} g^{-}}\right)+\mu(1) / r}
$$

The proof of theorem 3.11 is postponed to section 5.4 .
Corollary 3.12 We assume that $\mathcal{Y}^{+}\left(g^{+} / g^{-}\right) \vee \mathcal{Y}^{+}\left(g^{+} /\left(g^{-}\right)^{2}\right)<\infty$. Then there exists some parameters $0<\kappa_{0} \leq 1, \kappa_{1}<\infty$, and $\kappa_{2}>0$ such that for any $d \geq \kappa_{0}, \mu(1) \geq$ $\kappa_{1}$, and $h \leq \kappa_{2}$, the semigroups $\Phi_{p, n, \nu}^{1}$ and $\Phi_{p, n, m}^{2}$ satisfy the pair of conditions $($ Lip $(\Phi))$ and $(\operatorname{Cont}(\Phi))$ with some $\left(a_{p, n}^{i}, \tau_{n}^{i}\right)_{i=1,2, p \leq n, ~ s a t i s f y i n g ~ t h e ~ a s s u m p t i o n s ~ o f ~ l e m m a ~ 3.4 . ~ I n ~}^{n}$ particular, for any $N$-approximation models $\left(\gamma_{n}^{N}(1), \eta_{n}^{N}\right)$ of the $\operatorname{PHD}$ equation $\left(\gamma_{n}(1), \eta_{n}\right)$ satisfying condition (1.6), the uniform $\mathbb{L}_{r}$-mean error estimates presented in corollary 3.8 are satisfied.

## Proof:

There is no loss of generality to assume that $r(1-d)<1 / 2 \leq d$ and $\mu(1) \geq 1 \geq h$. Recalling that $m^{-} \geq \mu(1)$, one readily proves that

$$
\frac{m^{+}}{\mu(1)}=\frac{\gamma_{0}(1)}{\mu(1)}+\frac{1}{1-r(1-d)}\left(1+\frac{r}{\mu(1)} \mathcal{Y}^{+}(1)\right) \leq 2+\gamma_{0}(1)+2 r \mathcal{Y}^{+}(1):=\rho
$$

If we set $\delta(g):=\rho \vee \mathcal{Y}^{+}\left(\frac{g^{+}}{g^{-}}\right) \vee \mathcal{Y}^{+}\left(\frac{g^{+}}{\left(g^{-}\right)^{2}}\right)$, then we find the rather crude estimates

$$
a_{n, n+1}^{1} / r \leq(1-d)+\frac{2 h}{\mu(1)^{2}} \delta(g) \quad \text { and } \quad a_{n, n+1}^{2} / r \leq\left[\beta(M)(1-d)+\frac{2 h+\beta(M)}{\mu(1)}\right] \delta(g)
$$

as well as

$$
\tau_{n+1}^{1} / r \leq \frac{2 h}{\mu(1)} \delta(g)^{2} \quad \text { and } \quad \tau_{n+1}^{2} / r \leq \frac{1}{\mu(1)}\left[(1-d)+\frac{2 h}{\mu(1)^{2}} \delta(g)\right]
$$

from which we find that

$$
\begin{equation*}
\tau^{1} \tau^{2} \leq \frac{2 h r^{2}}{\mu(1)^{2}}\left[(1-d)+\frac{2 h}{\mu(1)^{2}} \delta(g)\right] \delta(g)^{2} \tag{3.17}
\end{equation*}
$$

Thus, there exists some $0<\kappa_{0} \leq 1$ and some $\kappa_{1}<\infty$ so that for any $d \geq \kappa_{0}$ and any $\mu(1) \geq \kappa_{1}$ we have

$$
\begin{aligned}
& a_{n, n+1}^{1} \leq r\left[(1-d)+\frac{2}{\mu(1)^{2}}\right] \delta(g):=e^{-\lambda_{1}}<1 \\
& a_{n, n+1}^{2} \leq r\left[(1-d)+\frac{3}{\mu(1)}\right] \delta(g):=e^{-\lambda_{2}}<1 \quad \text { with } \quad 0<\lambda_{2}<\lambda_{1}
\end{aligned}
$$

Finally, using (3.17) we find some $\kappa_{2}>0$ such that for any $h \leq \kappa_{2}$, we have that $\tau^{1} \tau^{2} \leq$ $\left(1-e^{-\lambda_{2}}\right)\left(e^{-\lambda_{2}}-e^{-\lambda_{1}}\right)$. The end of the proof is now a direct consequence of lemma 3.4 . This ends the proof of the corollary.

## 4 Stochastic particle approximations

### 4.1 Mean field interacting particle systems

### 4.1.1 Description of the models

The mean field type interacting particle system associated with the equation (1.2) relies on the fact that the one step mappings $\Gamma_{n+1}^{2}$ can be rewritten in the following form

$$
\begin{equation*}
\Gamma_{n+1}^{2}\left(\gamma_{n}(1), \eta_{n}\right)=\eta_{n} K_{n+1, \gamma_{n}} \quad \text { with } \quad \gamma_{n}=\gamma_{n}(1) \times \eta_{n} \tag{4.1}
\end{equation*}
$$

for some collection of Markov kernels $K_{n+1, \gamma}$ indexed by the time parameter $n$ and the set of measures $\gamma \in \mathcal{M}_{+}\left(E_{n}\right)$. We mention that the choice of the Markov transitions $K_{n, \gamma}$ is not unique. In the literature on mean field particle models, $K_{n, \gamma}$ are called a choice of McKean transitions. Some McKean interpretation models of the Bernoulli and the PHD filter models $(1.8)$ and 1.10 are discussed in section 2.2 (see for instance 2.10 ) and in section 2.1 (see for instance 2.2

These models provide a natural interpretation of the distribution laws $\eta_{n}$ as the laws of a non linear Markov chain $\bar{X}_{n}$ whose elementary transitions $\bar{X}_{n} \rightsquigarrow \bar{X}_{n+1}$ depends on the distribution $\eta_{n}=\operatorname{Law}\left(\bar{X}_{n}\right)$, as well as on the current mass process $\gamma_{n}(1)$. In contrast to traditional McKean model, the dependency on the mass process induce a dependency of all the flow of measures $\eta_{p}$, for $0 \leq p \leq n$. For a thorough description of these discrete generation and non linear McKean type models, we refer the reader to [6].

In further developments of the article, we always assume that the mappings

$$
\left(m, x_{n},\left(x^{i}\right)_{1 \leq i \leq N}\right) \mapsto K_{n+1, m \sum_{j=1}^{N} \delta_{x^{j}}}\left(x_{n}, A_{n+1}\right) \text { and } G_{n+1, m \sum_{j=1}^{N} \delta_{x j}^{j}}\left(x_{n}\right)
$$

are pointwise known, and of course measurable w.r.t. the corresponding product sigma fields, for any $n \geq 0, N \geq 1, A_{n+1} \in \mathcal{E}_{n+1}$, and any $x_{n} \in E_{n}$. In this situation, the mean field particle interpretation of this nonlinear measure-valued model is an $E_{n}^{N}$-valued Markov chain $\xi_{n}^{(N)}=\left(\xi_{n}^{(N, i)}\right)_{1 \leq i \leq N}$, with elementary transitions defined as

$$
\begin{align*}
\gamma_{n+1}^{N}(1) & =\gamma_{n}^{N}(1) \eta_{n}^{N}\left(G_{\left.n, \gamma_{n}^{N}\right)}\right.  \tag{4.2}\\
\mathbb{P}\left(\xi_{n+1}^{(N)} \in d x \mid \mathcal{F}_{n}^{(N)}\right) & =\prod_{i=1}^{N} K_{n+1, \gamma_{n}^{N}}\left(\xi_{n}^{(N, i)}, d x^{i}\right) \tag{4.3}
\end{align*}
$$

with the pair of occupation measures $\left(\gamma_{n}^{N}, \eta_{n}^{N}\right)$ defined below

$$
\eta_{n}^{N}:=\frac{1}{N} \sum_{j=1}^{N} \delta_{\xi_{n}^{(N, j)}} \quad \text { and } \quad \gamma_{n}^{N}(d x):=\gamma_{n}^{N}(1) \eta_{n}^{N}(d x)
$$

In the above displayed formula, $\mathcal{F}_{n}^{N}$ stands for the $\sigma$-field generated by the random sequence $\left(\xi_{p}^{(N)}\right)_{0 \leq p \leq n}$, and $d x=d x^{1} \times \ldots \times d x^{N}$ stands for an infinitesimal neighborhood of a point $x=\left(x^{1}, \ldots, x^{N}\right) \in E_{n}^{N}$. The initial system $\xi_{0}^{(N)}$ consists of $N$ independent and identically distributed random variables with common law $\eta_{0}$. As usual, to simplify the presentation, when there is no possible confusion we suppress the parameter $N$, so that we write $\xi_{n}$ and $\xi_{n}^{i}$ instead of $\xi_{n}^{(N)}$ and $\xi_{n}^{(N, i)}$.

### 4.1.2 Convergence analysis

The rationale behind the mean field particle model described in 4.3) is that $\eta_{n+1}^{N}$ is the empirical measure associated with $N$ independent variables with distributions $K_{n+1, \gamma_{n}^{N}}^{N}\left(\xi_{n}^{i}, d x\right)$, so as long as $\gamma_{n}^{N}$ is a good approximation of $\gamma_{n}$ then $\eta_{n+1}^{N}$ should be a good approximation of $\eta_{n+1}$. Roughly speaking, this induction argument shows that $\eta_{n}^{N}$ tends to $\eta_{n}$, as the population size $N$ tends to infinity.

These stochastic particle algorithms can be thought of in various ways: From the physical view point, they can be seen as microscopic particle interpretations of physical nonlinear measure-valued equations. From the pure mathematical point of view, they can also be interpreted as natural stochastic linearizations of nonlinear evolution semigroups. From the probabilistic point of view, they can be interpreted as a interacting recycling acceptancerejection sampling techniques. In this case, they can be seen as a sequential and interacting importance sampling technique.

By construction, the local fluctuation random fields $\left(W_{n}^{N}\right)_{n \geq 0}$ defined in (1.5) can be rewritten as follows $\eta_{n}^{N}=\eta_{n-1}^{N} K_{n, \gamma_{n-1}^{N}}+\frac{1}{\sqrt{N}} W_{n}^{N}$. Using Khintchine's inequality, we can check that 1.6 is met for any $r \geq 1$ and any $f_{n} \in \operatorname{Osc}_{1}\left(E_{n}\right)$, with the collection of universal constants $a_{r}$ defined as in (1.17) by taking $b=1$.

### 4.2 Interacting particle association systems

### 4.2.1 Description of the models

We let $\left(\mathcal{A}_{n}\right)_{n \geq 0}$ be a sequence of finite sets equipped with some finite positive measures $\left(\nu_{n}\right)_{n \geq 0}$. We further assume that the initial distribution $\gamma_{0}$ and the integral operators $Q_{n+1, \gamma_{n}}$ in (1.1) have the following form

$$
\gamma_{0}=\int \nu_{0}(d a) \eta_{0}^{(a)} \quad \text { and } \quad Q_{n+1, \gamma_{n}}=\int \nu_{n+1}(d a) Q_{n+1, \gamma_{n}}^{(a)}
$$

In the above display $\eta_{0}^{(a)}$ stands for a collection of measures on $E_{0}$, indexed by the parameter $a \in \mathcal{A}_{0}$, and $Q_{n+1, \gamma_{n}}^{(a)}$ is a collection of integral operators indexed by the parameter $a \in \mathcal{A}_{n+1}$. In this situation, we observe that

$$
\gamma_{0}(1)=\nu_{0}(1) \quad \text { and } \quad \eta_{0}=\int A_{0}(d a) \eta_{0}^{(a)} \quad \text { with } \quad A_{0}(d a):=\nu_{0}(d a) / \nu_{0}(1)
$$

We also assume that the following property is met

$$
\begin{equation*}
G_{n, \gamma}^{(a)}:=Q_{n+1, \gamma}^{(a)}(1) \propto G_{n}^{(a)} \quad \text { and } \quad Q_{n+1, \gamma}^{(a)}(f) / Q_{n+1, \gamma}^{(a)}(1):=M_{n+1}^{(a)}(f) \tag{4.4}
\end{equation*}
$$

for some function $G_{n}^{(a)}$ on $E_{n}$, and some Markov transitions $M_{n+1}^{(a)}$ from $E_{n}$ into $E_{n+1}$ whose values do not depend on the measures $\gamma$. For clarity of presentation, sometimes we write $\Psi_{G_{n}}^{(a)}$ instead of $\Psi_{G_{n}^{(a)}}$.

Definition 4.1 We consider the collection of probability measures $\eta_{n}^{\left(a_{n}\right)} \in \mathcal{P}\left(E_{n}\right)$, indexed by sequences of parameters $\mathrm{a}_{n}=\left(a_{0}, \ldots, a_{n}\right) \in \mathcal{A}_{[0, n]}:=\left(\mathcal{A}_{0} \times \ldots \times \mathcal{A}_{n}\right)$, and defined by the following equations

$$
\begin{equation*}
\eta_{n}^{\left(a_{n}\right)}=\left(\Phi_{n}^{\left(a_{n}\right)} \circ \ldots \circ \Phi_{1}^{\left(a_{1}\right)}\right)\left(\eta_{0}^{\left(a_{0}\right)}\right) \tag{4.5}
\end{equation*}
$$

with the mappings $\Phi_{n}^{(a)}: \mathcal{P}\left(E_{n-1}\right) \rightarrow \mathcal{P}\left(E_{n}\right)$ indexed by $a \in \mathcal{A}_{n}$ and defined by

$$
\Phi_{n}^{(a)}(\eta)=\Psi_{G_{n-1}}^{(a)}(\eta) M_{n}^{(a)}
$$

We illustrate these abstract conditions in the context of the multiple target tracking equation presented in (1.10). In this situation, it is convenient to add a pair of virtual observation states $c, c^{\prime}$ to $E_{n}^{Y}$. Using this notation, the above conditions are satisfied with the finite sets $\mathcal{A}_{n+1}$ and their counting measures $\nu_{n+1}$ defined below

$$
\mathcal{A}_{n+1}=\left\{Y_{n}^{i}, 1 \leq i \leq N_{n}^{Y}\right\} \cup\left\{c, c^{\prime}\right\} \quad \nu_{n+1}=\mathcal{Y}_{n}+\delta_{c}+\delta_{c^{\prime}} \in \mathcal{M}\left(\mathcal{A}_{n+1}\right)
$$

Using (1.10) and (1.12), we check that (4.4) is met with the couple of potential functions and Markov transitions defined by

$$
\left(G_{n}^{\left(y_{n}\right)}, M_{n+1}^{\left(y_{n}\right)}\right)=\left\{\begin{array}{lll}
\left(r_{n} d_{n} g_{n}\left(., y_{n}\right), M_{n+1}\right) & \text { for } & y_{n} \notin\left\{c, c^{\prime}\right\} \\
\left(r_{n}\left(1-d_{n}\right), M_{n+1}\right) & \text { for } & y_{n}=c \\
\left(1, \bar{\mu}_{n+1}\right) & \text { for } & y_{n}=c^{\prime}
\end{array}\right.
$$

In this case, we observe that $Q_{n+1, \gamma_{n}}^{\left(y_{n}\right)}\left(x_{n},.\right)=G_{n, \gamma_{n}}^{\left(y_{n}\right)}\left(x_{n}\right) M_{n+1}^{\left(y_{n}\right)}\left(x_{n},.\right)$ with the potential function $G_{n, \gamma_{n}}^{\left(y_{n}\right)}$ defined below

$$
G_{n, \gamma_{n}}^{\left(y_{n}\right)} / G_{n}^{\left(y_{n}\right)}=\left\{\begin{array}{cll}
{\left[h_{n}\left(y_{n}\right)+\gamma_{n}\left(d_{n} g_{n}\left(., y_{n}\right)\right)\right]^{-1}} & \text { for } & y_{n} \notin\left\{c, c^{\prime}\right\}  \tag{4.6}\\
1 & \text { for } & y_{n}=c \\
\mu_{n+1}(1) / \gamma_{n}(1) & \text { for } & y_{n}=c^{\prime}
\end{array}\right.
$$

Under our assumptions, using (1.2), we have the following result.
Proposition 4.2 The solution the equation (1.2) has the following form $\eta_{n}=\int A_{n}(d a) \eta_{n}^{(a)}$, with a total mass process $\gamma_{n}(1)$ and the association measures $A_{n} \in \mathcal{P}\left(\mathcal{A}_{[0, n]}\right)$ defined by the following recursive equations

$$
\gamma_{n+1}(1)=\gamma_{n}(1) \eta_{n}\left(G_{n, \gamma_{n}}\right) \quad \text { and } \quad A_{n+1}=\Omega_{n+1}\left(\gamma_{n}(1), A_{n}\right)
$$

With the mapping

$$
\Omega_{n+1}:(m, A) \in(] 0, \infty\left[\times \mathcal{P}\left(\mathcal{A}_{[0, n]}\right)\right) \mapsto \Omega_{n+1}(m, A) \in \mathcal{P}\left(\mathcal{A}_{[0, n+1]}\right)
$$

defined by the following formula

$$
\begin{equation*}
\Omega_{n+1}(m, A)(d(a, b)) \propto A(d a) \nu_{n+1}(d b) \eta_{n}^{(a)}\left(G_{n, m \int A(d a) \eta_{n}^{(a)}}^{(b)}\right) \tag{4.7}
\end{equation*}
$$

## Proof:

The proof of the above assertion is simply based on the fact that

$$
\begin{aligned}
\eta_{n+1} \propto \int \nu_{n+1}(d b) \eta_{n} Q_{n+1, \gamma_{n}}^{(b)} & =\int A_{n}(d a) \nu_{n+1}(d b) \eta_{n}^{(a)} Q_{n+1, \gamma_{n}}^{(b)} \\
& =\int A_{n}(d a) \nu_{n+1}(d b) \eta_{n}^{(a)}\left(G_{n, \gamma_{n}}^{(b)}\right) \eta_{n+1}^{(a, b)}
\end{aligned}
$$

This clearly implies that

$$
\left.\Gamma_{n}^{2}\left(m, \int A(d a) \eta_{n-1}^{(a)}\right)\right)=\int \Omega_{n}(m, A)(d(a, b)) \eta_{n}^{(a, b)}
$$

This ends the proof of the proposition.

By construction, we notice that for any discrete measure $A \in \mathcal{P}\left(\mathcal{A}_{[0, n-1]}\right)$, and any collection of measures $\eta^{(a)} \in \mathcal{P}\left(E_{n-1}\right)$, with $a \in \mathcal{A}_{[0, n-1]}$ we have the formula

$$
\left.\Gamma_{n}^{2}\left(m, \int A(d a) \eta^{(a)}\right)\right)=\int \Omega_{n}(m, A)(d(a, b)) \Phi_{n}^{(b)}\left(\eta^{(a)}\right)
$$

### 4.2.2 Particle approximation models

To get some feasible solution, we further assume that $\eta_{n}^{(a)}\left(G_{n, \gamma_{n}}^{(b)}\right)$ are explicitly known for any sequence of parameters $(a, b) \in\left(\mathcal{A}_{[0, n]} \times \mathcal{A}_{n+1}\right)$. This rather strong condition is satisfied for the multiple target tracking model discussed above as long as the quantities

$$
\eta_{n}^{\left(a_{0}, y_{0}, \ldots, y_{n-1}\right)}\left(r_{n} d_{n} g_{n}\left(., y_{n}\right)\right) \quad \eta_{n}^{\left(a_{0}, y_{0}, \ldots, y_{n-1}\right)}\left(r_{n}\left(1-d_{n}\right)\right) \quad \eta_{n}^{\left(a_{0}, y_{0}, \ldots, y_{n-1}\right)}\left(d_{n} g_{n}\left(., y_{n}\right)\right)
$$

are explicitly known. This condition is clearly met for linear gaussian target evolution and observation sensors as long as the survival and detection probabilities $s_{n}$ and $d_{n}$ are state independent, and spontaneous birth $\bar{\mu}_{n}$ and spawned targets branching rates $b_{n}$ are Gaussian mixtures. In this situation, the collection of measures $\eta_{n}^{\left(a_{0}, y_{0}, \ldots, y_{n-1}\right)}$ are gaussian distributions and the equation (4.5) coincides with the traditional updating-prediction transitions of the discrete generation Kalman-Bucy filter.

We let $A_{0}^{N}=\frac{1}{N} \sum_{i=1}^{N} \delta_{\mathrm{a}_{0}^{i}}$, be the empirical measure associated with $N$ independent and identically distributed random variables $\left(\mathrm{a}_{0}^{i}\right)_{1 \leq i \leq N}$ with common distribution $A_{0}$. By construction, we have

$$
\eta_{0}^{N}:=\int A_{0}^{N}(d a) \eta_{0}^{(a)}=\eta_{0}+\frac{1}{\sqrt{N}} W_{0}^{N}
$$

with some local sampling random fields satisfying (1.6). We further assume that $\gamma_{0}(1)$ is known and we set $\gamma_{0}^{N}=\gamma_{0}(1) \eta_{0}^{N}$.

$$
\gamma_{1}^{N}(1)=\gamma_{0}^{N}(1) \eta_{0}^{N}\left(G_{0, \gamma_{0}^{N}}\right) \quad \text { and } \quad \eta_{1}^{N}:=\int A_{1}^{N}(d a) \eta_{1}^{(a)}
$$

with the occupation measure $A_{1}^{N}=\frac{1}{N} \sum_{i=1}^{N} \delta_{a_{1}^{i}}$ associated with $N$ conditionally independent and identically distributed random variables $\mathrm{a}_{1}^{i}:=\left(a_{0,1}^{i}, a_{1,1}^{i}\right)$ with common law $\Omega_{1}\left(\gamma_{0}^{N}(1), A_{0}^{N}\right)$. By construction, we also have

$$
\eta_{1}^{N}:=\int \Omega_{1}\left(\gamma_{0}^{N}(1), A_{0}^{N}\right)(d a) \eta_{1}^{(a)}+\frac{1}{\sqrt{N}} W_{1}^{N}=\Gamma_{1}^{2}\left(\gamma_{0}^{N}(1), \eta_{0}^{N}\right)+\frac{1}{\sqrt{N}} W_{1}^{N}
$$

with some local sampling random fields satisfying 1.6). Iterating this procedure, we define by induction a sequence of $N$-particle approximation measures

$$
\gamma_{n}^{N}(1)=\gamma_{n-1}^{N}(1) \eta_{n-1}^{N}\left(G_{n-1, \gamma_{n-1}^{N}}\right) \quad \text { and } \quad \eta_{n}^{N}:=\int A_{n}^{N}(d a) \eta_{n}^{(a)}
$$

with the occupation measure $A_{n}^{N}=\frac{1}{N} \sum_{i=1}^{N} \delta_{\mathrm{a}_{n}^{i}}$ associated with $N$ conditionally independent and identically distributed random variables $\mathrm{a}_{n}^{i}:=\left(a_{0, n}^{i}, a_{1, n}^{i}, \ldots, a_{n, n}^{i}\right)$ with common law $\Omega_{n}\left(\gamma_{n-1}^{N}(1), A_{n-1}^{N}\right)$. Arguing as above, we find that

$$
\eta_{n}^{N}=\int \Omega_{n}\left(\gamma_{n-1}^{N}(1), A_{n-1}^{N}\right)(d a) \eta_{n}^{(a)}+\frac{1}{\sqrt{N}} W_{n}^{N}=\Gamma_{n}^{2}\left(\gamma_{n-1}^{N}(1), \eta_{n-1}^{N}\right)+\frac{1}{\sqrt{N}} W_{n}^{N}
$$

with some local sampling random fields satisfying (1.6).

### 4.2.3 Convergence analysis

The main objective of this section is to show that $N$-particle occupation measures $A_{n}^{N}$ converge in a sense to be given, as $N$ tends to $\infty$, to the association probability measures $A_{n}$. To this end we observe that the one step mapping $\Omega_{n+1}$ introduced in 4.7) can be rewritten in the following form

$$
\Omega_{n+1}(m, A)(F)=\frac{A \mathcal{Q}_{n+1, m A}(F)}{A \mathcal{Q}_{n+1, m A}(1)}
$$

with the collection of integral operators $\mathcal{Q}_{n+1, m A}$ from $\mathcal{A}_{[0, n]}$ into $\mathcal{A}_{[0, n+1]}$ defined below

$$
\mathcal{Q}_{n+1, B}\left(a, d\left(a^{\prime}, b\right)\right):=\delta_{a}\left(d a^{\prime}\right) \nu_{n+1}(d b) \eta_{n}^{\left(a^{\prime}\right)}\left(\mathcal{G}_{n, B}^{(b)}\right) \quad \text { where } \quad \mathcal{G}_{n, B}^{(b)}:=G_{n, \int B(d a)}^{(b)} \eta_{n}^{(a)}
$$

with $B=m A$. In the above display $d\left(a^{\prime}, b\right)=d a^{\prime} \times d b$ stands for an infinitesimal neighborhood of the point $\left(a^{\prime}, b\right) \in \mathcal{A}_{[0, n+1]}$, with $a=\left(a_{0}^{\prime}, \ldots, a_{n}^{\prime}\right) \in \mathcal{A}_{[0, n]}$ and $b \in \mathcal{A}_{n+1}$, and $a=\left(a_{0}, \ldots, a_{n}\right) \in \mathcal{A}_{[0, n]}$. It is important to point out that

$$
B_{n}:=\gamma_{n}(1) \times A_{n} \Longrightarrow B_{n+1}=B_{n} \mathcal{Q}_{n+1, B_{n}}
$$

Notice that the flow of measures $\left(B_{n}\right)_{n \geq 0}$ satisfies the same type of equation as in (1.1), with the a total mass evolution of the same form as 1.3 :

$$
B_{n+1}(1)=B_{n}(1) A_{n}\left(\mathcal{G}_{n, B_{n}}\right) \quad \text { with } \quad \mathcal{G}_{n, m A}:=\int \nu_{n+1}(d b) \mathcal{G}_{n, m A}^{(b)}
$$

Notice that

$$
\left[\mathcal{Q}_{n+1, B}(F)-\mathcal{Q}_{n+1, B^{\prime}}(F)\right](a)=\int \nu_{n+1}(d b) \quad\left[\eta_{n}^{(a)}\left(\mathcal{G}_{n, B}^{(b)}\right)-\eta_{n}^{(a)}\left(\mathcal{G}_{n, B^{\prime}}^{(b)}\right)\right] F(a, b)
$$

If we set $B=m A$ and $B^{\prime}=m^{\prime} A^{\prime}$ then condition $\left(H_{2}^{\prime}\right)$ is met as long as

$$
\left|\eta_{n}^{(a)}\left(\mathcal{G}_{n, B}^{(b)}\right)-\eta_{n}^{(a)}\left(\mathcal{G}_{n, B^{\prime}}^{(b)}\right)\right| \leq c(n)\left|m-m^{\prime}\right|+\int\left|\left[A-A^{\prime}\right](\varphi)\right| \Sigma_{n, B^{\prime}}^{(b)}(d \varphi)
$$

for some collection of bounded measures $\Sigma_{n, B^{\prime}}^{(b)}$ on $\mathcal{B}\left(\mathcal{A}_{n}\right)$ such that $\int \operatorname{osc}(\varphi) \Sigma_{n, B^{\prime}}^{(b)} \leq$ $\delta\left(\Sigma_{n}^{(b)}\right)$, for some finite constant $\delta\left(\Sigma_{n}^{(b)}\right)<\infty$, whose values do dot depend on the parameters $(m, A) \in\left(I_{n} \times \mathcal{P}\left(\mathcal{A}_{n}\right)\right)$. Under the assumptions 4.4), we have $\mathcal{G}_{n, B}^{(b)}(x)=\alpha_{n}^{(b)}(B) G_{n}^{(b)}(x)$, for some collection of parameters $\alpha_{n}^{(b)}(B)$ satisfying

$$
\left|\alpha_{n}^{(b)}(B)-\alpha_{n}^{(b)}\left(B^{\prime}\right)\right| \leq c(n)\left|m-m^{\prime}\right|+\int\left|\left[A-A^{\prime}\right](\varphi)\right| \Sigma_{n, B^{\prime}}^{(b)}(d \varphi)
$$

This condition is clearly satisfied for the PHD model discussed in (4.6), as long as the functions $h_{n}\left(y_{n}\right)+d_{n} g_{n}\left(., y_{n}\right)$ are uniformly bounded from above and below.

For instance, for $b=y_{n} \notin\left\{c, c^{\prime}\right\}$ we have $\alpha_{n}^{(b)}(B)=\left[h_{n}(b)+\int B(d a) \eta_{n}^{(a)}\left(d_{n} g_{n}(., b)\right)\right]^{-1}$. In this case, we can check that

$$
\left|\alpha_{n}^{(b)}(B)-\alpha_{n}^{(b)}\left(B^{\prime}\right)\right| \leq c(n)\left|\left[B-B^{\prime}\right]\left(\varphi_{n}^{(b)}\right)\right| \quad \text { with } \quad \varphi_{n}^{(b)}(a):=\eta_{n}^{(a)}\left(d_{n} g_{n}(., b)\right)
$$

In the same way, we show that the condition $\left(H_{1}\right)$ is also met for the PHD model. This, by construction of $A_{n}^{N}$ we find that

$$
A_{n}^{N}=\Omega_{n}\left(\gamma_{n-1}^{N}(1), A_{n-1}^{N}\right)+\frac{1}{\sqrt{N}} \mathcal{W}_{n}^{N}
$$

with some local sampling random fields satisfying (1.6).

### 4.3 Mixed particle association models

We consider the association mapping

$$
\Omega_{n+1}:(m, A, \eta) \in(] 0, \infty\left[\times \mathcal{A}_{[0, n]} \times \mathcal{P}\left(E_{n}\right)^{\mathcal{A}_{[0, n]}}\right) \mapsto \Omega_{n+1}(m, A, \eta) \in \mathcal{P}\left(\mathcal{A}_{[0, n+1]}\right)
$$

defined for any $(m, A) \in(] 0, \infty\left[\times \mathcal{A}_{[0, n]}\right)$ and any $\eta: a \in \operatorname{Supp}(A) \mapsto \eta^{(a)} \in \mathcal{P}\left(E_{n}\right)$ by

$$
\Omega_{n+1}(m, A, \eta)(d(a, b)) \propto A(d a) \nu_{n+1}(d b) \eta^{(a)}\left(G_{n, m \int A(d a)}^{(b)} \eta^{(a)}\right)
$$

By construction, for any $A \in \mathcal{P}\left(\mathcal{A}_{[0, n-1]}\right)$, and $a \in \operatorname{Supp}(A) \mapsto \eta^{(a)} \in \mathcal{P}\left(E_{n-1}\right)$, we have

$$
\left.\Gamma_{n}^{2}\left(m, \int A(d a) \eta^{(a)}\right)\right)=\int \Omega_{n}\left(m, A, \eta^{(\cdot)}\right)(d(a, b)) \Phi_{n}^{(b)}\left(\eta^{(a)}\right)
$$

We also mention that the updating-prediction transformation defined in (4.5)

$$
\begin{equation*}
\Phi_{n}^{(a)}(\eta)=\Psi_{G_{n-1}}^{(a)}(\eta) M_{n}^{(a)}=\eta \mathcal{K}_{n, \eta}^{(a)} \quad \text { with } \quad \mathcal{K}_{n, \eta}^{(a)}=\mathcal{S}_{n-1, \eta}^{(a)} M_{n}^{(a)} \tag{4.8}
\end{equation*}
$$

In the above displayed formula $\mathcal{S}_{n, \eta}^{(a)}$ stands for some updating Markov transition from $E_{n-1}$ into itself satisfying the compatibility condition $\eta S_{n-1, \eta}^{(a)}=\Psi_{G_{n-1}}^{(a)}(\eta)$.

We let $A_{0}^{N}=\frac{1}{N} \sum_{i=1}^{N} \delta_{\mathrm{a}_{0}^{i}}$, be the empirical measure associated with $N$ independent and identically distributed random variables $\left(\mathrm{a}_{0}^{i}\right)_{1 \leq i \leq N}$ with common distribution $A_{0}$. For any $\mathrm{a} \in \mathcal{A}_{0}$, we let

$$
\eta_{0}^{N}:=\int A_{0}^{N}(d \mathrm{a}) \eta_{0}^{\left(\mathrm{a}, N^{\prime}\right)} \quad \text { and } \quad \eta_{0}^{\left(\mathrm{a}, N^{\prime}\right)}=\frac{1}{N^{\prime}} \sum_{i=1}^{N^{\prime}} \delta_{\xi_{0}^{[\mathrm{a}, j]}}
$$

with the empirical measure $\eta_{0}^{\left(\mathrm{a}, N^{\prime}\right)}$ associated with $N^{\prime}$ random variables $\xi_{0}^{[\mathrm{a}]}=\left(\xi_{0}^{[\mathrm{a}, j]}\right)_{1 \leq j \leq N^{\prime}}$ with common law $\eta_{0}^{(a)}$. We further assume that $\gamma_{0}(1)$ is known and set

$$
\gamma_{0}^{N}:=\gamma_{0}(1) \eta_{0}^{N} \quad \text { and } \quad \gamma_{1}^{N}(1):=\gamma_{0}^{N}(1) \eta_{0}^{N}\left(G_{0, \gamma_{0}^{N}}\right)
$$

It is readily checked that the fluctuation random fields given below

$$
\mathcal{W}_{0}^{\left(a, N^{\prime}\right)}=\sqrt{N^{\prime}}\left(\eta_{0}^{\left(a, N^{\prime}\right)}-\eta_{0}^{(a)}\right)
$$

satisfies (1.6), with $N=N^{\prime}$, for any given $a \in \mathcal{A}_{0}$. Using the fact that

$$
\int A_{0}^{N}(d \mathrm{a}) \eta_{0}^{\left(\mathrm{a}, N^{\prime}\right)}=\int A_{0}^{N}(d \mathrm{a}) \eta_{0}^{(\mathrm{a})}+\frac{1}{\sqrt{N^{\prime}}} \int A_{0}^{N}(d \mathrm{a}) \mathcal{W}_{0}^{\left(a, N^{\prime}\right)}
$$

we conclude that $\eta_{0}^{N}:=\eta_{0}+\frac{1}{\sqrt{N}} W_{0}^{N}$, with some local sampling random fields $W_{0}^{N}$ satisfying the same estimates as in 1.6 by replacing $1 / \sqrt{N}$ by the $\operatorname{sum}\left(1 / \sqrt{N}+1 / \sqrt{N^{\prime}}\right)$.

Using (4.8), for any $\mathrm{a}_{1}=\left(a_{0}, a_{1}\right)$ we find that

$$
\Phi_{1}^{\left(a_{1}\right)}\left(\eta_{0}^{\left(a_{0}, N^{\prime}\right)}\right)=\eta_{0}^{\left(a_{0}, N^{\prime}\right)} \mathcal{K}_{n, \eta_{0}^{\left(a_{0}, N^{\prime}\right)}}^{\left(a_{1}\right)}
$$

We let $A_{1}^{N}=\frac{1}{N} \sum_{i=1}^{N} \delta_{\mathrm{a}_{1}^{i}}$ be the occupation measure associated with $N$ conditionally independent and identically distributed random variables $\mathrm{a}_{1}^{i}:=\left(a_{0,1}^{i}, a_{1,1}^{i}\right)$ with common law

$$
\Omega_{1}\left(\gamma_{0}^{N}(1), A_{0}^{N}, \eta_{0}^{\left(\cdot, N^{\prime}\right)}\right)
$$

In the above displayed formula $\eta_{0}^{\left(\cdot, N^{\prime}\right)}$ stands for the mapping $a_{0} \in \mathcal{A}_{0} \mapsto \eta_{0}^{\left(a_{0}, N^{\prime}\right)} \in \mathcal{P}\left(E_{0}\right)$.
We consider a sequence of conditionally independent random variables $\xi_{1}^{\left[a_{0}, a_{1}, j\right]}$ with distribution $\mathcal{K}_{n, \eta_{0}^{\left(a_{0}, N^{\prime}\right)}}^{\left(a_{1}\right)}\left(\xi_{0}^{\left[a_{0}, j\right]},.\right)$, with $1 \leq j \leq N^{\prime}$, and we set

$$
\eta_{1}^{\left(\left(a_{0}, a_{1}\right), N^{\prime}\right)}=\frac{1}{N^{\prime}} \sum_{i=1}^{N^{\prime}} \delta_{\xi_{1}^{\left[\left(a_{0}, a_{1}\right), j\right]}} \quad \text { and } \quad \eta_{1}^{N}:=\int A_{1}^{N}(d \mathrm{a}) \eta_{1}^{\left(\mathrm{a}, N^{\prime}\right)}
$$

Arguing as before, for any given $\mathrm{a}_{1}:=\left(a_{0}, a_{1}\right) \in \operatorname{Supp}\left(A_{1}^{N}\right)$, the sequence of random fields

$$
\mathcal{W}_{1}^{\left(\mathrm{a}_{1} N^{\prime}\right)}:=\sqrt{N}\left(\eta_{1}^{\left(\left(a_{0}, a_{1}\right), N^{\prime}\right)}-\Phi_{1}^{\left(a_{1}\right)}\left(\eta_{0}^{\left(a_{0}, N^{\prime}\right)}\right)\right)
$$

satisfies (1.6), with $N=N^{\prime}$. Thus, we conclude that

$$
\eta_{1}^{N}=\Gamma_{1}^{2}\left(\gamma_{0}^{N}(1), \eta_{0}^{N}\right)+\frac{1}{\sqrt{N}} W_{1}^{N}
$$

with some local sampling random fields $W_{1}^{N}$ satisfying the same estimates as in 1.6 by replacing $1 / \sqrt{N}$ by the sum $\left(1 / \sqrt{N}+1 / \sqrt{N^{\prime}}\right)$. Iterating this procedure, we define by induction a sequence of $N$-particle approximation measures

$$
\gamma_{n}^{N}(1)=\gamma_{n-1}^{N}(1) \eta_{n-1}^{N}\left(G_{n-1, \gamma_{n-1}^{N}}\right) \quad \text { and } \quad \eta_{n}^{N}:=\int A_{n}^{N}(d a) \eta_{n}^{\left(a, N^{\prime}\right)}
$$

with the occupation measure $A_{n}^{N}=\frac{1}{N} \sum_{i=1}^{N} \delta_{\mathrm{a}_{n}^{i}}$ associated with $N$ conditionally independent and identically distributed random variables $\mathrm{a}_{n}^{i}:=\left(a_{0, n}^{i}, a_{1, n}^{i}, \ldots, a_{n, n}^{i}\right)$ with common law $\Omega_{n}\left(\gamma_{n-1}^{N}(1), A_{n-1}^{N}, \eta_{n-1}^{\left(\cdot, N^{\prime}\right)}\right)$. Arguing as above, we find that
$\eta_{n}^{N}=\int \Omega_{n}\left(\gamma_{n-1}^{N}(1), A_{n-1}^{N}, \eta_{n-1}^{\left(\cdot, N^{\prime}\right)}\right)(d(a, b)) \Phi_{n}^{(b)}\left(\eta_{n-1}^{\left(a, N^{\prime}\right)}\right)=\Gamma_{n}^{2}\left(\gamma_{n-1}^{N}(1), \eta_{n-1}^{N}\right)+\frac{1}{\sqrt{N}} W_{n}^{N}$
with some local sampling random fields satisfying the same estimates as in (1.6) by replacing $1 / \sqrt{N}$ by the sum $\left(1 / \sqrt{N}+1 / \sqrt{N^{\prime}}\right)$. As before, the $N$-particle occupation measures $A_{n}^{N}$ converge as $N$ tends to $\infty$ to the association probability measures $A_{n}$.

## 5 Appendix

### 5.1 Proof of corollary 3.6

For constant mappings $s_{n}=\mu_{n+1}(1)$, the mappings $\Phi_{n+1, \nu_{n}}^{1}$ and $\Phi_{n+1, m_{n}}^{2}$ are given by

$$
\Phi_{n+1, \nu_{n}}^{1}(u)=s_{n} \quad \text { and } \quad \Phi_{n+1, m_{n}}^{2}(\eta)=\Psi_{g_{n}^{(s)}}(\eta) M_{n+1}^{(s)}
$$

with the likelihood function $g_{n}^{(s)}$ and the Markov transitions $M_{n+1}^{(s)}$ defined in 2.7. Firstly, we observe that $r_{n}\left(s_{n}\right):=\sup _{x, x^{\prime} \in E_{n}} g_{n}^{(s)}(x) / g_{n}^{(s)}\left(x^{\prime}\right)$. We also notice that the second component mapping $\Phi_{n+1, m_{n}}^{2}$ does not depends on the parameter $m_{n}$, and it induces a FeynmanKac semigroup of the same form as the one discussed in section 3.2.

Under the premise of the proposition, the semigroup of associated with the Markov transitions $M_{n}$ satisfies the mixing property stated in the l.h.s. of (3.11) for some integer $m \geq 1$ and some parameter $\left.\left.\epsilon_{p}(m) \in\right] 0,1\right]$. In this situation, we also have that

$$
M_{p, p+m}^{(s)}(x, .) \geq \epsilon_{p}^{(s)}(m) M_{p, p+m}^{(s)}\left(x^{\prime}, .\right)
$$

with some positive parameter

$$
\epsilon_{p}^{(s)}(m) \geq \epsilon_{p}(m) / \prod_{p \leq k<p+m} r_{k}\left(s_{k}\right) r_{k}(1) \quad \text { and } \quad r_{n}\left(s_{n}\right):=\frac{s_{n} g_{n}^{+}+\left(1-s_{n}\right)}{s_{n} g_{n}^{-}+\left(1-s_{n}\right)}\left(\leq r_{n}(1)\right)
$$

To prove this claim, firstly we observe that $M_{p, p+m}^{(s)}(x,.) \ll M_{p, p+m}^{(s)-}(x,$.$) and$

$$
\prod_{p \leq k<p+m} r_{k}\left(s_{k}\right)^{-1} \leq d M_{p, p+m}^{(s)}(x, .) / d M_{p, p+m}^{(s)-}(x, .) \leq \prod_{p \leq k<p+m} r_{k}(1)
$$

with the semigroup $M_{p, n}^{(s)-}$ associated with the Markov transition

$$
M_{p, p+1}^{(s)-}(x, .)=\alpha_{p+1} M_{p+1}(x, .)+\left(1-\alpha_{k+1}\right) \bar{\mu}_{k+1} \quad \text { with } \quad \alpha_{p+1}:=\frac{s_{k} g_{k}^{-}}{s_{k} g_{k}^{-}+\left(1-s_{k}\right)}
$$

Using the geometric representation

$$
M_{p, n}^{(s)-}(x, .)=\left(\prod_{p<k \leq n} \alpha_{k}\right) M_{p, n}(x, .)+\sum_{p<k \leq n}\left(1-\alpha_{k}\right)\left(\prod_{k<l \leq n} \alpha_{l}\right) \bar{\mu}_{k} M_{k, n}
$$

it can be verified that

$$
M_{p, p+m}^{(s)-}(x, .) \geq \epsilon_{p}(m) M_{p, p+m}^{(s)-}\left(x^{\prime}, .\right) \geq \epsilon_{p}(m)\left(\prod_{p \leq k<p+m} g_{k}^{-} / g_{k}^{+}\right) M_{p, p+m}^{(s)}\left(x^{\prime}, .\right)
$$

from which we conclude that

$$
M_{p, p+m}^{(s)}(x, .) \geq \epsilon_{p}^{(s)}(m) M_{p, p+m}^{(s)}\left(x^{\prime}, .\right) \quad \text { with } \quad \epsilon_{p}^{(s)}(m) \geq \epsilon_{p}(m) / \prod_{p \leq k<p+m} r_{k}\left(s_{k}\right) r_{k}(1)
$$

We end the proof of the proposition combing the proposition 3.5 with the couple of estimates presented in (3.12) and (3.13). This ends the proof of the corollary.

### 5.2 Proof of theorem 3.7

The formulae presented in (2.6) can be rewritten in terms of matrix operations as follows

$$
\left[\gamma_{n+1}(1), 1-\gamma_{n+1}(1)\right]=\left[\widehat{\gamma}_{n}(1), 1-\widehat{\gamma}_{n}(1)\right]\left[\begin{array}{cc}
\Psi_{g_{n}}\left(\eta_{n}\right)\left(s_{n}\right) & 1-\Psi_{g_{n}}\left(\eta_{n}\right)\left(s_{n}\right) \\
\mu_{n+1}(1) & 1-\mu_{n+1}(1)
\end{array}\right]
$$

and

$$
\left[\widehat{\gamma}_{n}(1), 1-\widehat{\gamma}_{n}(1)\right]=\frac{\left[\gamma_{n}(1), 1-\gamma_{n}(1)\right]\left[\begin{array}{cc}
\eta_{n}\left(g_{n}\right) & 0 \\
0 & 1
\end{array}\right]}{\left[\gamma_{n}(1), 1-\gamma_{n}(1)\right]\left[\begin{array}{cc}
\eta_{n}\left(g_{n}\right) & 0 \\
0 & 1
\end{array}\right]\left[\begin{array}{l}
1 \\
1
\end{array}\right]}
$$

With a slight abuse of notation, we set

$$
\vartheta_{n}:=\left[\gamma_{n}(1), 1-\gamma_{n}(1)\right] \quad \widehat{\vartheta}_{n}:=\left[\widehat{\gamma}_{n}(1), 1-\widehat{\gamma}_{n}(1)\right] \quad \text { and } \quad 1=\left[\begin{array}{l}
1 \\
1
\end{array}\right]
$$

We also denote by $\mathcal{M}_{n+1, \eta_{n}}$ and $\mathcal{D}_{n, \eta_{n}}$ the stochastic and the diagonal matrices defined by

$$
\mathcal{M}_{n+1, \eta_{n}}:=\left[\begin{array}{cc}
\Psi_{g_{n}}\left(\eta_{n}\right)\left(s_{n}\right) & 1-\Psi_{g_{n}}\left(\eta_{n}\right)\left(s_{n}\right)  \tag{5.1}\\
\mu_{n+1}(1) & 1-\mu_{n+1}(1)
\end{array}\right] \quad \text { and } \quad \mathcal{D}_{n, \eta_{n}}:=\left[\begin{array}{cc}
\eta_{n}\left(g_{n}\right) & 0 \\
0 & 1
\end{array}\right]
$$

In this notation, the above recursion can be rewritten in a more compact form

$$
\vartheta_{n+1}=\widehat{\vartheta}_{n} \mathcal{M}_{n+1, \eta_{n}} \quad \text { and } \quad \widehat{\vartheta}_{n}=\frac{\vartheta_{n} \mathcal{D}_{n, \eta_{n}}}{\vartheta_{n} \mathcal{D}_{n, \eta_{n}} 1} \Longrightarrow \vartheta_{n+1}=\frac{\vartheta_{n} \mathcal{Q}_{n+1, \eta_{n}}}{\vartheta_{n} \mathcal{Q}_{n+1, \eta_{n}} 1}
$$

with the product of matrices $\mathcal{Q}_{n+1, \eta_{n}}=\mathcal{D}_{n, \eta_{n}} \mathcal{M}_{n+1, \eta_{n}}$.

$$
\forall u \in I_{p}(\subset[0,1]) \quad\left[\Phi_{p, n, \nu}^{1}(u), 1-\Phi_{p, n, \nu}^{1}(u)\right]=\frac{[u, 1-u] \mathcal{Q}_{p, n, \nu}}{[u, 1-u] \mathcal{Q}_{p, n, \nu}(1)}
$$

with the matrix semigroup

$$
\mathcal{Q}_{p, n, \nu}=\mathcal{Q}_{p+1, \nu_{p}} \mathcal{Q}_{p+2, \nu_{p+1}} \ldots \mathcal{Q}_{n, \nu_{n-1}}
$$

These semigroups are again of the same form as the Feynman-Kac models discussed in section 3.2 with a two point state space. When $\left.\mu_{n+1}(1) \in\right] 0,1\left[\right.$ and $0<s_{n}^{-} \leq s_{n}^{+}<1$, we have for any $n \geq 0$ and any $i, i^{\prime}, j \in\{1,2\}$

$$
\mathcal{M}_{n+1, \nu_{n}}(i, j) \geq \epsilon_{n} \mathcal{M}_{n+1, \nu_{n}}\left(i^{\prime}, j\right) \quad \text { and } \quad \sup _{i, i^{\prime} \in\{1,2\}} \frac{\mathcal{Q}_{n+1, \nu_{n}}(1)(i)}{\mathcal{Q}_{n+1, \nu_{n}}(1)\left(i^{\prime}\right)} \leq \delta_{n}^{\prime}(g)
$$

The first assertion is a direct consequence of the proposition 3.5 with the couple of estimates presented in $(3.12$ ) and $(3.13)$.

Using 2.3 , we find that $\Phi_{n+1, m_{n}}^{2}$ induces a Feynman-Kac models of the same form as the one discussed in section 3.2, More precisely, we have that

$$
\Phi_{n+1, m_{n}}^{2}(\eta)=\Psi_{\widehat{G}_{n, m_{n}}}(\eta) \widehat{M}_{n+1, m_{n}}
$$

with the potential functions $G_{n, m_{n}}$ and the Markov transitions $\widehat{M}_{n+1, m_{n}}$ defined in 2.4 and (2.5). Notice that

$$
\sup _{x, x^{\prime} \in E_{n}} \frac{\widehat{G}_{n, m_{n}}(x)}{\widehat{G}_{n, m_{n}}\left(x^{\prime}\right)} \leq \delta_{n}(s g)
$$

and for any $x \in E_{n}$ and any $n \geq 0$

$$
\delta_{n}(s g)^{-1} \widehat{M}_{n+1, m_{n}}^{-}(x, .) \leq \widehat{M}_{n+1, m_{n}}(x, .) \leq \delta_{n}(s g) \widehat{M}_{n+1, m_{n}}^{-}(x, .)
$$

with the Markov transitions $\widehat{M}_{n+1, m_{n}}^{-}$defined as $\widehat{M}_{n+1, m_{n}}^{-}$by replacing the functions $\left(s_{n}, g_{n}\right)$ by their lower bounds $\left(s_{n}^{-}, g_{n}^{-}\right)$. To prove this claim, we use the fact that for any positive function $f$ we have
$\frac{d \widehat{M}_{n+1, m_{n}}(f)}{d \widehat{M}_{n+1, m_{n}}^{-}(f)}=\frac{m_{n} g_{n}^{-} s_{n}^{-}+\left(1-m_{n}\right) \mu_{n+1}(1)}{m_{n} g_{n} s_{n}+\left(1-m_{n}\right) \mu_{n+1}(1)} \times \frac{m_{n} g_{n} s_{n} M_{n+1}(f)+\left(1-m_{n}\right) \mu_{n+1}(1) \bar{\mu}_{n+1}(f)}{m_{n} g_{n}^{-} s_{n}^{-}+\left(1-m_{n}\right) \mu_{n+1}(1) \bar{\mu}_{n+1}(f)}$
and the two series of inequalities

$$
\delta_{n}(s g)^{-1} \leq \frac{m_{n} g_{n}^{-} s_{n}^{-}+\left(1-m_{n}\right) \mu_{n+1}(1)}{m_{n} g_{n} s_{n}+\left(1-m_{n}\right) \mu_{n+1}(1)} \leq 1
$$

and

$$
1 \leq \frac{m_{n} g_{n} s_{n} M_{n+1}(f)+\left(1-m_{n}\right) \mu_{n+1}(1) \bar{\mu}_{n+1}(f)}{m_{n} g_{n}^{-} s_{n}^{-}+\left(1-m_{n}\right) \mu_{n+1}(1) \bar{\mu}_{n+1}(f)} \leq \delta_{n}(s g)
$$

With a slight abuse of notation, we write $\widehat{M}_{p, n}$, and respectively $\widehat{M}_{p, n}^{-}$, the semigroup associated with the Markov transitions $\widehat{M}_{n+1, m_{n}}$, and resp. $\widehat{M}_{n+1, m_{n}}^{-}$. Using the same argument as in the proof of corollary 3.6 it follows that

$$
\widehat{M}_{p, p+m}^{-}(x, .) \geq \epsilon_{p}(m) \widehat{M}_{p, p+m}^{-}\left(x^{\prime}, .\right)
$$

from which we conclude that

$$
\widehat{M}_{p, p+m}(x, .) \geq \widehat{\epsilon}_{p}(m) \widehat{M}_{p, p+m}\left(x^{\prime}, .\right) \quad \text { with } \quad \widehat{\epsilon}_{p}(m) \geq \epsilon_{p}(m) \prod_{0 \leq k<m} \delta_{p+k}(s g)^{-2}
$$

using proposition 3.5 with the couple of estimates presented in (3.12) and (3.13), we check that (3.3) is satisfied with

$$
a_{p, n}^{2} \leq 2 \rho_{p}(m) \prod_{k=0}^{\lfloor n / m\rfloor-1}\left(1-\epsilon_{p+k m}^{(m)}\right)
$$

and some parameters

$$
\epsilon_{p}^{(m)}:=\widehat{\epsilon}_{p}(m)^{2} \prod_{0<k<m} \delta_{p+k}(s g)^{-1} \geq \epsilon_{p}(m)^{2} \delta_{p}(s g)^{-4} \prod_{0<k<m} \delta_{p+k}(s g)^{-5}
$$

and

$$
\rho_{p}(m):=\widehat{\epsilon}_{p}(m)^{-1} \prod_{0 \leq k<m} \delta_{p+k}(s g) \leq \epsilon_{p}(m)^{-1} \prod_{0 \leq k<m} \delta_{p+k}(s g)^{3}
$$

This ends the proof of the first assertion of the theorem. Next, we discuss condition $(\operatorname{Cont}(\Phi))$. We observe that

$$
\Phi_{n+1, \nu}^{1}(u)=\frac{u \nu\left(g_{n} s_{n}\right)+(1-u) \mu_{n+1}(1)}{u \nu\left(g_{n}\right)+(1-u)}
$$

After some manipulations

$$
\begin{aligned}
& \Phi_{n+1, \nu}^{1}(u)-\Phi_{n+1, \nu^{\prime}}^{1}(u) \\
& =\frac{u \nu^{\prime}\left(g_{n}\right)}{u \nu^{\prime}\left(g_{n}\right)+(1-u)}\left[\Psi_{g_{n}}(\nu)-\Psi_{g_{n}}\left(\nu^{\prime}\right)\right]\left(s_{n}\right) \\
& \quad+\frac{u}{u \nu\left(g_{n}\right)+(1-u)} \frac{(1-u)}{u \nu^{\prime}\left(g_{n}\right)+(1-u)}\left[\Psi_{g_{n}}(\nu)\left(s_{n}\right)-\mu_{n+1}(1)\right]\left[\nu-\nu^{\prime}\right]\left(g_{n}\right)
\end{aligned}
$$

Recalling that the mapping $\theta_{a}(x)=a x /(a x+(1-x))$ in increasing on $[0,1]$ and using the fact that

$$
\Psi_{g_{n}}(\nu)=\nu S_{n, \nu} \Longrightarrow \Psi_{g_{n}}(\nu)-\Psi_{g_{n}}\left(\nu^{\prime}\right)=\frac{g_{n}^{+}}{\nu\left(g_{n}\right)}\left(\nu-\nu^{\prime}\right) S_{n, \nu^{\prime}}
$$

with the Markov transition

$$
S_{n, \nu^{\prime}}\left(x, d x^{\prime}\right)=\frac{g_{n}(x)}{g_{n}^{+}(x)} \delta_{x}\left(d x^{\prime}\right)+\left(1-\frac{g_{n}(x)}{g_{n}^{+}(x)}\right) \Psi_{g_{n}}\left(\nu^{\prime}\right)\left(d x^{\prime}\right)
$$

we prove

$$
\begin{equation*}
\left|\Psi_{g_{n}}(\nu)\left(s_{n}\right)-\Psi_{g_{n}}\left(\nu^{\prime}\right)\left(s_{n}\right)\right| \leq \frac{g_{n}^{+}}{g_{n}^{-}}\left|\left(\nu-\nu^{\prime}\right) S_{n, \nu^{\prime}}\left(s_{n}\right)\right| \tag{5.2}
\end{equation*}
$$

and for any $u \in I_{n}=\left[m_{n}^{-}, m_{n}^{+}\right]$

$$
\begin{aligned}
& \left|\Phi_{n+1, \nu}^{1}(u)-\Phi_{n+1, \nu^{\prime}}^{1}(u)\right| \\
& =\frac{m_{n}^{+} g_{n}^{+}}{m_{n}^{+} g_{n}^{+}+\left(1-m_{n}^{+}\right)} \frac{g_{n}^{+}}{g_{n}^{-}}\left|\left(\nu-\nu^{\prime}\right) S_{n, \nu^{\prime}}\left(s_{n}\right)\right| \\
& \quad+\frac{m_{n}^{+} g_{n}^{+}}{m_{n}^{+} g_{n}^{+}+\left(1-m_{n}^{+}\right)} \frac{\left(1-m_{n}^{-}\right)}{m_{n}^{-} g_{n}^{-}+\left(1-m_{n}^{-}\right)}\left\|s_{n}-\mu_{n+1}(1)\right\|\left|\left[\nu-\nu^{\prime}\right]\left(g_{n} / g_{n}^{-}\right)\right|
\end{aligned}
$$

This implies that

$$
\begin{aligned}
\tau_{n+1}^{1} \leq & \frac{m_{n}^{+} g_{n}^{+}}{m_{n}^{+} g_{n}^{+}+\left(1-m_{n}^{+}\right)} \frac{g_{n}^{+}}{g_{n}^{-}}\left(s_{n}^{+}-s_{n}^{-}\right) \\
& \quad+\frac{m_{n}^{+} g_{n}^{+}}{m_{n}^{+} g_{n}^{+}+\left(1-m_{n}^{+}\right)} \frac{\left(1-m_{n}^{-}\right)}{m_{n}^{-} g_{n}^{-}+\left(1-m_{n}^{-}\right)}\left\|s_{n}-\mu_{n+1}\right\|\left(\frac{g_{n}^{+}}{g_{n}^{-}}-1\right) \\
\leq & \frac{g_{n}^{+}}{g_{n}^{-}}\left[\left(s_{n}^{+}-s_{n}^{-}\right)+\left\|s_{n}-\mu_{n+1}(1)\right\|\right]
\end{aligned}
$$

Using (2.3) we also find that

$$
\Phi_{n+1, m}^{2}(\eta)(f)=\frac{m \eta\left(s_{n} g_{n} M_{n+1}(f)\right)+(1-m) \mu_{n+1}(f)}{m \eta\left(s_{n} g_{n}\right)+(1-m) \mu_{n+1}(1)}
$$

It is also readily check that

$$
\left[\Phi_{n+1, m}^{2}(\eta)-\Phi_{n+1, m^{\prime}}^{2}(\eta)\right](f)=\frac{\mu_{n+1}(1) \eta\left(g_{n} s_{n}\right)\left[\Psi_{g_{n} s_{n}}(\eta) M_{n+1}-\bar{\mu}_{n+1}\right](f)\left(m-m^{\prime}\right)}{\left[m \eta\left(s_{n} g_{n}\right)+(1-m) \mu_{n+1}(1)\right]\left[m^{\prime} \eta\left(s_{n} g_{n}\right)+\left(1-m^{\prime}\right) \mu_{n+1}(1)\right]}
$$

from which we conclude that

$$
\tau_{n+1}^{2} \leq \sup \left\{\frac{\mu_{n+1}}{s_{n}^{-} g_{n}^{-}}, \frac{s_{n}^{+} g_{n}^{+}}{\mu_{n+1}(1)}\right\} \leq \delta_{n}^{\prime}(g) \sup \left\{\frac{\mu_{n+1}}{s_{n}^{-}}, \frac{s_{n}^{+}}{\mu_{n+1}(1)}\right\}
$$

This ends the proof of the theorem.

### 5.3 Proof of proposition 3.3

The proof of proposition 3.3 is based on the following technical lemma.
Lemma 5.1 We assume that the regularity conditions $(\operatorname{Lip}(\Phi))$ and $(\operatorname{Cont}(\Phi))$ are satisfied. In this situation, for any $p \leq n, u, u^{\prime} \in I_{p}, \eta, \eta^{\prime} \in \mathcal{P}\left(E_{p}\right)$ and $f \in O s c_{1}\left(E_{n}\right)$ and any flow of masses and probability measures $m=\left(m_{n}\right)_{n \geq 0} \in \prod_{n \geq 0} I_{n}$ and $\nu:=\left(\nu_{n}\right)_{n \geq 0} \in \prod_{n \geq 0} \mathcal{P}\left(E_{n}\right)$ we have the following estimates

$$
\left|\Phi_{p, n, \nu^{\prime}}^{1}\left(u^{\prime}\right)-\Phi_{p, n, \nu}^{1}(u)\right| \leq a_{p, n}^{1}\left|u-u^{\prime}\right|+\sum_{p \leq q<n} \bar{a}_{q, n}^{1} \int\left|\left[\nu_{q}-\nu_{q}^{\prime}\right](\varphi)\right| \Omega_{q+1, \nu_{q}^{\prime}}^{1}(d \varphi)
$$

$$
\left|\Phi_{p, n, m^{\prime}}^{2}\left(\eta^{\prime}\right)(f)-\Phi_{p, n, m}^{2}(\eta)(f)\right| \leq a_{p, n}^{2} \int\left|\left[\eta-\eta^{\prime}\right](\varphi)\right| \Omega_{p, n, \eta^{\prime}}^{2}(f, d \varphi)+\sum_{p \leq q<n} \bar{a}_{q, n}^{2}\left|m_{q}-m_{q}^{\prime}\right|
$$

with the collection of parameters $\bar{a}_{p, n}^{i}, i=1,2$, defined in (3.6).

## Proof:

We use the decomposition

$$
\Phi_{p, n, \nu^{\prime}}^{1}\left(u^{\prime}\right)-\Phi_{p, n, \nu}^{1}(u)=\sum_{p \leq q \leq n}\left[\Phi_{q, n, \nu}^{1}\left(\Phi_{p, q, \nu^{\prime}}^{1}\left(u^{\prime}\right)\right)-\Phi_{q-1, n, \nu}^{1}\left(\Phi_{p, q-1, \nu^{\prime}}^{1}\left(u^{\prime}\right)\right)\right]
$$

and the fact that

$$
\begin{aligned}
\Phi_{q-1, n, \nu}^{1}\left(\Phi_{p, q-1, \nu^{\prime}}^{1}\left(u^{\prime}\right)\right) & =\Phi_{q, n, \nu}^{1}\left(\Phi_{q-1, q, \nu}^{1}\left[\Phi_{p, q-1, \nu^{\prime}}^{1}\left(u^{\prime}\right)\right]\right) \\
\Phi_{q, n, \nu}^{1}\left(\Phi_{p, q, \nu^{\prime}}^{1}\left(u^{\prime}\right)\right) & =\Phi_{q, n, \nu}^{1}\left(\Phi_{q-1, q, \nu^{\prime}}^{1}\left[\Phi_{p, q-1, \nu^{\prime}}^{1}\left(u^{\prime}\right)\right]\right)
\end{aligned}
$$

as well as

$$
\left|\Phi_{p, n, \nu}^{1}\left(u^{\prime}\right)-\Phi_{p, n, \nu}^{1}(u)\right| \leq a_{p, n}^{1}\left|u-u^{\prime}\right|
$$

and

$$
\left|\Phi_{q, n, \nu}^{1}\left(\Phi_{p, q, \nu^{\prime}}^{1}\left(u^{\prime}\right)\right)-\Phi_{q-1, n, \nu}^{1}\left(\Phi_{p, q-1, \nu^{\prime}}^{1}\left(u^{\prime}\right)\right)\right| \leq \bar{a}_{q-1, n}^{1} \int\left|\left[\nu_{q-1}-\nu_{q-1}^{\prime}\right](\varphi)\right| \Omega_{q, \nu_{q-1}^{\prime}}(d \varphi)
$$

to show that

$$
\left|\Phi_{p, n, \nu^{\prime}}^{1}\left(u^{\prime}\right)-\Phi_{p, n, \nu}^{1}(u)\right| \leq a_{p, n}^{1}\left|u-u^{\prime}\right|+\sum_{p<q \leq n} \bar{a}_{q-1, n}^{1} \int\left|\left[\nu_{q-1}-\nu_{q-1}^{\prime}\right](\varphi)\right| \Omega_{q, \nu_{q-1}^{\prime}}^{1}(d \varphi)
$$

The proof of the second assertion follows the same line of arguments. This ends the proof of the lemma.

Now we come to the proof of proposition 3.3.
Proof of proposition 3.3:
We fix a parameter $p \geq 0$, and we let $\left(m_{n}\right)_{n \geq p},\left(m_{n}^{\prime}\right)_{n \geq p} \in \prod_{n \geq p} I_{n}$ and $\left(\nu_{n}\right)_{n \geq p}$, and $\left(\nu_{n}^{\prime}\right)_{n \geq p} \in \prod_{n \geq p} \mathcal{P}\left(E_{n}\right)$ be defined by the following recursive formulae

$$
\begin{array}{llll}
\forall q>p & m_{q}^{\prime}=\Phi_{q, \nu_{q-1}^{\prime}}^{1}\left(m_{q-1}^{\prime}\right) & \text { and } & \nu_{q}^{\prime}=\Phi_{q, m_{q-1}^{\prime}}^{2}\left(\nu_{q-1}^{\prime}\right) \\
\forall q>p & m_{q}=\Phi_{q, \nu_{q-1}}^{1}\left(m_{q-1}\right) & \text { and } & \nu_{q}=\Phi_{q, m_{q-1}}^{2}\left(\nu_{q-1}\right)
\end{array}
$$

with the initial condition $\left(\nu_{p}, \nu_{p}^{\prime}\right)=\left(\eta, \eta^{\prime}\right)$ and $\left(m_{p}, m_{p}^{\prime}\right)=\left(u, u^{\prime}\right)$, for $q=p$.
By construction, we have $\left(\nu_{q}^{\prime}, \nu_{q}\right)=\left(\Phi_{p, q, m^{\prime}}^{2}\left(\eta^{\prime}\right), \Phi_{p, q, m}^{2}(\eta)\right)$ as well as $\left(m_{q}^{\prime}, m_{q}\right)=$ $\left(\Phi_{p, q, \nu^{\prime}}^{1}\left(u^{\prime}\right), \Phi_{p, q, \nu}^{1}(u)\right)$. In this case, using lemma 5.1 it follows that

$$
\begin{aligned}
& \left.\mid\left[\Gamma_{p, n}^{2}\left(m^{\prime}, \eta^{\prime}\right)-\Gamma_{p, n}^{2}(m, \eta)\right](f)\right] \\
& \leq a_{p, n}^{2} \int\left|\left[\eta-\eta^{\prime}\right](\varphi)\right| \Omega_{p, n, \eta^{\prime}}^{2}(f, d \varphi)+\sum_{p \leq q<n} \bar{a}_{q, n}^{2}\left|\Gamma_{p, q}^{1}\left(m^{\prime}, \eta^{\prime}\right)-\Gamma_{p, q}^{1}(m, \eta)\right|
\end{aligned}
$$

and

$$
\begin{aligned}
& \left|\Gamma_{p, n}^{1}\left(m^{\prime}, \eta^{\prime}\right)-\Gamma_{p, n}^{1}(m, \eta)\right| \\
& \leq a_{p, n}^{1}\left|m-m^{\prime}\right|+\sum_{p \leq q<n} \bar{a}_{q, n}^{1} \int\left|\left[\Gamma_{p, q}^{2}\left(m^{\prime}, \eta^{\prime}\right)-\Gamma_{p, q}^{2}(m, \eta)\right](\varphi)\right| \bar{\Omega}_{p, q, m^{\prime}, \eta^{\prime}}^{1}(d \varphi)
\end{aligned}
$$

with the probability measure $\bar{\Omega}_{p, q, m^{\prime}, \eta^{\prime}}^{1}=\Omega_{q+1, \Gamma_{p, q}^{2}\left(m^{\prime}, \eta^{\prime}\right)}^{1}$. Combining these two estimates, we arrive at the following inequality

$$
\begin{aligned}
& \left.\mid\left[\Gamma_{p, n}^{2}\left(m^{\prime}, \eta^{\prime}\right)-\Gamma_{p, n}^{2}(m, \eta)\right](f)\right] \\
& \leq a_{p, n}^{2} \int\left|\left[\eta-\eta^{\prime}\right](\varphi)\right| \Omega_{p, n, \eta^{\prime}}^{2}(f, d \varphi)+\left[\sum_{p \leq q<n} a_{p, q}^{1} \bar{a}_{q, n}^{2}\right]\left|m-m^{\prime}\right| \\
& +\sum_{p \leq r<q<n} \bar{a}_{r, q}^{1} \bar{a}_{q, n}^{2} \int\left|\left[\Gamma_{p, r}^{2}\left(m^{\prime}, \eta^{\prime}\right)-\Gamma_{p, r}^{2}(m, \eta)\right](\varphi)\right| \bar{\Omega}_{p, r, m^{\prime}, \eta^{\prime}}^{1}(d \varphi)
\end{aligned}
$$

This implies that

$$
\begin{aligned}
& \left.\mid\left[\Gamma_{p, n}^{2}\left(m^{\prime}, \eta^{\prime}\right)-\Gamma_{p, n}^{2}(m, \eta)\right](f)\right] \\
& \leq b_{p, n}^{\prime}\left|m-m^{\prime}\right|+a_{p, n}^{2} \int\left|\left[\eta-\eta^{\prime}\right](\varphi)\right| \Omega_{p, n, \eta^{\prime}}^{2}(f, d \varphi) \\
& +\sum_{p \leq r_{1}<n} b_{r_{1}, n} \int\left|\left[\Gamma_{p, r_{1}}^{2}\left(m^{\prime}, \eta^{\prime}\right)-\Gamma_{p, r_{1}}^{2}(m, \eta)\right](\varphi)\right| \bar{\Omega}_{p, r_{1}, m^{\prime}, \eta^{\prime}}^{1}(d \varphi)
\end{aligned}
$$

Our next objective is to show that

$$
\begin{aligned}
& \left.\mid\left[\Gamma_{p, n}^{2}\left(m^{\prime}, \eta^{\prime}\right)-\Gamma_{p, n}^{2}(m, \eta)\right](f)\right] \\
& \leq \alpha_{p, n}^{k}\left|m-m^{\prime}\right|+\beta_{p, n}^{k} \int\left|\left[\eta-\eta^{\prime}\right](\varphi)\right| \Theta_{p, n, \eta^{\prime}}^{k}(f, d \varphi) \\
& +\sum_{p \leq r_{1}<r_{2}<\ldots<r_{k}<n}\left(\prod_{1 \leq l \leq k} b_{r_{l}, r_{l+1}}\right) \int\left|\left[\Gamma_{p, r_{1}}^{2}\left(m^{\prime}, \eta^{\prime}\right)-\Gamma_{p, r_{1}}^{2}(m, \eta)\right](\varphi)\right| \bar{\Omega}_{p, r_{1}, m^{\prime}, \eta^{\prime}}^{1}(d \varphi)
\end{aligned}
$$

with the convention $r_{k+1}=n$, for any $k \leq(n-p)$, some Markov transitions $\Theta_{p, n, m^{\prime} \eta^{\prime}}^{k}(f, d \varphi)$ and the parameters

$$
\begin{aligned}
& \alpha_{p, n}^{k}=b_{p, n}^{\prime}+\sum_{l=1}^{k-1} \sum_{p \leq r_{1}<\ldots r_{l}<n} b_{p, r_{1}}^{\prime} \prod_{1 \leq k \leq l} b_{r_{k}, r_{k+1}} \\
& \beta_{p, n}^{k}=a_{p, n}^{2}+\sum_{l=1}^{k-1} \sum_{p \leq r_{1}<\ldots . r_{l}<n} a_{p, r_{1}}^{2} \prod_{1 \leq k \leq l} b_{r_{k}, r_{k+1}}, \quad \text { with the convention } r_{l+1}=n .
\end{aligned}
$$

We proceed by induction on the parameter $k$. Firstly, we observe that the result is satisfied for $k=1$ with $\left(\alpha_{p, n}^{1}, \beta_{p, n}^{1}\right)=\left(b_{p, n}^{\prime}, a_{p, n}^{2}\right)$ and $\Theta_{p, n, \eta^{\prime}}^{1}=\Omega_{p, n, \eta^{\prime}}^{2}$. We further assume that the result is satisfied at rank $k$. In this situation, using the fact that

$$
\begin{aligned}
& \left.\mid\left[\Gamma_{p, r_{1}}^{2}\left(m^{\prime}, \eta^{\prime}\right)-\Gamma_{p, r_{1}}^{2}(m, \eta)\right](\varphi)\right] \\
& \leq b_{p, r_{1}}^{\prime}\left|m-m^{\prime}\right|+a_{p, r_{1}}^{2} \int\left|\left[\eta-\eta^{\prime}\right]\left(\varphi^{\prime}\right)\right| \Omega_{p, r_{1}, \eta^{\prime}}^{2}\left(\varphi, d \varphi^{\prime}\right) \\
& +\sum_{p \leq r_{0}<r_{1}} b_{r_{0}, r_{1}} \int\left|\left[\Gamma_{p, r_{0}}^{2}\left(m^{\prime}, \eta^{\prime}\right)-\Gamma_{p, r_{0}}^{2}(m, \eta)\right](\varphi)\right| \bar{\Omega}_{p, r_{0}, m^{\prime}, \eta^{\prime}}^{1}(d \varphi)
\end{aligned}
$$

we conclude that

$$
\begin{aligned}
& \left.\mid\left[\Gamma_{p, n}^{2}\left(m^{\prime}, \eta^{\prime}\right)-\Gamma_{p, n}^{2}(m, \eta)\right](f)\right] \\
& \leq \alpha_{p, n}^{k+1}\left|m-m^{\prime}\right|+\beta_{p, n}^{k+1} \int\left|\left[\eta-\eta^{\prime}\right](\varphi)\right| \Theta_{p, n, m^{\prime} \eta^{\prime}}^{k+1}(f, d \varphi) \\
& +\sum_{p \leq r_{0}<r_{1}<r_{2}<\ldots<r_{k}<n} b_{r_{0}, r_{1}} b_{r_{1}, r_{2}} \ldots b_{r_{k}, n} \int\left|\left[\Gamma_{p, r_{0}}^{2}\left(m^{\prime}, \eta^{\prime}\right)-\Gamma_{p, r_{0}}^{2}(m, \eta)\right](\varphi)\right| \bar{\Omega}_{p, r_{0}, m^{\prime}, \eta^{\prime}}^{1}(d \varphi) \\
& \text { with } \\
& \alpha_{p, n}^{k+1}=\alpha_{p, n}^{k}+\sum_{p \leq r_{1}<r_{2}<\ldots<r_{k}<n} b_{p, r_{1}}^{\prime} b_{r_{1}, r_{2}} \ldots b_{r_{k}, n} \\
& \beta_{p, n}^{k+1}=\beta_{p, n}^{k}+\sum_{p \leq r_{1}<r_{2}<\ldots<r_{k}<n} a_{p, r_{1}}^{2} b_{r_{1}, r_{2}} \ldots b_{r_{k}, n}
\end{aligned}
$$

and the Markov transition

$$
\begin{aligned}
\beta_{p, n}^{k+1} \Theta_{p, n, m^{\prime} \eta^{\prime}}^{k+1}(f, d \varphi)= & \beta_{p, n}^{k} \Theta_{p, n, \eta^{\prime}}^{k}(f, d \varphi) \\
& +\sum_{p \leq r_{1}<r_{2}<\ldots<r_{k}<n} a_{p, r_{1}}^{2} b_{r_{1}, r_{2}} \ldots b_{r_{k}, n}\left(\bar{\Omega}_{p, r_{1}, m^{\prime}, \eta^{\prime}}^{1} \Omega_{p, r_{1}, \eta^{\prime}}^{2}\right)(d \varphi)
\end{aligned}
$$

We end the proof of the proposition using the fact that

$$
\begin{aligned}
\left|\Gamma_{p, n}^{1}\left(m^{\prime}, \eta^{\prime}\right)-\Gamma_{p, n}^{1}(m, \eta)\right| \leq & {\left[a_{p, n}^{1}+\sum_{p \leq q<n} c_{p, q}^{2,1} \bar{a}_{q, n}^{1}\right]\left|m-m^{\prime}\right| } \\
& +\sum_{p \leq q<n} \bar{a}_{q, n}^{1} c_{p, q}^{2,2} \int\left|\left[\eta-\eta^{\prime}\right]\left(\varphi^{\prime}\right)\right|\left[\bar{\Omega}_{p, q, m^{\prime}, \eta^{\prime}}^{1} \Theta_{p, q, \eta^{\prime}}\right]\left(d \varphi^{\prime}\right)
\end{aligned}
$$

This proof of the proposition is now completed.

### 5.4 Proof of theorem 3.11

For any $\eta \in \mathcal{P}(E)$ and any $u, u^{\prime} \in I_{n}$, we have

$$
\left|\Phi_{n+1, \eta}^{1}(u)-\Phi_{n+1, \eta}^{1}\left(u^{\prime}\right)\right| \leq\left|u-u^{\prime}\right|\left[r(1-d)+r d h \mathcal{Y}_{n}\left(\frac{g^{+}}{\left.\left[h+d m^{-} g^{-}\right)\right]^{2}}\right)\right]
$$

This implies that condition $(3.2$ is satisfied with

$$
a_{n, n+1}^{1} \leq r(1-d)+r d h \mathcal{Y}_{n}\left(\frac{g^{+}}{\left.\left[h+d m^{-} g^{-}\right)\right]^{2}}\right)
$$

In the same way, for any $\eta, \eta^{\prime} \in \mathcal{P}(E)$ and any $u \in I_{n}$, we have

$$
\Phi_{n+1, \eta}^{1}(u)-\Phi_{n+1, \eta^{\prime}}^{1}(u)=r d h u \int \mathcal{Y}_{n}(d y) \frac{1}{\left[h+\operatorname{du\eta (g(\cdot ,y))][h+du\eta ^{\prime }(g(\cdot ,y))]}\left(\eta-\eta^{\prime}\right)(g(., y)), ~(g)\right.}
$$

This shows 3.4 with $\tau_{n+1}^{1} \leq r d h m^{+} \mathcal{Y}_{n}\left(\frac{g^{+}-g^{-}}{\left[h+d m^{-} g^{-}\right]^{2}}\right)$, and the probability measures

$$
\Omega_{n, \eta^{\prime}}^{1}(d \varphi) \propto \int \mathcal{Y}_{n}(d y) \frac{g^{+}(y)-g^{-}(y)}{\left[h+d m^{-} g^{-}(y)\right]^{2}} \delta_{\frac{g(\cdot, y)}{g^{+}(y)-g^{-}(y)}}(d \varphi)
$$

Now, we come to the analysis of the mappings

$$
\Phi_{n+1, u}^{2}(\eta) \propto r(1-d) u \eta M+\int \mathcal{Y}_{n}(d y) w_{u}(\eta, y) \Psi_{g(\cdot, y)}(\eta) M+\mu(1) \bar{\mu}
$$

with the weight functions

$$
w_{u}(\eta, y):=\frac{r d u \eta(g(., y))}{h+\operatorname{du\eta }(g(., y))}=r\left(1-\frac{h}{h+\operatorname{du\eta }(g(., y))}\right)
$$

Notice that

$$
w^{-}(y):=\frac{r d m^{-} g^{-}(y)}{h+d m^{-} g^{-}(y)} \leq w_{u}(\eta, y) \leq w^{+}(y):=\frac{r d m^{+} g^{+}(y)}{h+d m^{+} g^{+}(y)}
$$

To have a more synthetic formula, we extend the observation state space with two auxiliary points $c_{1}, c_{2}$ and we set

$$
\mathcal{Y}_{n}^{c}=\mathcal{Y}_{n}+\delta_{c_{1}}+\delta_{c_{2}}
$$

we extend the likelihood and the weight functions by setting $g\left(x, c_{1}\right)=g\left(x, c_{2}\right)=1$, and

$$
\begin{aligned}
w^{-}\left(c_{1}\right) & :=r(1-d) m^{-} \leq w_{u}\left(\eta, c_{1}\right):=r(1-d) u \leq w^{+}\left(c_{1}\right):=r(1-d) m^{+} \\
w_{u}\left(\eta, c_{2}\right) & =w^{+}\left(c_{2}\right)=w^{-}\left(c_{2}\right):=\mu(1)
\end{aligned}
$$

In this notation, we find that

$$
\Phi_{n+1, u}^{2}(\eta) \propto \int \mathcal{Y}_{n}^{c}(d y) w_{u}(\eta, y) \Psi_{g(\cdot, y)}(\eta) M_{y}
$$

with the collection of Markov transitions $M_{y}$ defined below

$$
\forall y \notin\left\{c_{2}\right\} \quad M_{y}=M \quad \text { and } \quad M_{c_{2}}=\bar{\mu}
$$

Notice that the normalizing constants $\mathcal{Y}_{n}^{c}\left(w_{u}(\eta,).\right)$ satisfy the following lower bounds

$$
\mathcal{Y}_{n}^{c}\left(w_{u}(\eta, .)\right) \geq \mathcal{Y}_{n}^{c}\left(w^{-}\right)=r(1-d) m^{-}+\mathcal{Y}_{n}\left(w^{-}\right)+\mu(1)
$$

We analyze the Lipschitz properties of the mappings $\Phi_{n+1, u}^{2}$ using the following decomposition

$$
\Phi_{n+1, u}^{2}(\eta)-\Phi_{n+1, u}^{2}\left(\eta^{\prime}\right)=\Delta_{n+1, u}\left(\eta, \eta^{\prime}\right)+\Delta_{n+1, u}^{\prime}\left(\eta, \eta^{\prime}\right)
$$

with the signed measures

$$
\Delta_{n+1, u}\left(\eta, \eta^{\prime}\right)=\int \mathcal{Y}_{n}^{c}(d y) \frac{w_{u}(\eta, y)}{\mathcal{Y}_{n}^{c}\left(w_{u}(\eta, .)\right)}\left[\Psi_{g(\cdot, y)}(\eta) M_{y}-\Psi_{g(\cdot, y)}\left(\eta^{\prime}\right) M_{y}\right]
$$

and

$$
\Delta_{n+1, u}^{\prime}\left(\eta, \eta^{\prime}\right)=\frac{1}{\mathcal{Y}_{n}^{c}\left(w_{u}(\eta, .)\right)} \int \mathcal{Y}_{n}^{c}(d y)\left[w_{u}(\eta, y)-w_{u}\left(\eta^{\prime}, y\right)\right]\left(\Psi_{g(\cdot, y)}\left(\eta^{\prime}\right) M_{y}-\Phi_{n+1, u}^{2}\left(\eta^{\prime}\right)\right)
$$

Arguing as in the proof of theorem 3.7 given in the appendix (see for instance (5.2) , one checks that

$$
\begin{aligned}
& \left|\Delta_{n+1, u}\left(\eta, \eta^{\prime}\right)(f)\right| \\
& \leq \frac{1}{\mathcal{Y}_{n}^{c}\left(w^{-}\right)}\left(r(1-d) m^{+}\left|\left(\eta-\eta^{\prime}\right)(M(f))\right|+\int \mathcal{Y}_{n}(d y) w^{+}(y) \frac{g^{+}(y)}{g^{-}(y)}\left|\left(\eta-\eta^{\prime}\right)\left(S_{\eta^{\prime}}^{y} M(f)\right)\right|\right)
\end{aligned}
$$

for some collection of Markov transitions $S_{\eta^{\prime}}^{y}$ from $E$ into itself. It is also readily checked that

$$
\left|\Delta_{n+1, u}^{\prime}\left(\eta, \eta^{\prime}\right)(f)\right| \leq \frac{h r d m^{+}}{\mathcal{Y}_{n}^{c}\left(w^{-}\right)} \int \mathcal{Y}_{n}(d y) \frac{1}{\left(h+m^{-} d g^{-}(y)\right)^{2}}\left|\left(\eta-\eta^{\prime}\right)(g(., y))\right|
$$

This clearly implies that condition (3.3) is satisfied with

$$
a_{n, n+1}^{2} \leq \frac{1}{\mathcal{Y}_{n}^{c}\left(w^{-}\right)}\left(\beta(M)\left[r(1-d) m^{+}+\mathcal{Y}_{n}\left(\frac{w^{+} g^{+}}{g^{-}}\right)\right]+h r d m^{+} \mathcal{Y}_{n}\left(\frac{g^{+}-g^{-}}{\left(h+m^{-} d g^{-}\right)^{2}}\right)\right)
$$

We analyze the continuity properties of the mappings $u \mapsto \Phi_{n+1, u}^{2}(\eta)$ using the following decomposition

$$
\begin{aligned}
& \Phi_{n+1, u}^{2}(\eta)-\Phi_{n+1, u^{\prime}}^{2}(\eta) \\
& =\frac{1}{\mathcal{Y}_{n}^{c}\left(w_{u}(\eta, .)\right)} \int \mathcal{Y}_{n}^{c}(d y)\left[w_{u}(\eta, y)-w_{u^{\prime}}(\eta, y)\right]\left(\Psi_{g(\cdot, y)}(\eta) M_{y}-\Phi_{n+1, u^{\prime}}^{2}(\eta)\right)
\end{aligned}
$$

This implies that

$$
\left|\left[\Phi_{n+1, u}^{2}(\eta)-\Phi_{n+1, u^{\prime}}^{2}(\eta)\right](f)\right| \leq \frac{1}{\mathcal{Y}_{n}^{c}\left(w^{-}\right)}\left[r(1-d)+h r d \mathcal{Y}_{n}\left(\frac{g^{+}}{\left(h+d m^{-} g^{-}\right)^{2}}\right)\right]\left|u-u^{\prime}\right|
$$

This shows that condition (3.5) is satisfied with

$$
\tau_{n+1}^{2} \leq \frac{1}{\mathcal{Y}_{n}^{c}\left(w^{-}\right)}\left[r(1-d)+h r d \mathcal{Y}_{n}\left(\frac{g^{+}}{\left(h+d m^{-} g^{-}\right)^{2}}\right)\right]
$$

This ends the proof of the theorem.

## References

[1] Y. Bar-Shalom and T. Fortmann, Tracking and Data Association. Academic Press, San Diego, 1988.
[2] S. Blackman, Multiple Target Tracking with Radar Applications. Artech House, Norwood, 1986.
[3] Clark, D.E., Bell, J., Convergence results for the particle PHD filter, Signal Processing, IEEE Transactions, Volume 54, Issue 7, pp. 2652-2661, (2006).
[4] Fr. Caron, P. Del Moral, A. Doucet, and M. Pace On the Conditional Distributions of Spatial Point Processes HAL-INRIA RR-7232, March 2010.
[5] P. Del Moral, A. Doucet and S. S. Singh A Backward Particle Interpretation of Feynman-Kac Formulae M2AN, vol 44, no. 5, pp. 947-976, (2010).
[6] Del Moral, P., Feynman-Kac formulae. Genealogical and interacting particle systems with applications, Probability and its Applications, Springer Verlag, New York (2004).
[7] Del Moral P. and Doucet, A. Particle Motions in Absorbing Medium with Hard and Soft Obstacles. Stochastic Analysis and Applications, vol. 22, no. 5, (2004).
[8] P. Del Moral, Fr. Patras, S. Rubenthaler A Mean Field Theory of Nonlinear Filtering HAL-INRIA RR-6437 [39p], (Feb. 2008). Oxford Handbook of Nonlinear Filtering, Eds D. Crisan and B. Rosovsky (2010).
[9] P. Del Moral, L. Miclo Particle approximations of Lyapunov exponents connected to Schrödinger operators and Feynman?Kac semigroups ESAIM: P\&S 7 171-208 (2003).
[10] Doucet A., de Freitas J.F., Gordon N.J., Sequential Monte-Carlo Methods in Practice, Springer Verlag New York (2001).
[11] E.B. Dynkin. An Introduction to Branching Measure-Valued Processes, CRM Monograph Series 6, American Mathematical Society, Providence, RI, 1994.
[12] I. Goodman, R. Mahler, and H. Nguyen. Mathematics of Data Fusion, Kluwer Academic Publishers (1997).
[13] J. Goutsias, R. Mahler, and H. Nguyen (eds.). Random Sets Theory and Applications, Springer-Verlag New York (1997).
[14] A.M. Johansen, S.S. Singh, A. Doucet, and B.N. Vo Convergence of the sequential Monte Carlo implementation of the PHD filter. Methodology and Computing in Applied Probability, vol. 8, no. 2, pp. 265-291, 2006.
[15] M. L. Kleptsyna and A. Y. Veretennikov. On discrete time ergodic filters with wrong initial data. Probab. Theory Related Fields, 141(3-4):411?444, 2008.
[16] R. Mahler. A theoretical foundation for the Stein-Winter Probability Hypothesis Density (PHD) multi-target tracking approach, Proc. MSS Nat'l Symp. on Sensor and Data Fusion, Vol. I, San Antonio TX 2000.
[17] R. Mahler. Multi-target Bayes filtering via first-order multi-target moments, IEEE Trans. Aerospace $\mathfrak{E}$ Electronic Systems, Vol. 39, No. 4, pp. 1152-1178 (2003).
[18] R. Mahler, Statistical Multisource-Multitarget Information Fusion. Artech House, (2007).
[19] S. Mori, C. Chong, E. Tse and R. Wishner. Tracking and identifying multiple targets without apriori identifications, IEEE Trans. Automatic Control, vol. 21, pp. 401-409, (1986).
[20] S. Sarkka, A. Vehtari, and J. Lampinen. Rao-Blackwellized Monte Carlo data association for multiple target tracking. Proceedings of FUSION 2004 The 7th International Conference on Information Fusion, Stockholm (2004).
[21] S. Sarkka, A. Vehtari, and J. Lampinen. Rao-Blackwellized Particle Filter for Multiple Target Tracking. Information Fusion Journal, Volume 8, Issue 1, Pages 2-15, (2007).
[22] S.S. Singh, B.N. Vo, A. Baddeley and S. Zuyev. Filters for spatial point processes. SIAM J. Control Optim., vol. 48, no. 4, pp. 2275-2295, (2009).
[23] Sidenbladh H., Multi-target particle filtering for the Probability Hypothesis Density, Proc. Int'l Conf. on Information Fusion, Cairns, Australia, pp. 800?806, (2003).
[24] B.-N Vo, S.S. Singh, and Doucet A.. Sequential Monte Carlo methods for Bayesian Multi-target filtering with Random Finite Sets, IEEE Trans. Aerospace Elec. Systems, vol. 41, pp. 1224-1245 (2005).
[25] Vo, B.-N.; Ma, W.-K., A closed-form solution for the probability hypothesis density filter, Proc. Information Fusion, 8th International Conference on Volume 2, 25-28 (2005).
[26] B.-N. Vo, and W.-K. Ma, "The Gaussian mixture Probability Hypothesis Density filter," IEEE Trans. Signal Processing, IEEE Trans. Signal Processing, Vol. 54, No. 11, pp. 4091-4104, 2006.
[27] B.-T. Vo, Random Finite Sets in Multi-Object Filtering, PhD Thesis, University of Western Australia, 2008.
[28] William Ng, Jack Li, Simon Godsill, and Jaco Vermaak, Tracking variable number of targets using Sequential Monte Carlo Methods. Proceedings of the IEEE Statistical Signal Processing, (2005).
[29] William Ng, Jack Li, Simon Godsill, and Jaco Vermaak. A review of recent results in multiple target tracking. Proceedings of the International Symposium on Image and Signal Processing and Analysis, (2005).
[30] William Ng, Jack Li, Simon Godsill, and Jaco Vermaak. Tracking variable number of targets using Sequential Monte Carlo Method, Proceedings of the European Signal Processing Conference (2005).
[31] Zajic T., Ravichandran R., Mahler R., Mehra R., and Noviskey M., Joint tracking and identification with robustness against unmodeled targets, Signal Processing, Sensor Fusion and Target Recognition XII, SPIE Proc., vol. 5096, pp. 279?290, (2003).


[^0]:    ${ }^{*}$ The work of F. Caron, P. Del Moral, and M. Pace is supported by the French Agence Nationale de la Recherche under the project Propagation 2009. The work of B.-N. Vo is supported by Australian Research Council under the Future Fellowship FT0991854
    ${ }^{\dagger}$ Centre INRIA Bordeaux et Sud-Ouest \& Institut de Mathématiques de Bordeaux , Université Bordeaux, 351 cours de la Libération 33405 Talence cedex, France, Francois.Caron@inria.fr, Pierre.Del-Moral@inria.fr, Michele.Pace@inria.fr
    ${ }^{\ddagger}$ School of Electrical Electronic \& Computer Engineering M018 The University of Western Australia, ba-ngu.vo@uwa.edu.au.

